

laminar Geschw. prof.

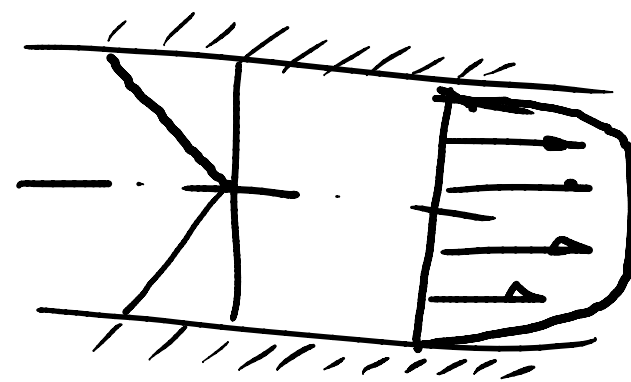
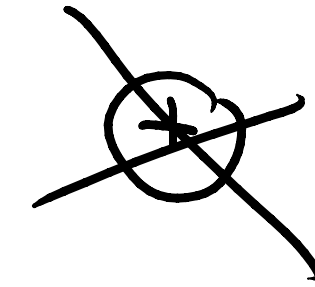
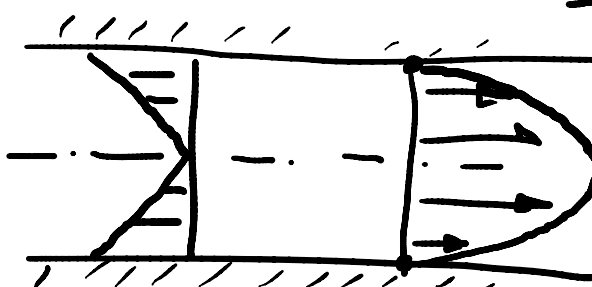
nicht l. l.

$$\rho \frac{D\vec{u}}{Dt} = -\nabla p + \eta \Delta \vec{u}$$

$\equiv 0$, für linear.

quasistationäre, laminare Schichtstr.

Schleppströmung



$$\Delta p = \eta \Delta \vec{u}$$



↳ Superposition.



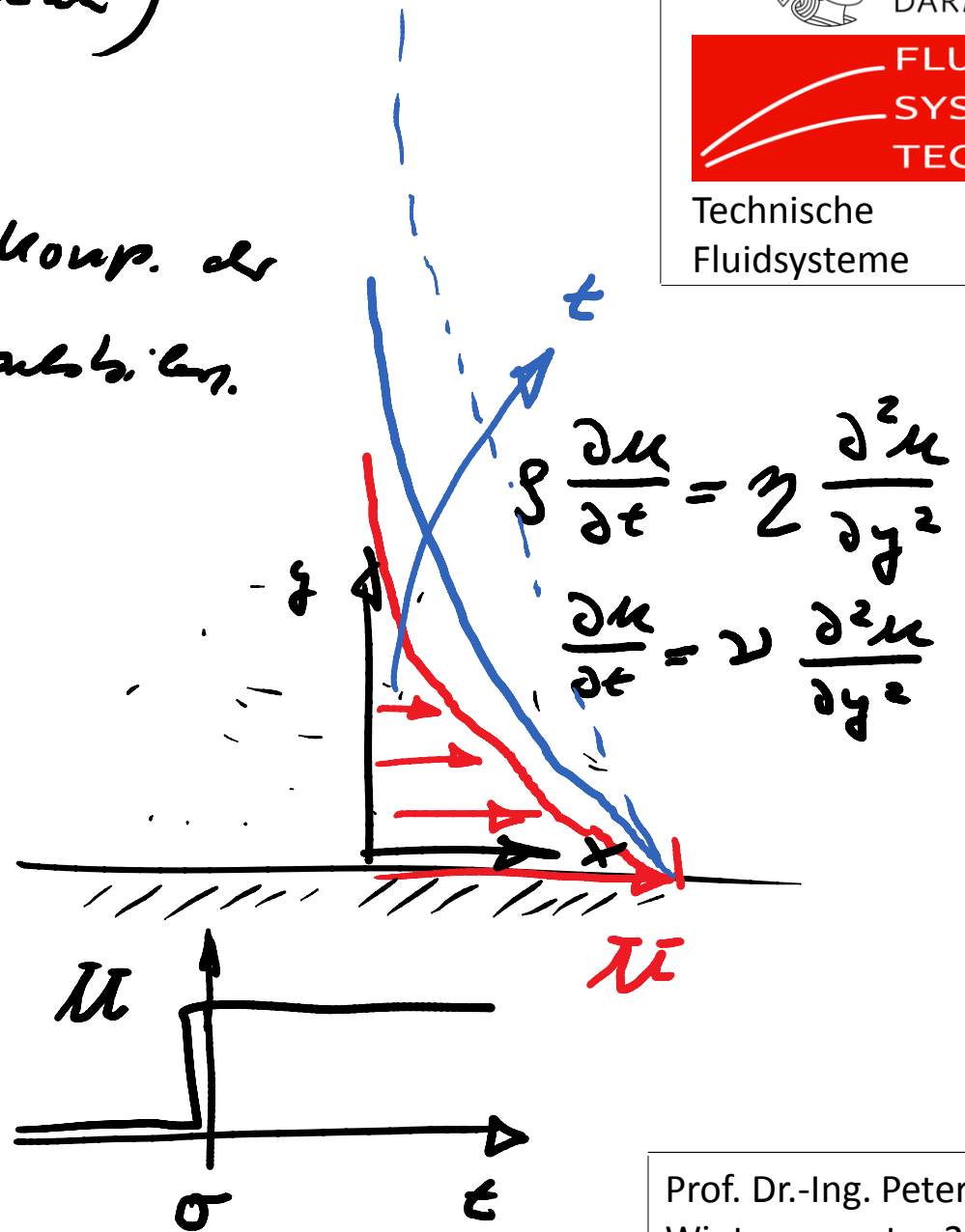
Druck

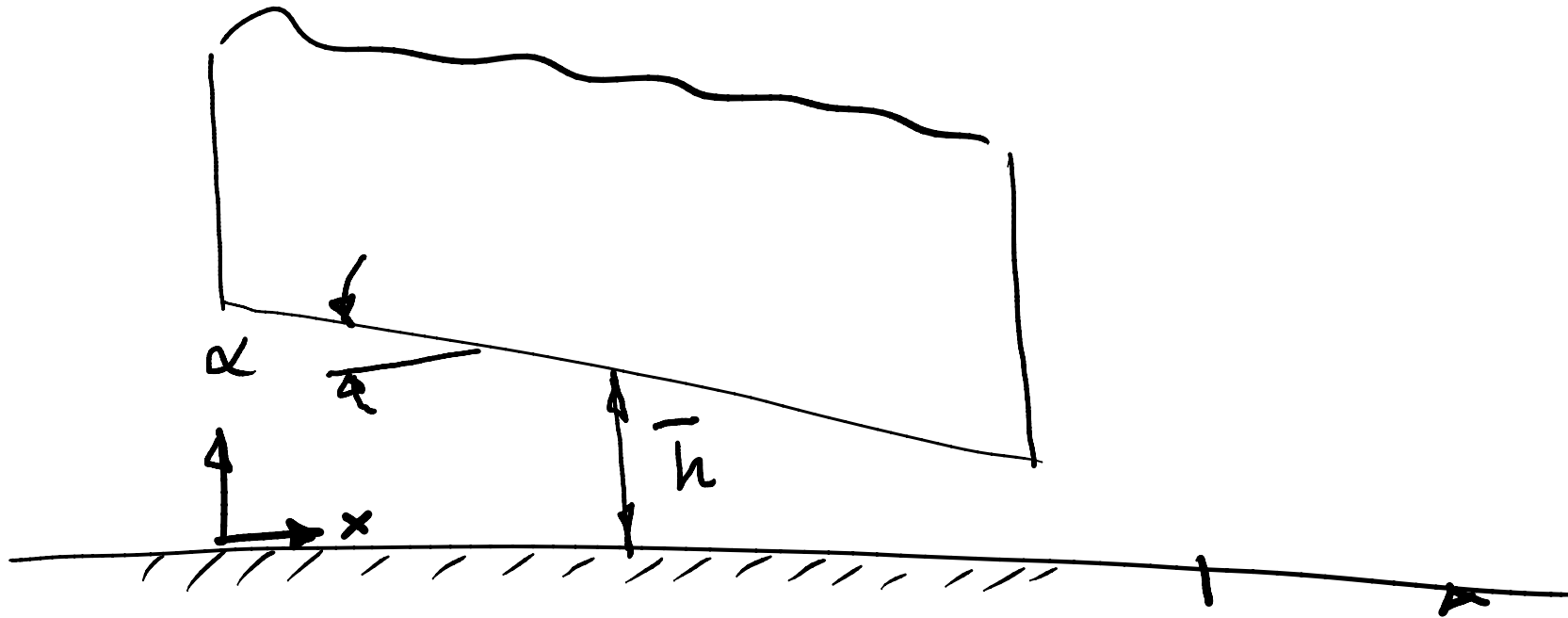
Instationäre Schichtströmung (laminar)

$$\rho \frac{D\mu}{Dt} = -\frac{dP}{dx} + \eta \left(\frac{\partial^2 \mu}{\partial y^2} + \frac{\partial^2 \mu}{\partial x^2} \right) \quad \text{x-Komp. der Turbulenzlsg.}$$

$$\rho \frac{\partial \mu}{\partial t} + \underbrace{\rho u \frac{\partial \mu}{\partial x}}_{\equiv 0} + \underbrace{\rho v \frac{\partial \mu}{\partial y}}_{\equiv 0} = \dots$$

$$\rho \frac{\partial \mu}{\partial t} = -\frac{dP}{dx} + \frac{d}{dy} \left(\tau \right)$$





\bar{h} mittlere Schichtdicke

U

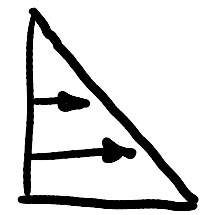
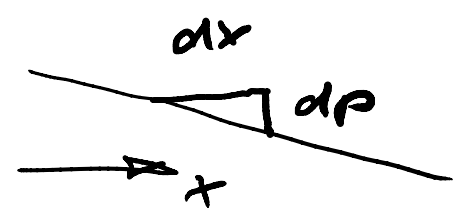
Wenn $\alpha \frac{U \bar{h}}{\nu} = \alpha Re \ll 1$, dann können die
Konvektive Terme in der Bewegungsgleichung vernachlässigt werden.

An jedem Ort x stellt sich eine
Durchschlepp-Schicht ein.

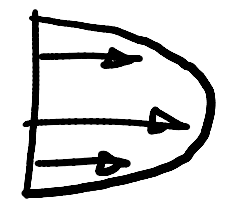


$$\frac{\mu}{\mu} = \frac{\gamma}{h(x)} - \left(\frac{dp}{dx} \right) \frac{h^2(x)}{2\mu} \left(1 - \frac{\gamma}{h(x)} \right) \frac{\gamma}{h(x)} = \dot{\gamma}_x$$

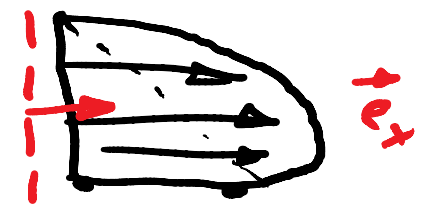
Schleppstr.



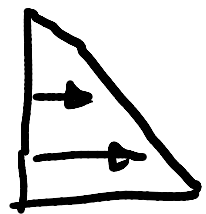
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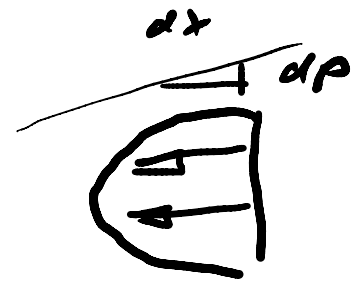
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$$\frac{dp}{dx} > 0$$



+



=



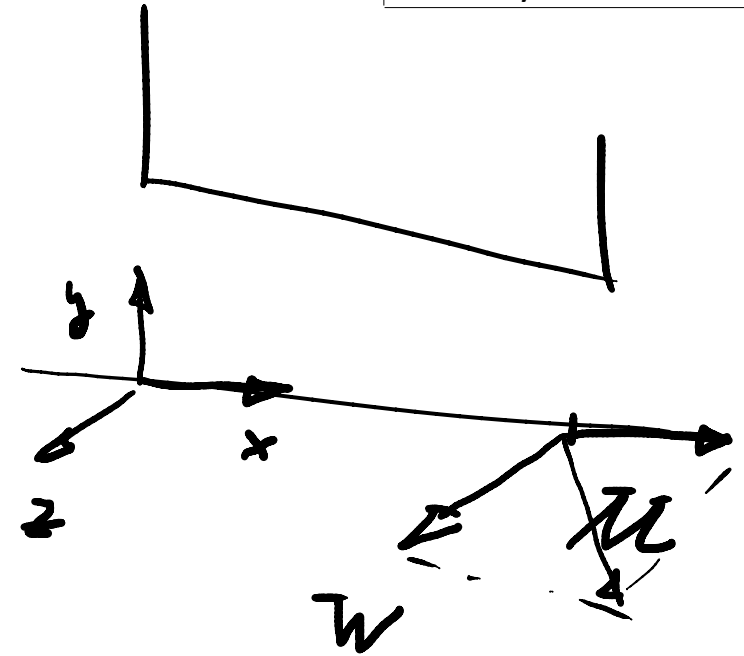


$$\dot{V}_x := \int_0^{h(x)} u \, dy = \frac{1}{2} \mu h - \frac{\partial p}{\partial x} \frac{h^3(x)}{12\eta}$$

$$\dot{V}_z = \frac{1}{2} W h - \frac{\partial p}{\partial z} \frac{h^3(x)}{12\eta}$$

$$\vec{V} = \dot{V}_x \vec{e}_x + \dot{V}_z \vec{e}_z$$

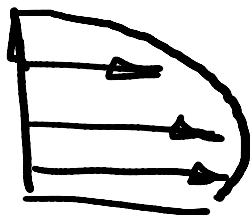
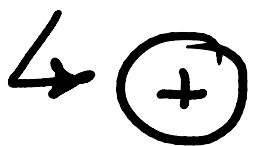
$$\frac{\partial \dot{V}_x}{\partial x} + \frac{\partial \dot{V}_z}{\partial z} \stackrel{!}{=} 0 \quad \text{aus Kontinuitätsgesetz.}$$





$$\frac{\partial}{\partial x} \left(\underbrace{\frac{h^3}{12}}_{\text{D} \rightarrow x} \frac{\partial p}{\partial x} \right) + \frac{\partial}{\partial z} \left(\underbrace{\frac{h^3}{12}}_{\text{D} \rightarrow z} \frac{\partial p}{\partial z} \right) = 6 \left(\underbrace{\frac{\partial(h\mu)}{\partial x}}_{\text{D} \rightarrow x} + \underbrace{\frac{\partial(h\nu)}{\partial z}}_{\text{D} \rightarrow z} + \underbrace{2 \frac{\partial h}{\partial t}}_{\text{D} \rightarrow t} \right) \quad (*)$$

Reynoldsche Gleichung für die
hydrodynamische Schmierung. } Poissongleichung für
 $P(x, z, t)$



Die Informationen über
das Geschwindigkeitsprofil
ist in (x) enthalten.

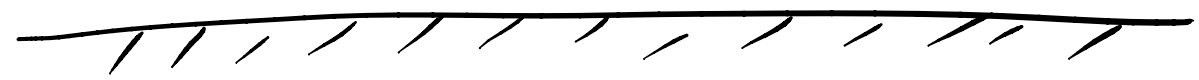


3D \rightarrow 2D Problem.



$$h = \bar{h} + \hat{h} \sin(\Omega t)$$

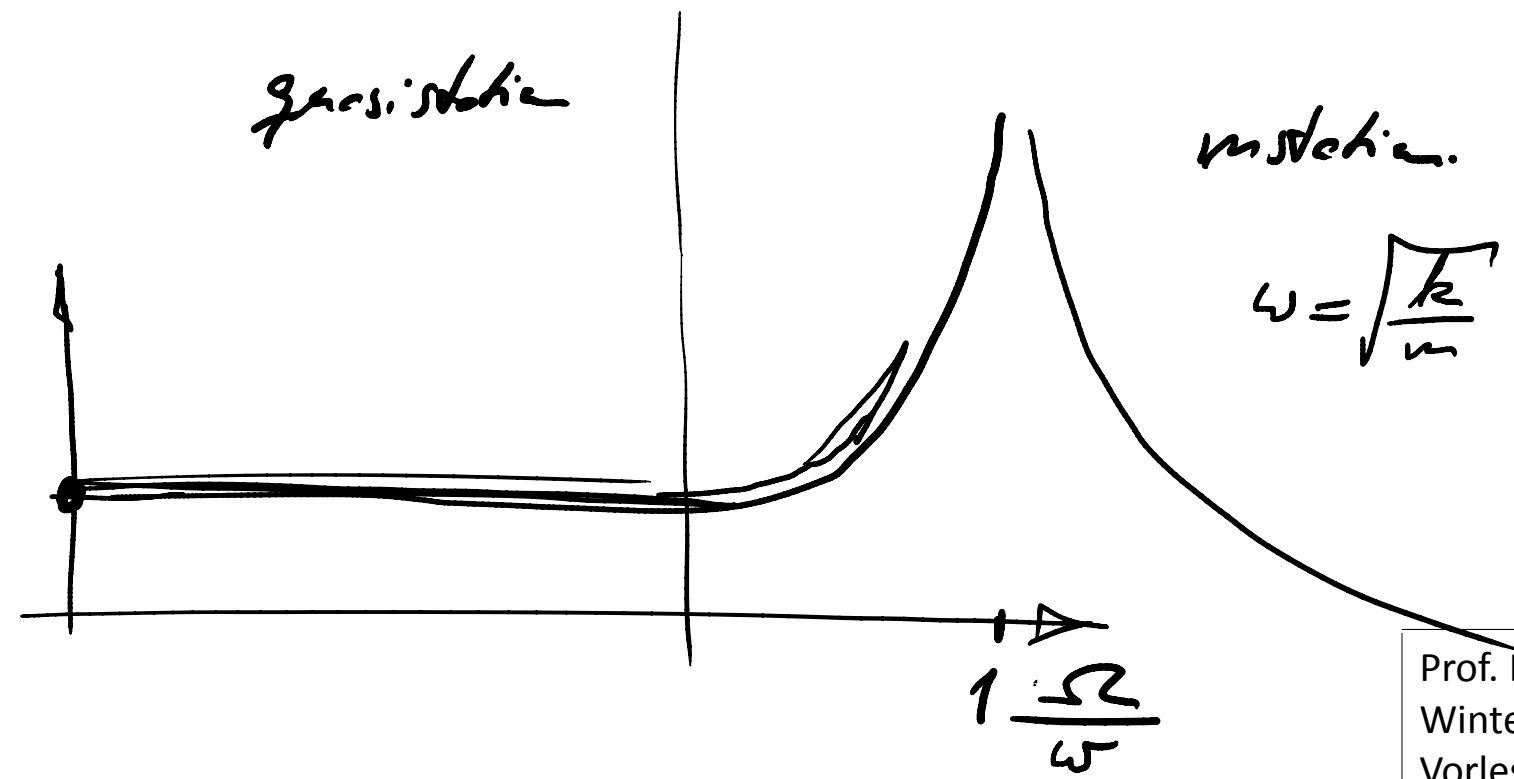
$$\frac{\Omega \hat{h}}{2} \ll 1$$



stationär

quasi-stationär

instationär

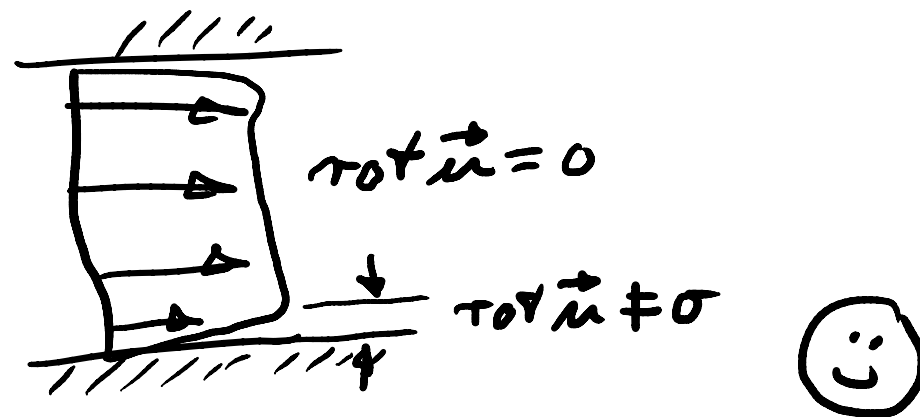




$\frac{\Omega h^2}{\nu} \gg 1 \rightsquigarrow$ Reibung hat
keine Zeit sich im
Spalt zu entwickeln.

\rightsquigarrow Blodprofil

▷ Φ ist bekannt.



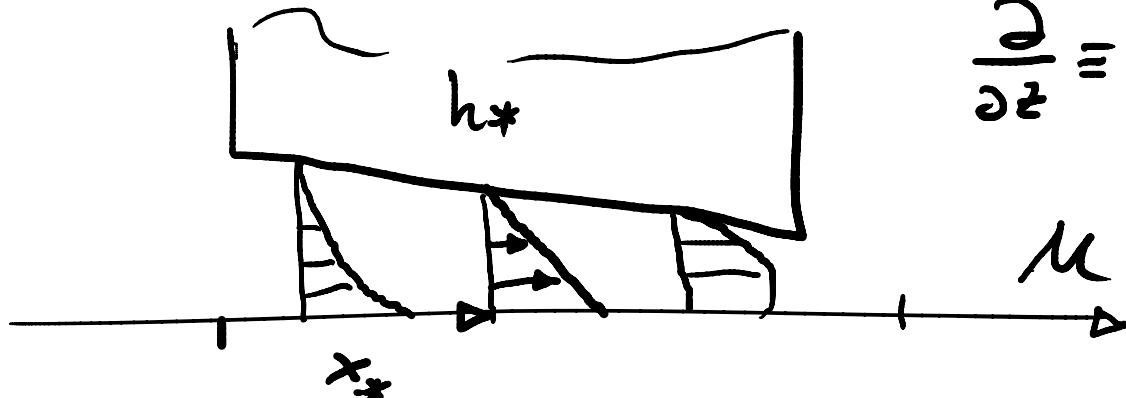
$$q' = \frac{\partial \phi}{\partial t} + \frac{1}{2} \nabla \phi \cdot \nabla \phi + \frac{p}{\rho} + gz$$

Bernoulli (gleich für $\tau \cdot \vec{n} = 0$)

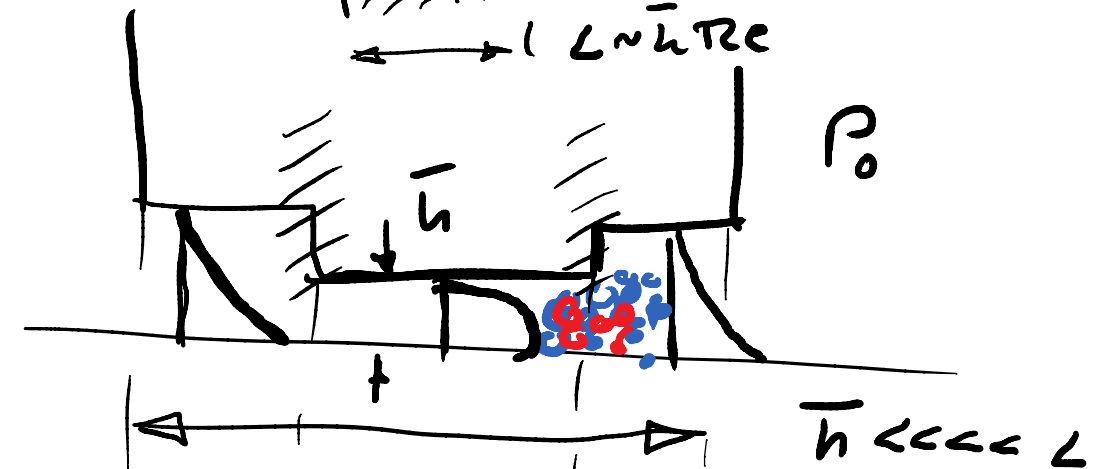
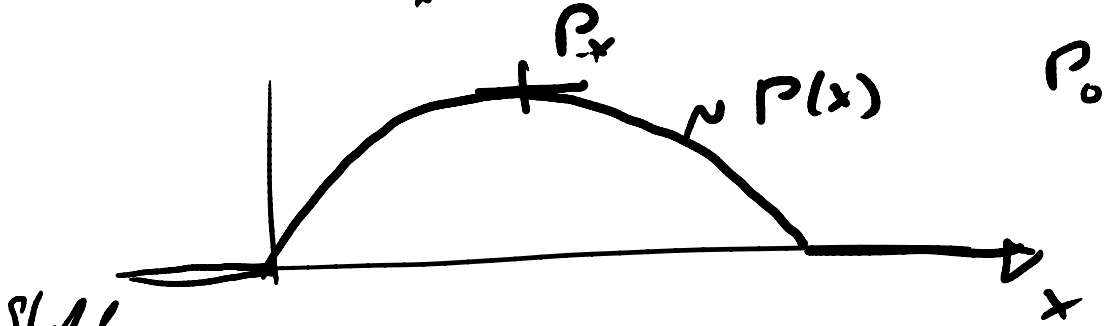
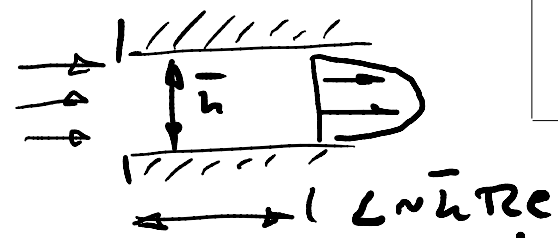
↳ $p(x, z, t)$.

$$\left. \begin{aligned} \vec{u} &= \nabla \Phi \\ \nabla \cdot \vec{u} &= 0 \end{aligned} \right\} \Delta \phi = 0.$$

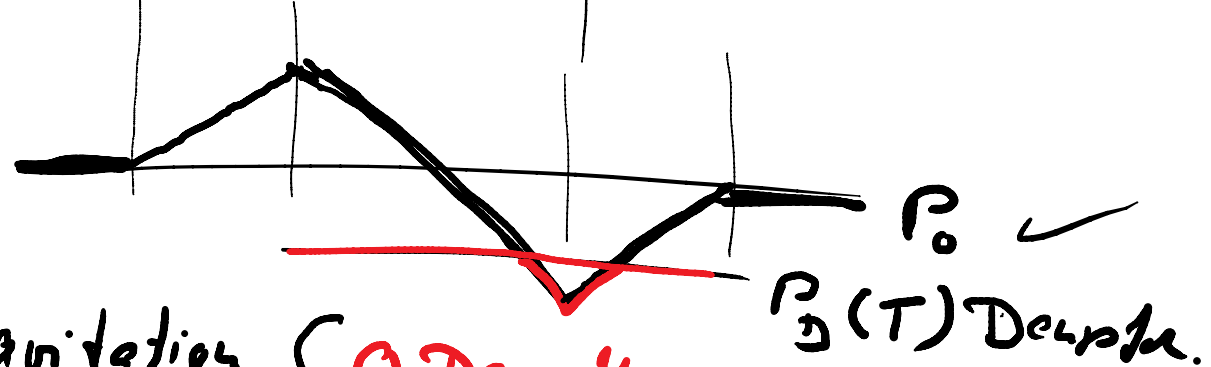
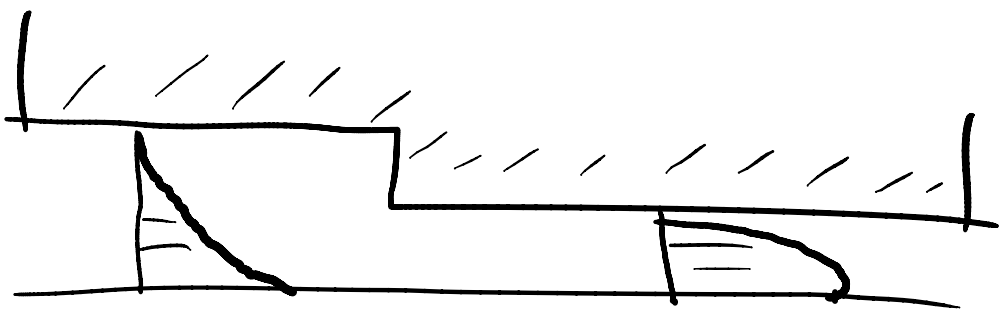
Keilström



$$\frac{\partial}{\partial z} \equiv 0, \quad w \equiv 0, \quad \frac{\partial h}{\partial t} \equiv 0$$



Stufenström



Kavitation

- Dampfbl. (red)
- Luftbl. (blue)
- Pseudo Kav. (black)



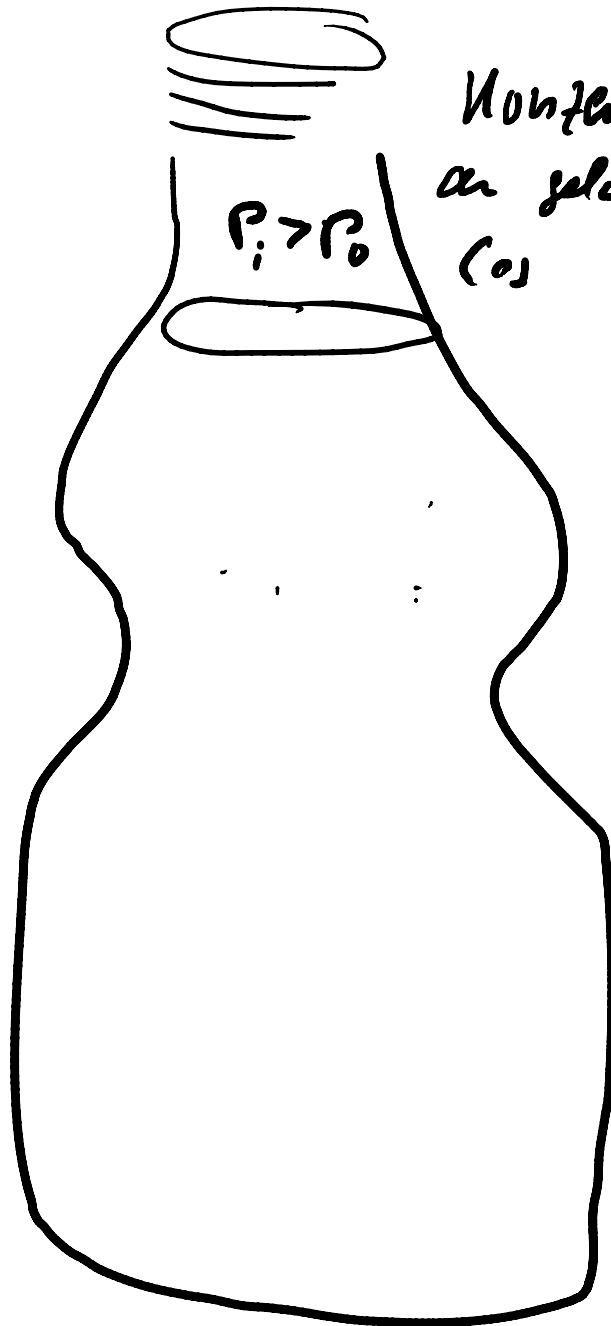
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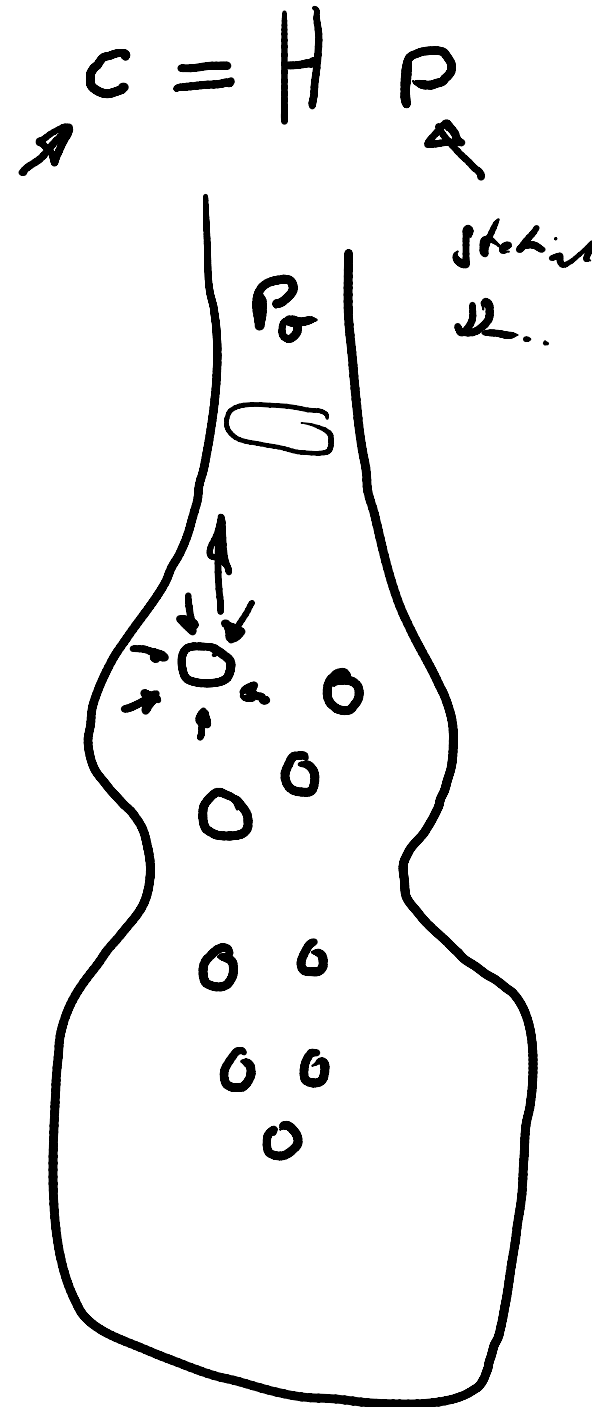
Technische
Fluidsysteme

Prof. Dr.-Ing. Peter Pelz
Wintersemester 2012/13
Vorlesung 7 F 92

Losungsfeld



Kontraktion
an gelöste
(O₂)



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Kavitiation ⊖ Funktion wird gestört

⊖ laut

⊖ Notwendigkeit.

Kavitieren ⊕ Medizin

⊕ Reinigung z.B.



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