

Recent developments in the study of dense matter



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Status report on project B05:
Nuclear matter equation of state for astrophysical applications

Marc Leonhardt

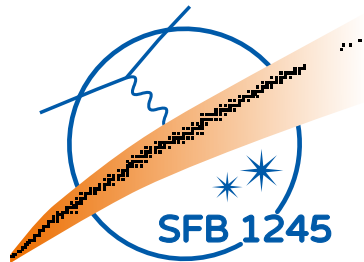
Institut für Kernphysik, Technische Universität Darmstadt

with

Martin Pospiech and Jens Braun,

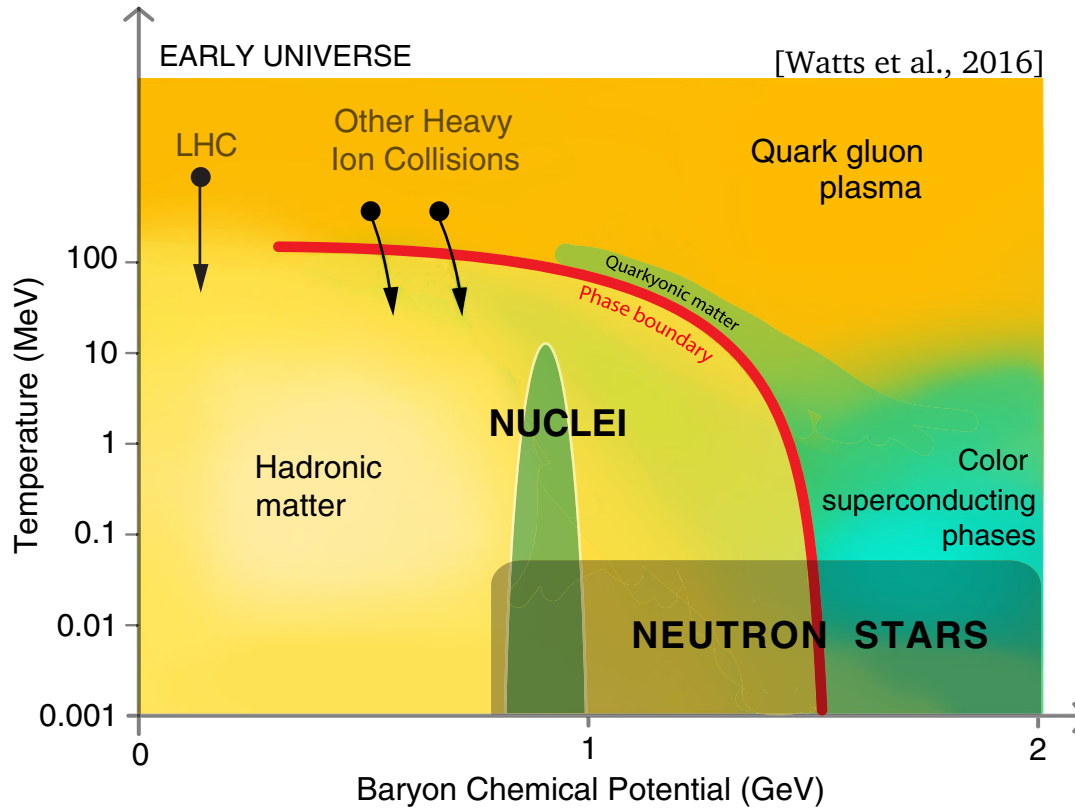
Christian Drischler, Corbinian Wellenhofer and Kai Hebeler

2nd CRC 1245 Workshop
Schloss Waldthausen, 2017

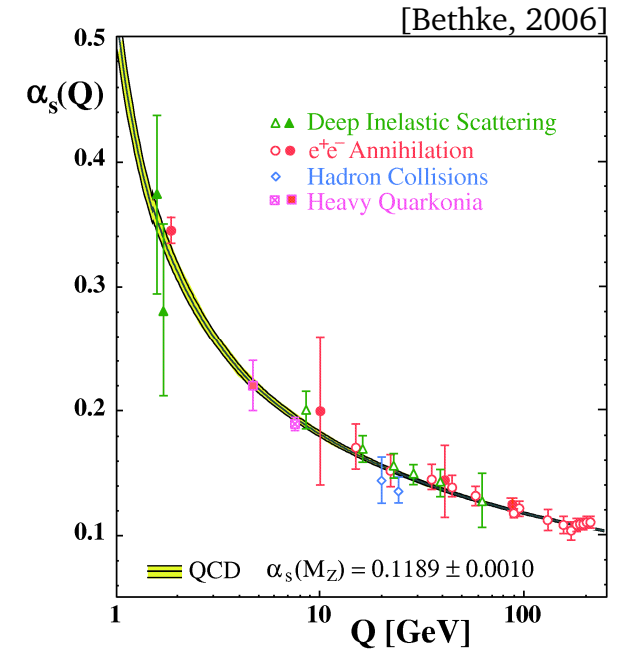
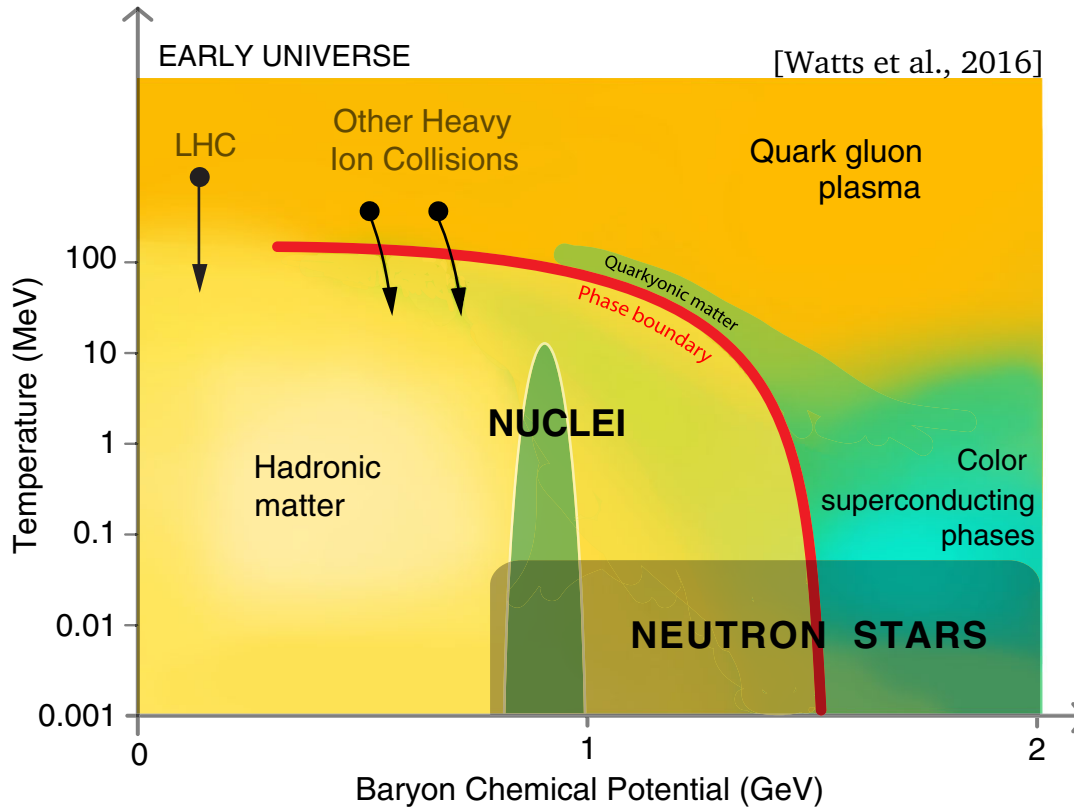


HIC | **FAIR**
for
Helmholtz International Center

QCD phase diagram: Neutron stars and the cold dense EoS



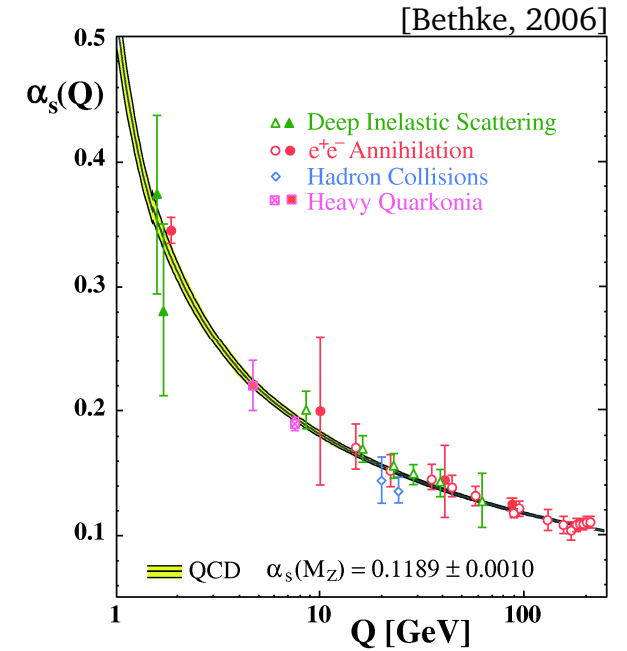
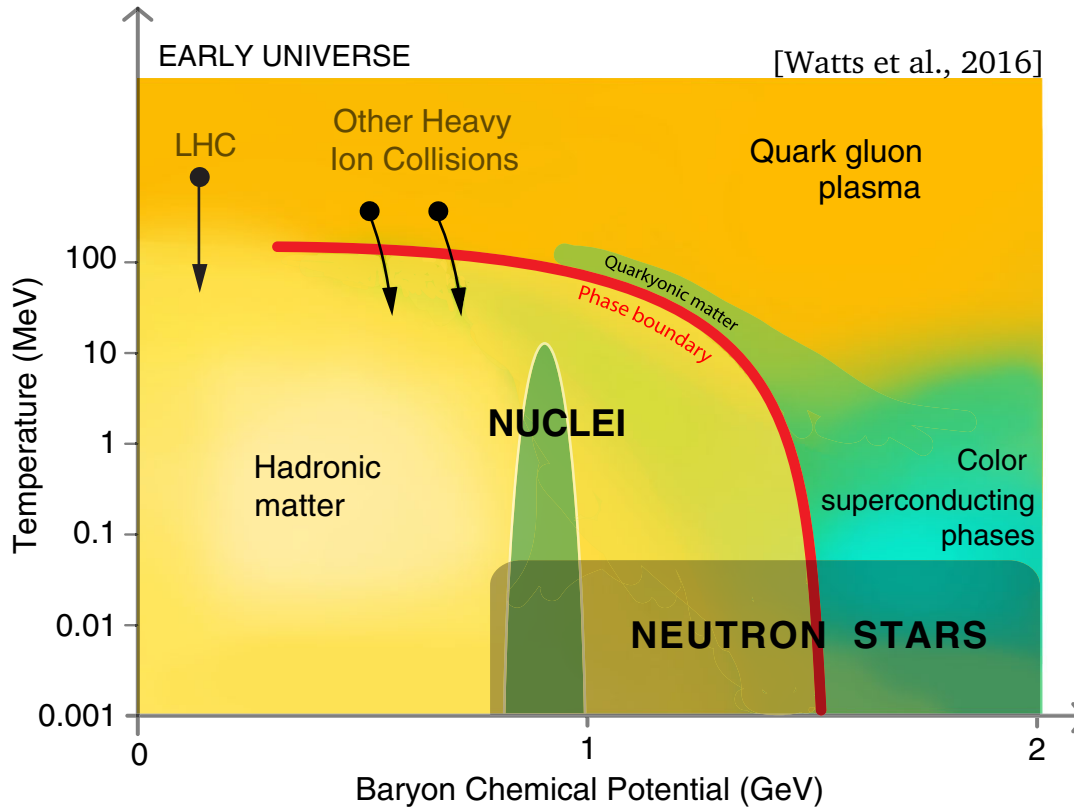
QCD phase diagram: Neutron stars and the cold dense EoS



$\lesssim 60 \rho_{\text{sat}}$ [Kurkela et al., 2014]

⌋ Perturbative methods

QCD phase diagram: Neutron stars and the cold dense EoS



$\lesssim 60 \rho_{\text{sat}}$ [Kurkela et al., 2014]

⌋ Perturbative methods

← Chiral effective field theory

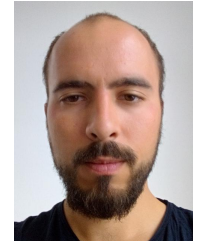
← Functional methods

Outline

- **Chiral effective field theory (at lower densities)**
 - Efficient Monte-Carlo framework for MBPT
 - Nuclear thermodynamics from χ EFT interactions
- **Functional renormalization group (at higher densities)**
 - Fierz-complete four-quark interactions in hot and dense QCD (2 flavors)
 - Ground state properties and phases
- Conclusions and outlook



Christian
Drischler



Corbinian
Wellenhofer



Martin
Pospiech



Marc
Leonhardt

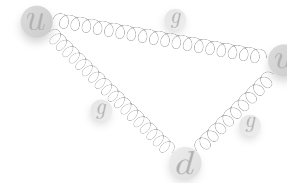
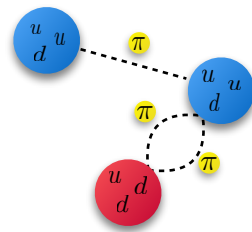
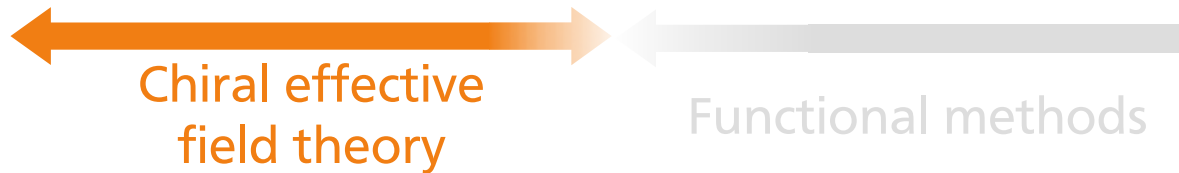
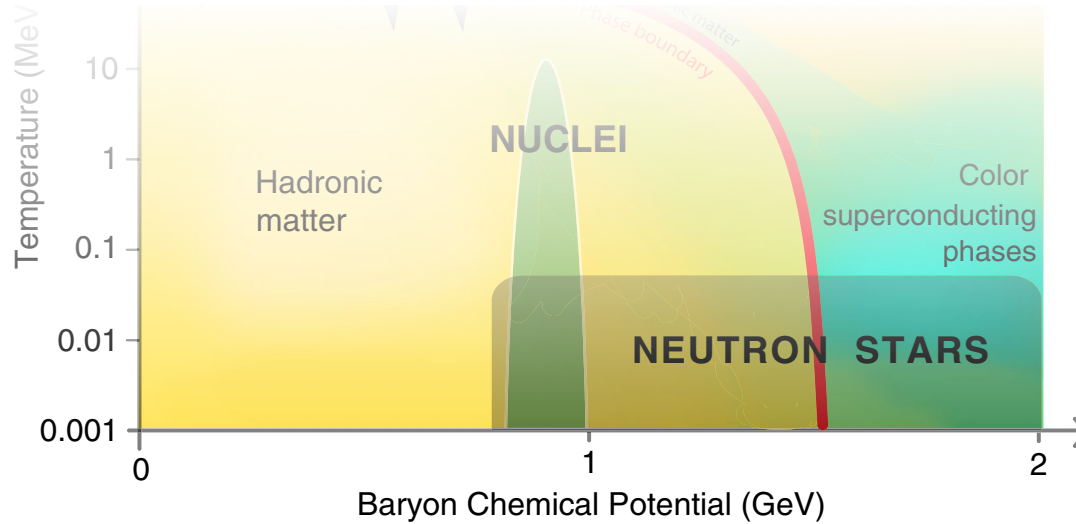
[Drischler, Hebeler, Schwenk, in preparation]

[Wellenhofer, Holt, Kaiser, Weise; '14, '15, '16]

[Braun, ML, Pospiech, arXiv:1705.00074]

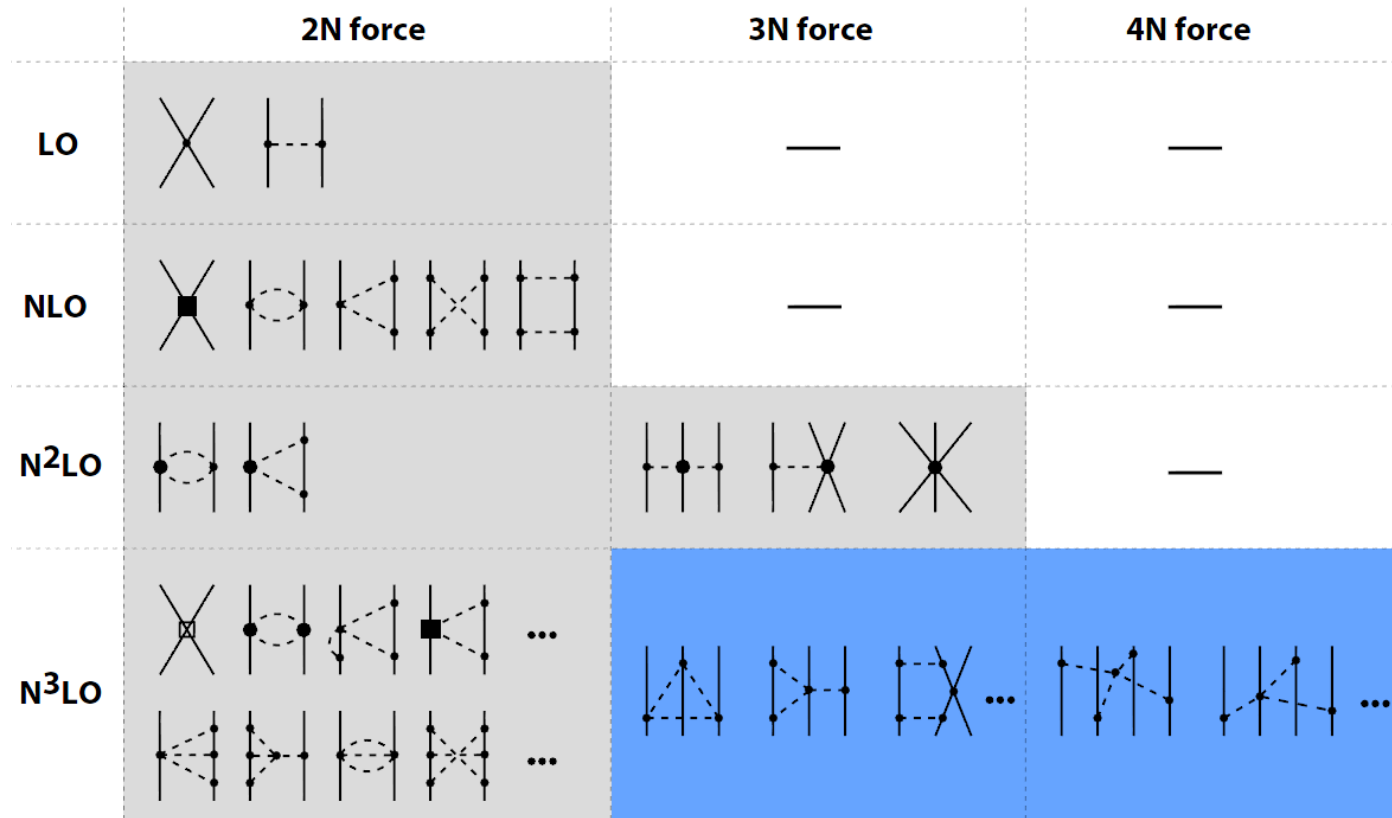
[Braun, ML, Pospiech, in preparation]

QCD phase diagram: Neutron stars and the cold dense EoS



Nuclear matter EOS for astrophysical applications

Chiral effective field theory

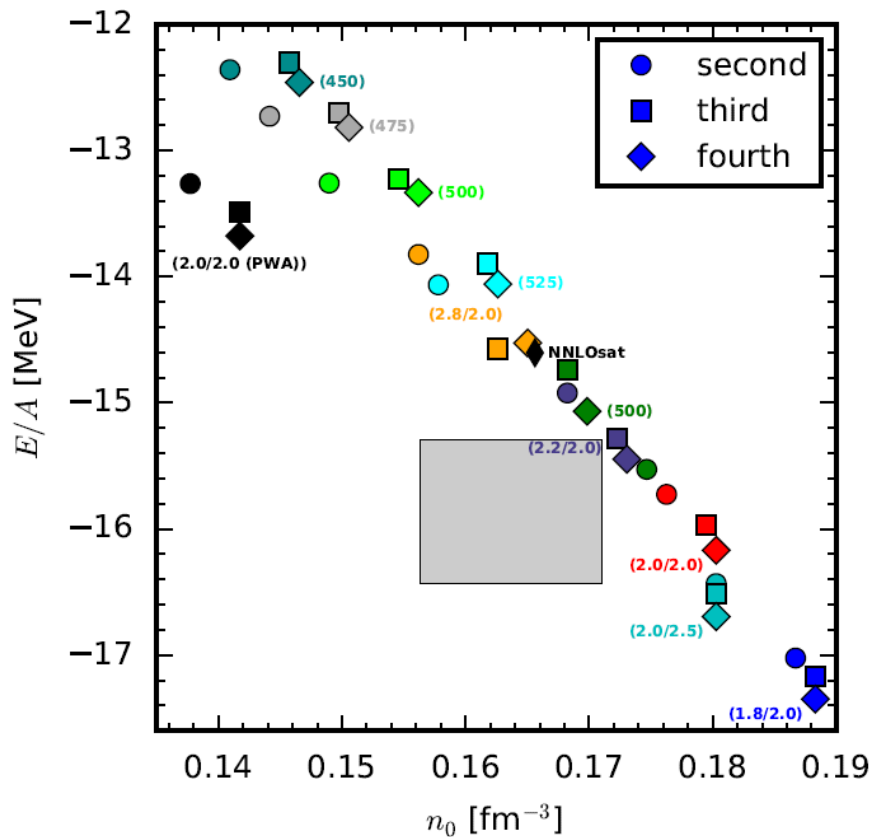


... and ongoing work at **N⁴LO**, **N⁵LO**, ...

Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Kaiser, Machleidt, Meißner, ...

Nuclear matter EOS for astrophysical applications

Efficient Monte-Carlo framework for MBPT

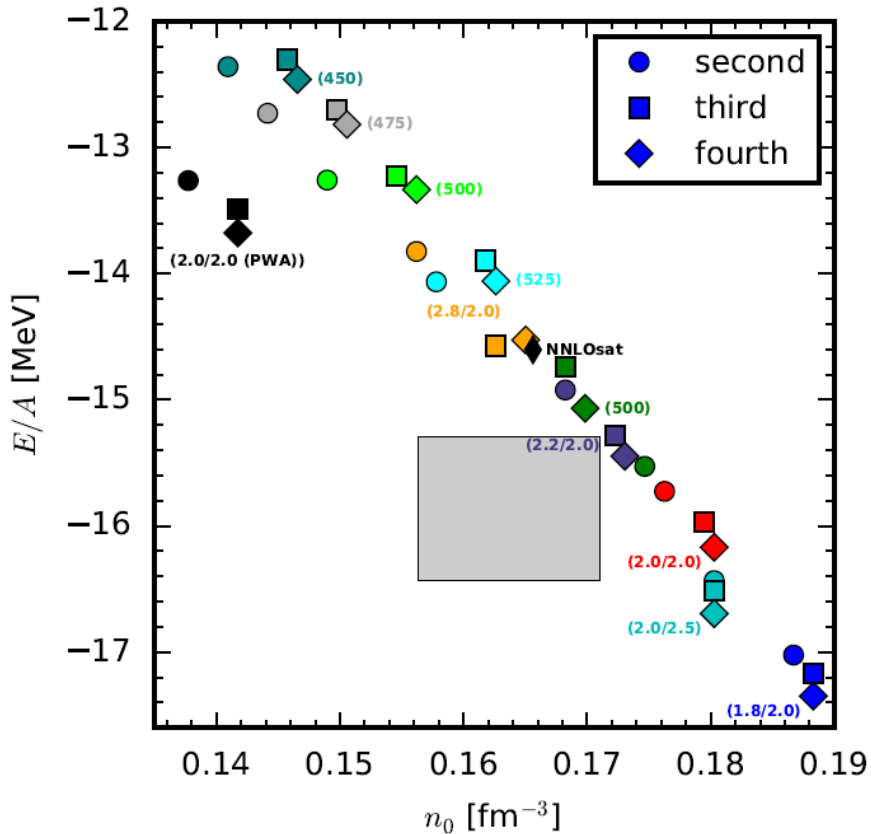


- based on **analytical expressions**
 - NN, 3N, 4N forces @ $N^3\text{LO}$ (no PW's)
 - MBPT for up to **4th** order (*automatic* code generation)

[Drischler, Hebeler, Schwenk, in preparation]

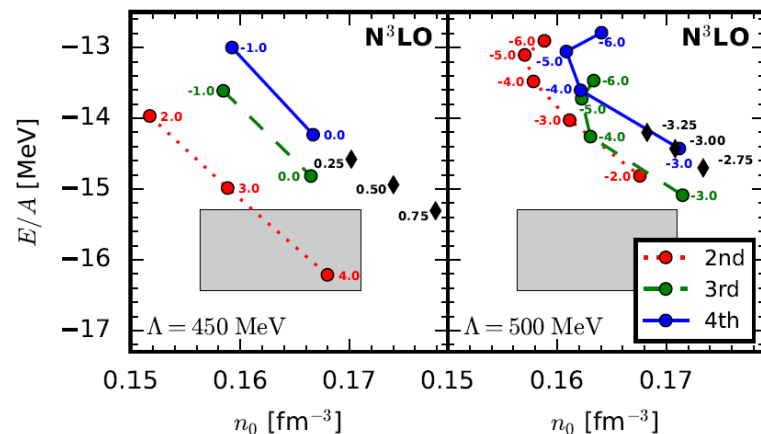
Nuclear matter EOS for astrophysical applications

Efficient Monte-Carlo framework for MBPT



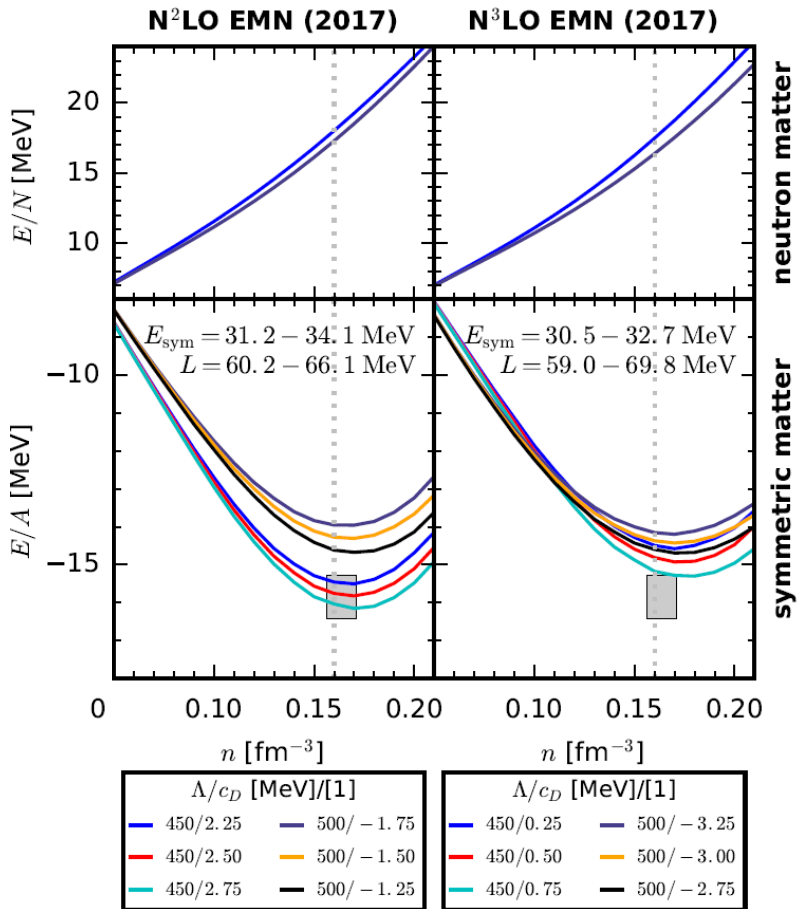
[Drischler, Hebeler, Schwenk, in preparation]

- based on **analytical expressions**
- NN, 3N, 4N forces @ N³LO (no PW's)
- MBPT for up to 4th order (*automatic* code generation)
- **aim:** guiding fits of next-generation interactions in terms of **saturation**
- fit 3N LECs c_D/c_E @ N³LO to ³H and study saturation

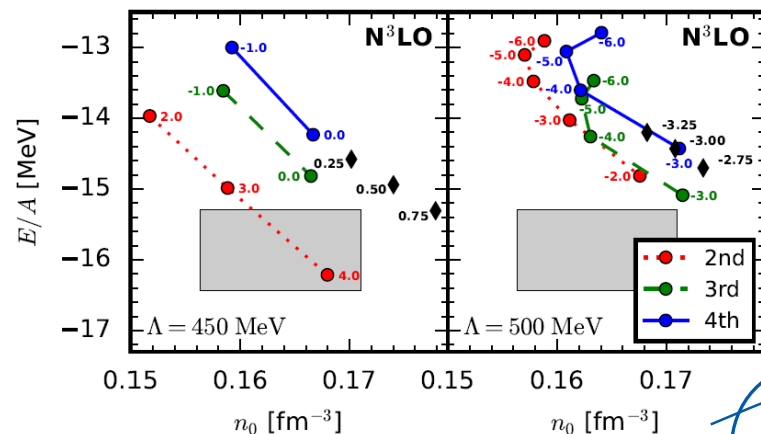


Nuclear matter EOS for astrophysical applications

Efficient Monte-Carlo framework for MBPT



- based on **analytical expressions**
- NN, 3N, 4N forces @ N³LO (no PW's)
- MBPT for up to **4th** order (*automatic* code generation)
- **aim:** guiding fits of next-generation interactions in terms of **saturation**
- fit 3N LECs c_D/c_E @ N³LO to ³H and study saturation



[Drischler, Hebeler, Schwenk, in preparation]

Nuclear thermodynamics

from chiral effective field theory interactions

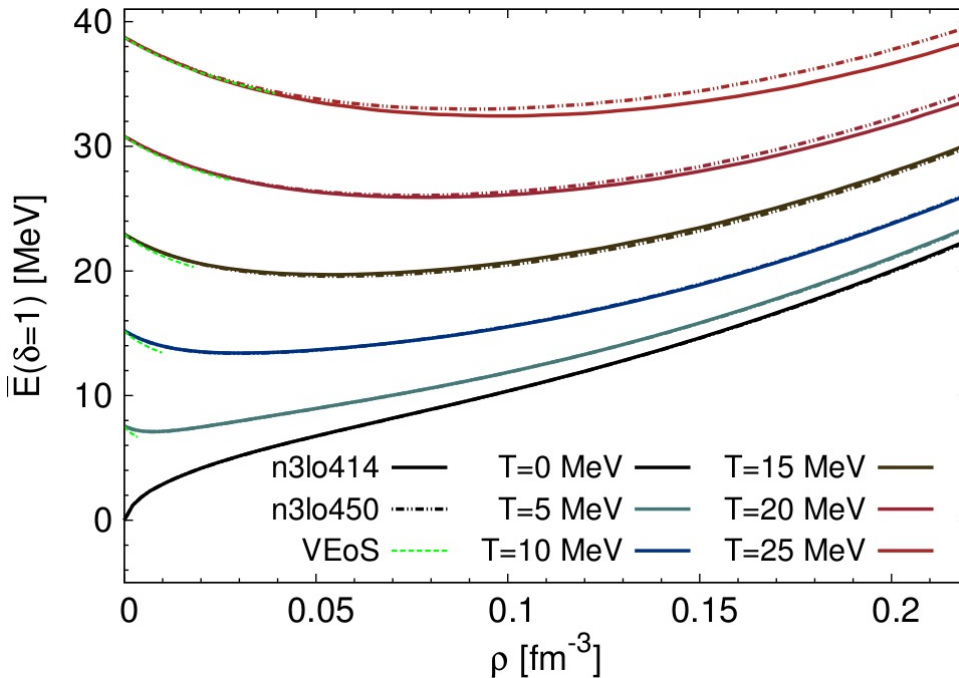


Figure: Internal energy of pure neutron matter ($\delta=1$; VEOs: virial expansion)

- ◆ Compute **thermodynamic properties of nuclear matter**
 - Needed for neutron star and supernova simulations
 - Large parameter space: **temperature T** , nucleon density ρ , **isospin asymmetry $\delta=1-2Y$** (where Y is the proton fraction)
- ◆ **Good benchmark results**
e.g.: good agreement with **virial expansion** at low densities (see Figure)
- ◆ **Future work:**
single-particle properties, improved calculations (better uncertainty estimates), ...

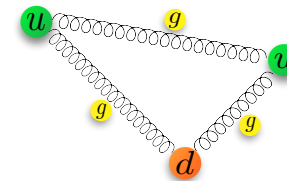
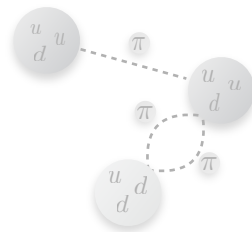
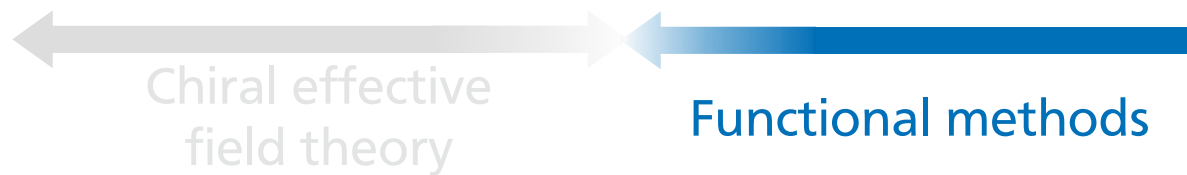
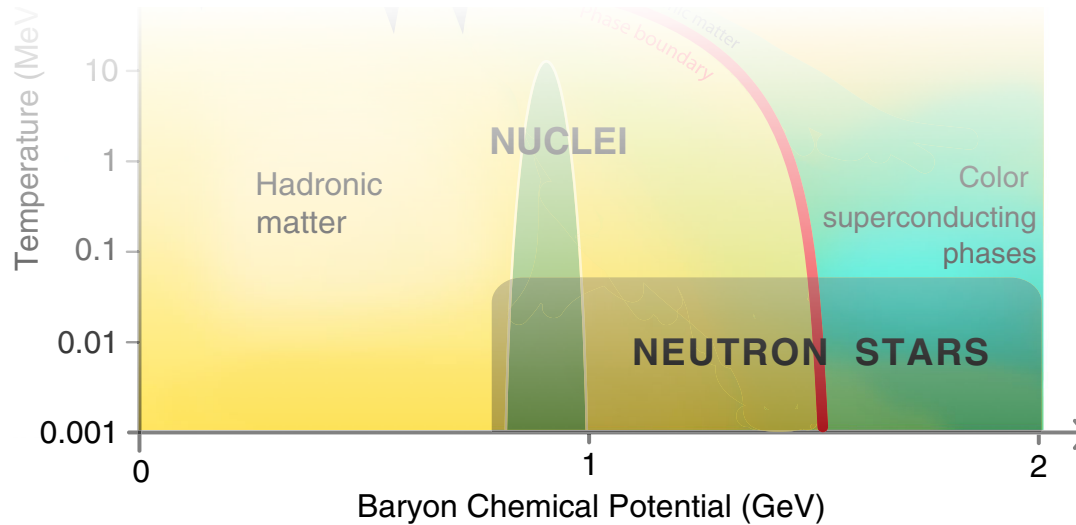
[Wellenhofer, Holt, Kaiser, Weise, PRC **89**, 064009 (2014)]

[Wellenhofer, Holt, Kaiser, PRC **92**, 015801 (2015)]

[Wellenhofer, Holt, Kaiser, PRC **93**, 055802 (2016)]

[Schwenk, Horowitz, Phys. Lett. **B638**, 153-159 (2006)]

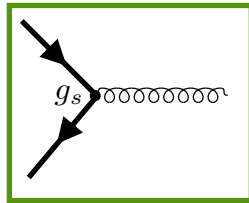
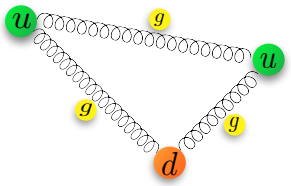
QCD phase diagram: Neutron stars and the cold dense EoS



Functional renormalization group (FRG)

From high to low energies in QCD

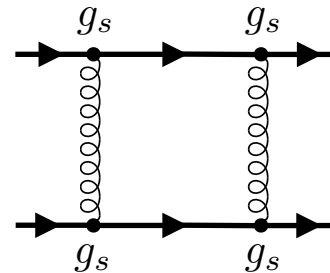
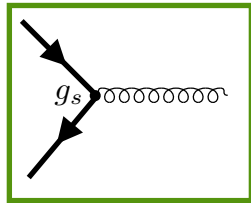
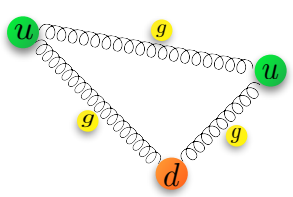
RG steps



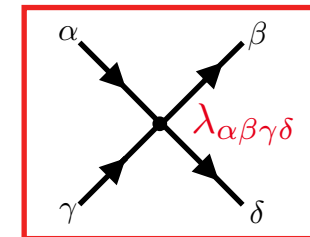
Functional renormalization group (FRG)

From high to low energies in QCD

RG steps



Four-quark
interactions



1. Important effective interactions
2. Encode information on the ground state and spontaneous symmetry breaking
3. Full treatment possible and crucial

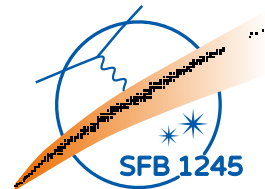
Four-quark interactions and symmetry breaking

Aspects of the low-energy sector in QCD



TECHNISCHE
UNIVERSITÄT
DARMSTADT

$$S[\bar{\psi}, \psi] = \int_x \left\{ \bar{\psi} i \not{\partial} \psi + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} [(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5 \vec{\tau}\psi)^2] \right\}$$



Four-quark interactions and symmetry breaking

Aspects of the low-energy sector in QCD



TECHNISCHE
UNIVERSITÄT
DARMSTADT

$$S[\bar{\psi}, \psi] = \int_x \left\{ \bar{\psi} i \not{\partial} \psi + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} [(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5 \vec{\tau}\psi)^2] \right\}$$

Partial bosonization

$$\begin{aligned} \sigma &\sim \bar{\psi}\psi \\ \vec{\pi} &\sim \bar{\psi}\gamma_5 \vec{\tau}\psi \end{aligned}$$

Symmetry of the
ground state

$$U_B \sim \frac{1}{\bar{\lambda}_{(\sigma-\pi)}} (\sigma^2 + \pi^2) + \dots$$



Four-quark interactions and symmetry breaking

Aspects of the low-energy sector in QCD

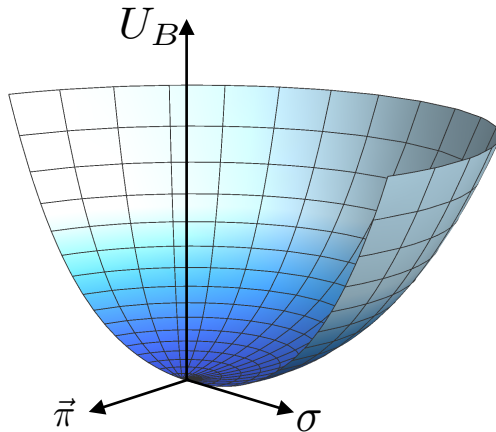
$$S[\bar{\psi}, \psi] = \int_x \left\{ \bar{\psi} i \not{\partial} \psi + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} [(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5 \vec{\tau}\psi)^2] \right\}$$

Partial bosonization

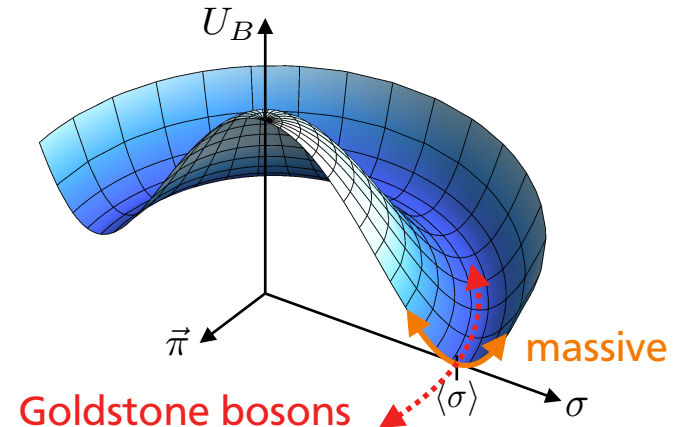
$$\begin{aligned} \sigma &\sim \bar{\psi}\psi \\ \vec{\pi} &\sim \bar{\psi}\gamma_5 \vec{\tau}\psi \end{aligned}$$

Symmetry of the
ground state

$$U_B \sim \frac{1}{\bar{\lambda}_{(\sigma-\pi)}} (\sigma^2 + \pi^2) + \dots$$



$$\frac{1}{\bar{\lambda}_{(\sigma-\pi)} \Big|_{k_0}} \longrightarrow 0$$



Four-quark interactions and symmetry breaking

Aspects of the low-energy sector in QCD

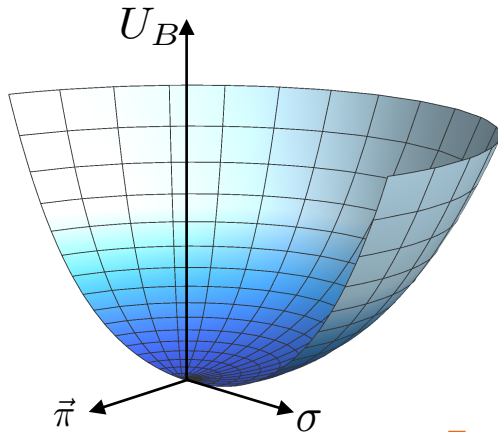
$$S[\bar{\psi}, \psi] = \int_x \left\{ \bar{\psi} i \not{\partial} \psi + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} [(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5 \vec{\tau}\psi)^2] \right\}$$

Partial bosonization

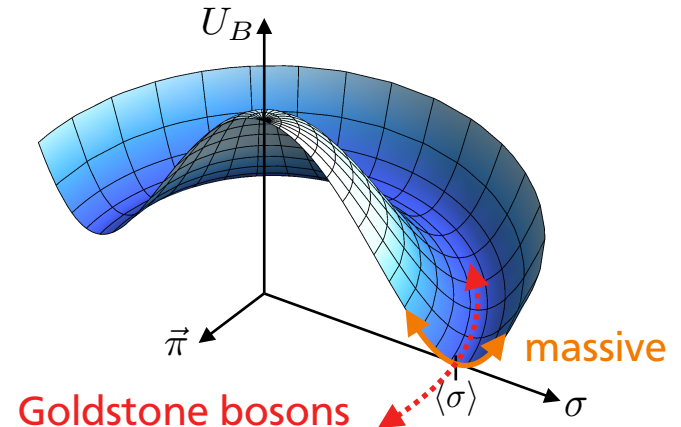
$$\begin{aligned} \sigma &\sim \bar{\psi}\psi \\ \vec{\pi} &\sim \bar{\psi}\gamma_5 \vec{\tau}\psi \end{aligned}$$

Symmetry of the
ground state

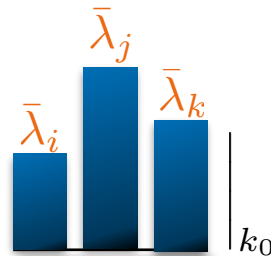
$$U_B \sim \frac{1}{\bar{\lambda}_{(\sigma-\pi)}} (\sigma^2 + \pi^2) + \dots$$



$$\frac{1}{\bar{\lambda}_{(\sigma-\pi)}} \Big|_{k_0} \longrightarrow 0$$



Assessing relative
interaction strengths

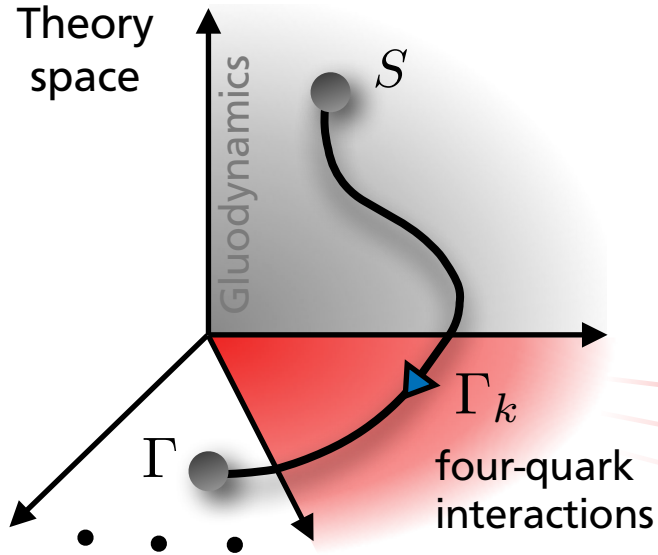


Formation of
specific condensates

Four-quark interactions and symmetries

Fierz-complete basis of interactions

Studies so far: Incomplete subsets,
typically 1-3 four-quark interaction channels.



Ambiguities!

Neglected contributions!



$$[(\bar{\psi}\gamma_{\mu}\psi)^2]$$

$$[(\bar{\psi}\gamma_{\mu}\gamma_5 T^a \psi)^2]$$

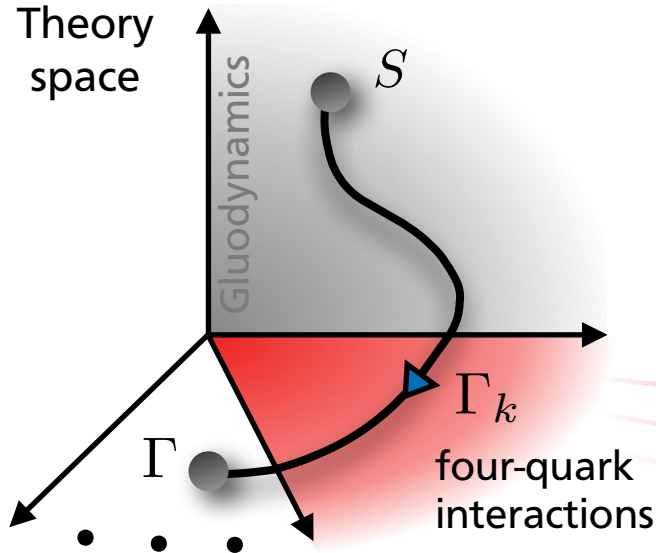
⋮

*Everything that is not
forbidden is allowed.*

Four-quark interactions and symmetries

Fierz-complete basis of interactions

Studies so far: Incomplete subsets,
typically 1-3 four-quark interaction channels.



Ambiguities!

Neglected contributions!



$$[(\bar{\psi}\gamma_{\mu}\psi)^2]$$

$$[(\bar{\psi}\gamma_{\mu}\gamma_5 T^a \psi)^2]$$

⋮

*Everything that is not
forbidden is allowed.*

Only constrained by symmetries!

T (time reversal) ✓ P (parity) ✓ ~~⊗~~ (charge conj.)

Lorentz group: $SO(1, 3) \longrightarrow SO(3)$

Flavor space: $SU_L(2) \otimes SU_R(2) \otimes U_V(1)$

Color space: $SU(N_c)$

Four-quark interactions and symmetries

Fierz-complete basis of interactions

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ \bar{\psi}(iZ_{\parallel}\gamma_0\partial_0 + iZ_{\perp}\gamma_i\partial_i - iZ_{\mu}\mu\gamma_0)\psi + \frac{1}{2} \sum_i \bar{\lambda}_i \mathcal{L}_{(\bar{\psi}\psi)^2}^i \right\}$$

In total **20 channels** meet
symmetry constraints

Fierz identities 

Fierz-complete basis:
10 channels

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{(\sigma-\pi)} = (\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2$$

↔ formation of
chiral condensate

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{\text{CSC}} \sim (i\bar{\psi}\gamma_5\tau_A t_c^{A'} \mathcal{C}\bar{\psi}^T)(i\psi^T \mathcal{C}\gamma_5\tau_A t_c^{A'} \psi) \quad J^P = 0^+$$

↔ formation of
diquark condensate

[Rapp, Schäfer, Shuryak, Velkovsky, 1998]

+ 8 more interaction channels

RG flow of four-quark interactions

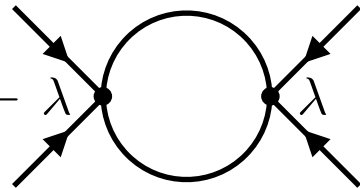
Qualitative behavior and the effect of external parameters



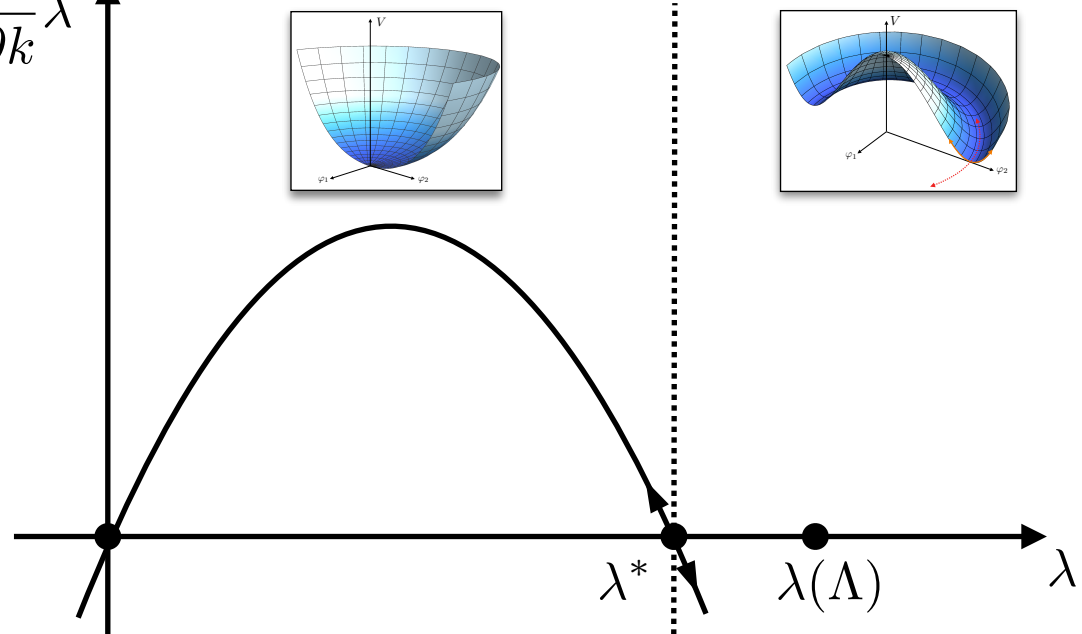
TECHNISCHE
UNIVERSITÄT
DARMSTADT

RG flow equation:

$$k \frac{\partial}{\partial k} \lambda = 2\lambda - \lambda$$



$$k \frac{\partial}{\partial k} \lambda$$



RG flow of four-quark interactions

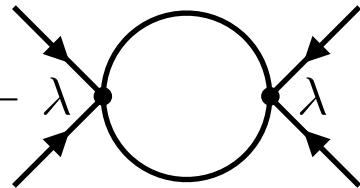
Qualitative behavior and the effect of external parameters



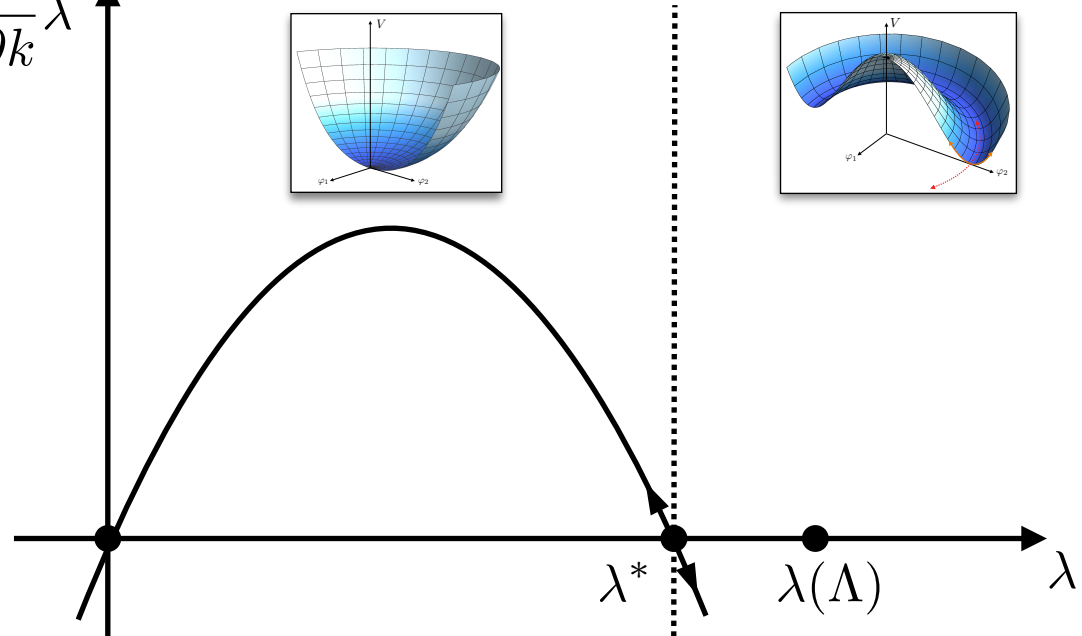
TECHNISCHE
UNIVERSITÄT
DARMSTADT

RG flow equation:

$$k \frac{\partial}{\partial k} \lambda = 2\lambda -$$



$$k \frac{\partial}{\partial k} \lambda$$



Scale fixing procedure:

- Only $\lambda_{(\sigma-\pi)}(\Lambda)$ assumes **finite** value as inspired by gluon-induced four-quark flows [Braun, 2011]
- UV value tuned so that specific **symmetry-breakdown scale** k_0 is obtained which corresponds to a quark mass of 300 MeV in the vacuum limit



RG flow of four-quark interactions

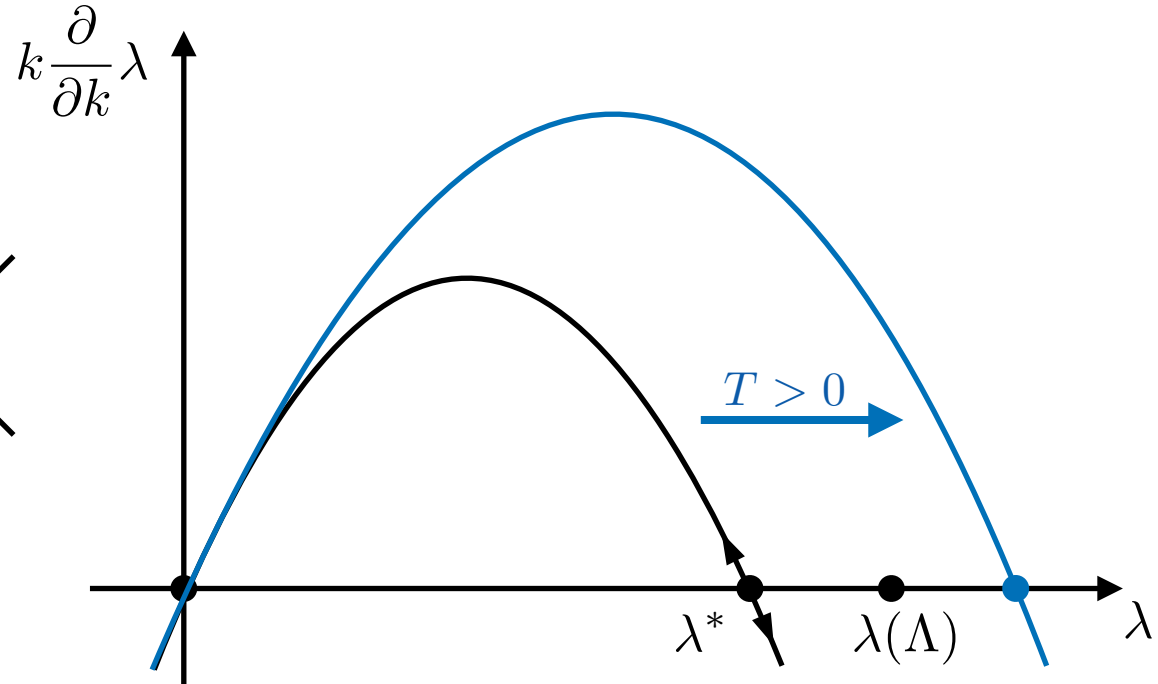
Qualitative behavior and the effect of external parameters



TECHNISCHE
UNIVERSITÄT
DARMSTADT

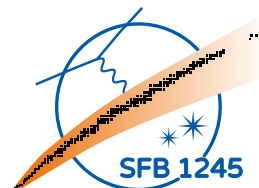
RG flow equation:

$$k \frac{\partial}{\partial k} \lambda = 2\lambda - \lambda$$



Scale fixing procedure:

- Only $\lambda_{(\sigma-\pi)}(\Lambda)$ assumes **finite** value as inspired by gluon-induced four-quark flows [Braun, 2011]
- UV value tuned so that specific **symmetry-breakdown scale** k_0 is obtained which corresponds to a quark mass of 300 MeV in the vacuum limit



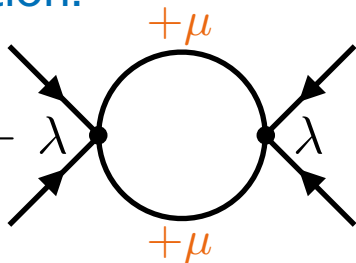
RG flow of four-quark interactions

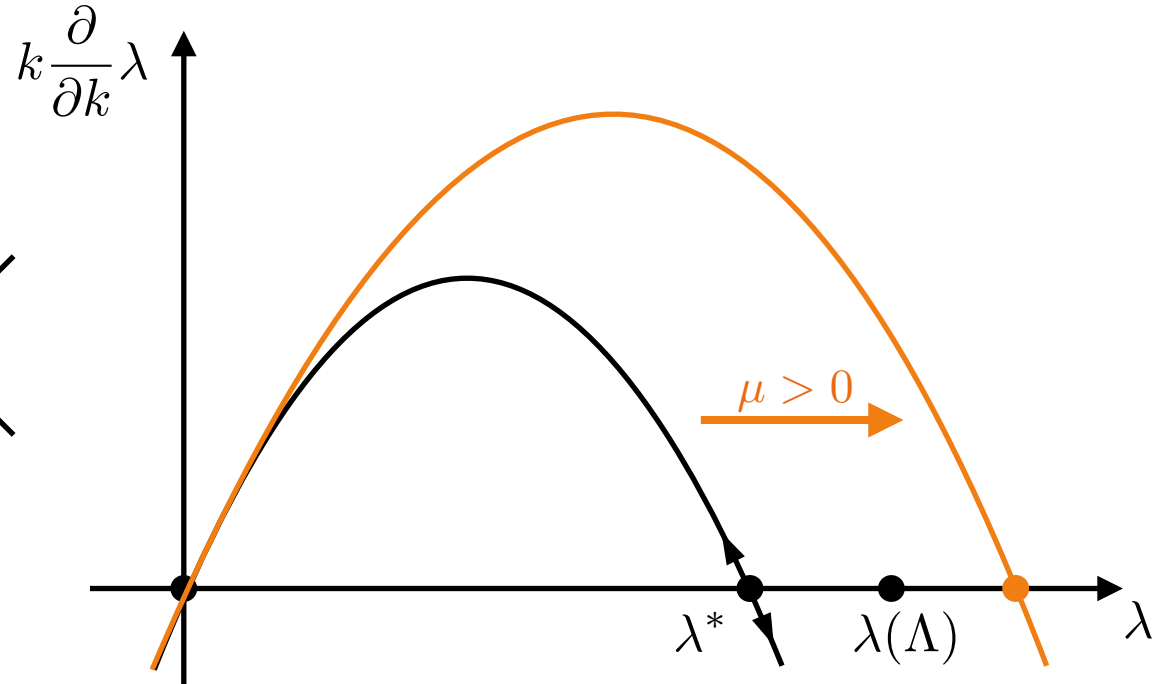
Qualitative behavior and the effect of external parameters



TECHNISCHE
UNIVERSITÄT
DARMSTADT

RG flow equation:

$$k \frac{\partial}{\partial k} \lambda = 2\lambda - \lambda$$




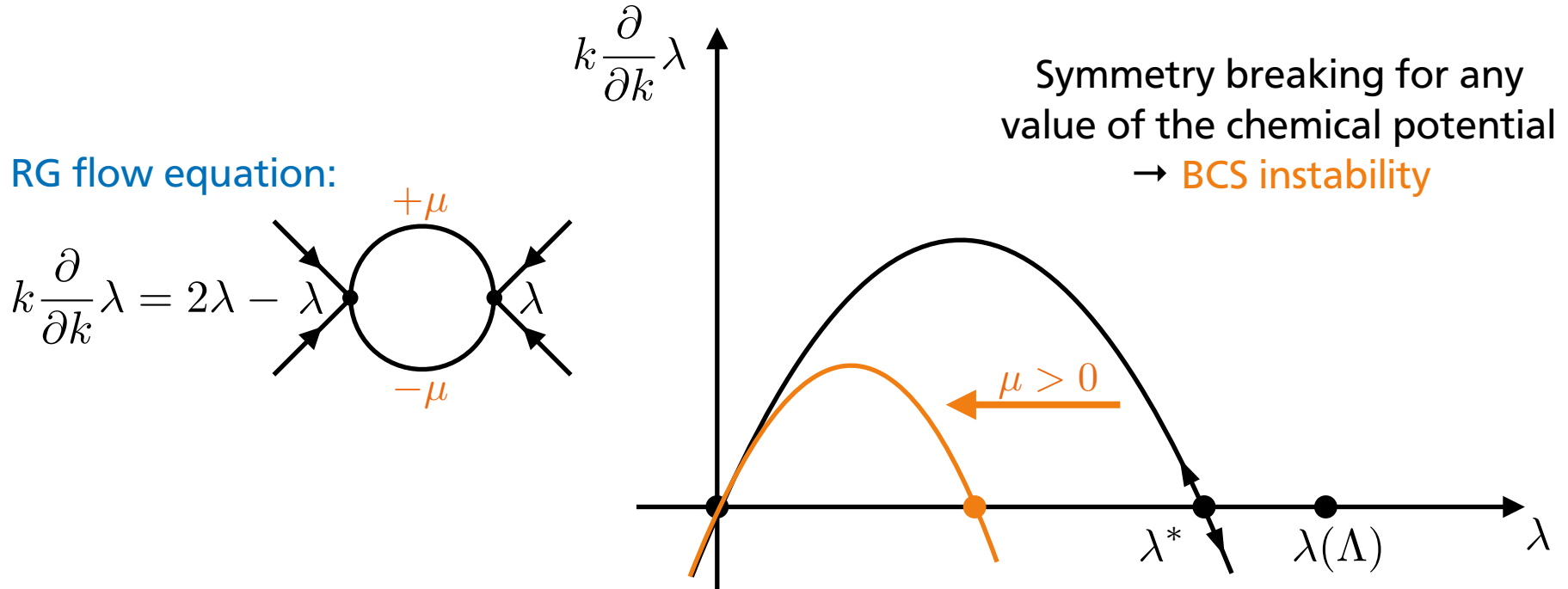
Scale fixing procedure:

- Only $\lambda_{(\sigma-\pi)}(\Lambda)$ assumes **finite** value as inspired by gluon-induced four-quark flows [Braun, 2011]
- UV value tuned so that specific **symmetry-breakdown scale** k_0 is obtained which corresponds to a quark mass of 300 MeV in the vacuum limit



RG flow of four-quark interactions

Qualitative behavior and the effect of external parameters



Scale fixing procedure:

- Only $\lambda_{(\sigma-\pi)}(\Lambda)$ assumes **finite** value as inspired by gluon-induced four-quark flows [Braun, 2011]
- UV value tuned so that specific **symmetry-breakdown scale** k_0 is obtained which corresponds to a quark mass of 300 MeV in the vacuum limit

RG flow of four-quark interactions

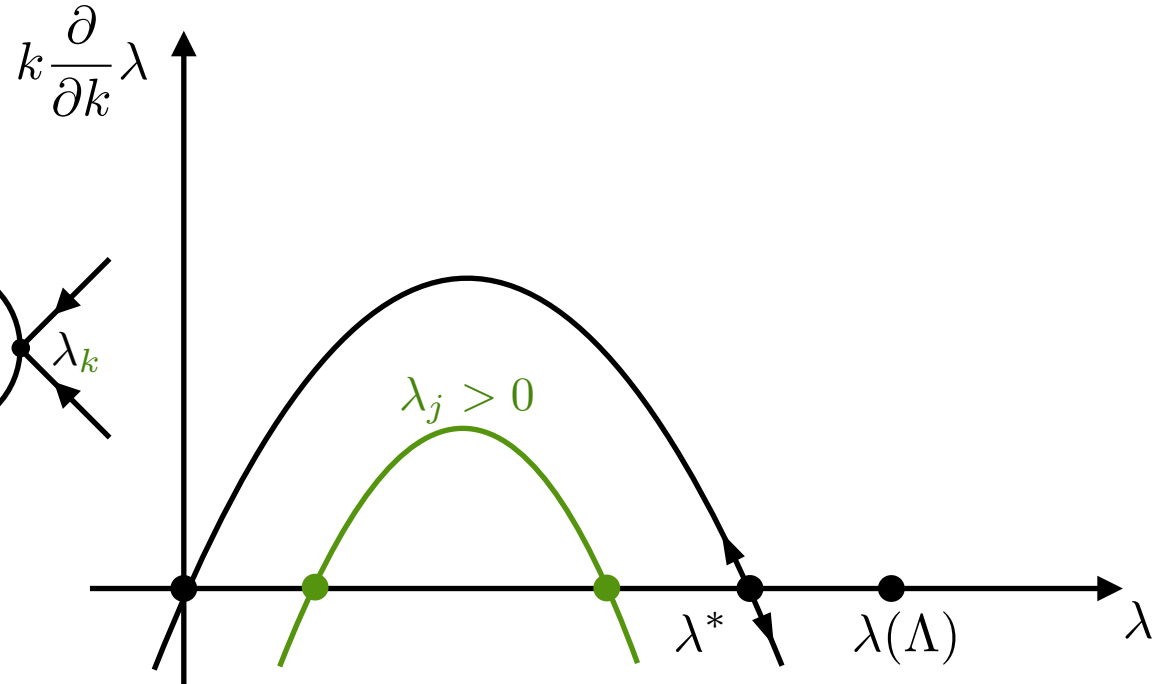
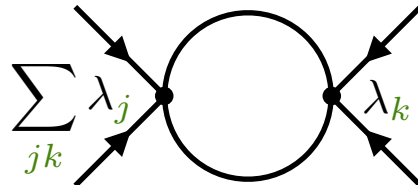
Qualitative behavior and the effect of external parameters



TECHNISCHE
UNIVERSITÄT
DARMSTADT

RG flow equation:

$$k \frac{\partial}{\partial k} \lambda_i = 2\lambda_i - \sum_{jk} \lambda_j$$



Scale fixing procedure:

- Only $\lambda_{(\sigma-\pi)}(\Lambda)$ assumes **finite** value as inspired by gluon-induced four-quark flows [Braun, 2011]
- UV value tuned so that specific **symmetry-breakdown scale** k_0 is obtained which corresponds to a quark mass of 300 MeV in the vacuum limit

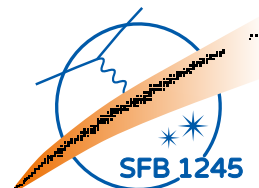
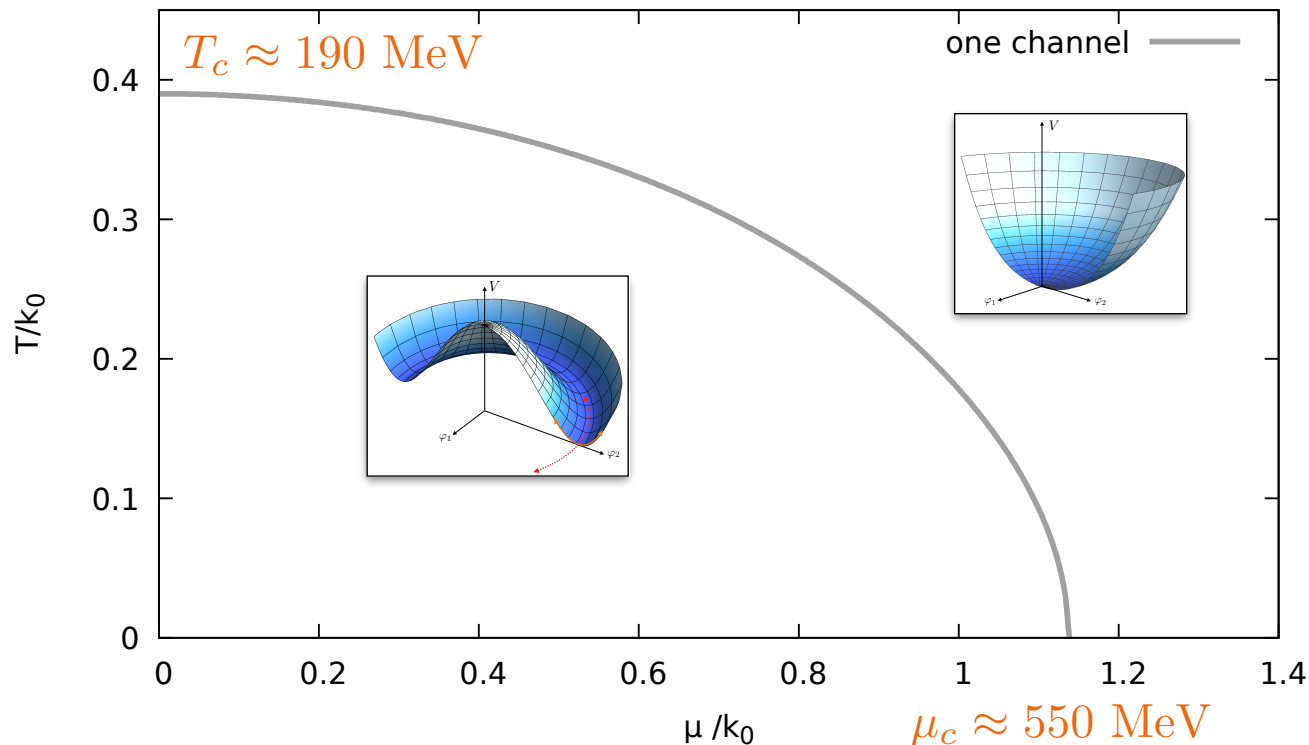
Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking



Ansatz one-channel approximation

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} [(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2] \right\}$$



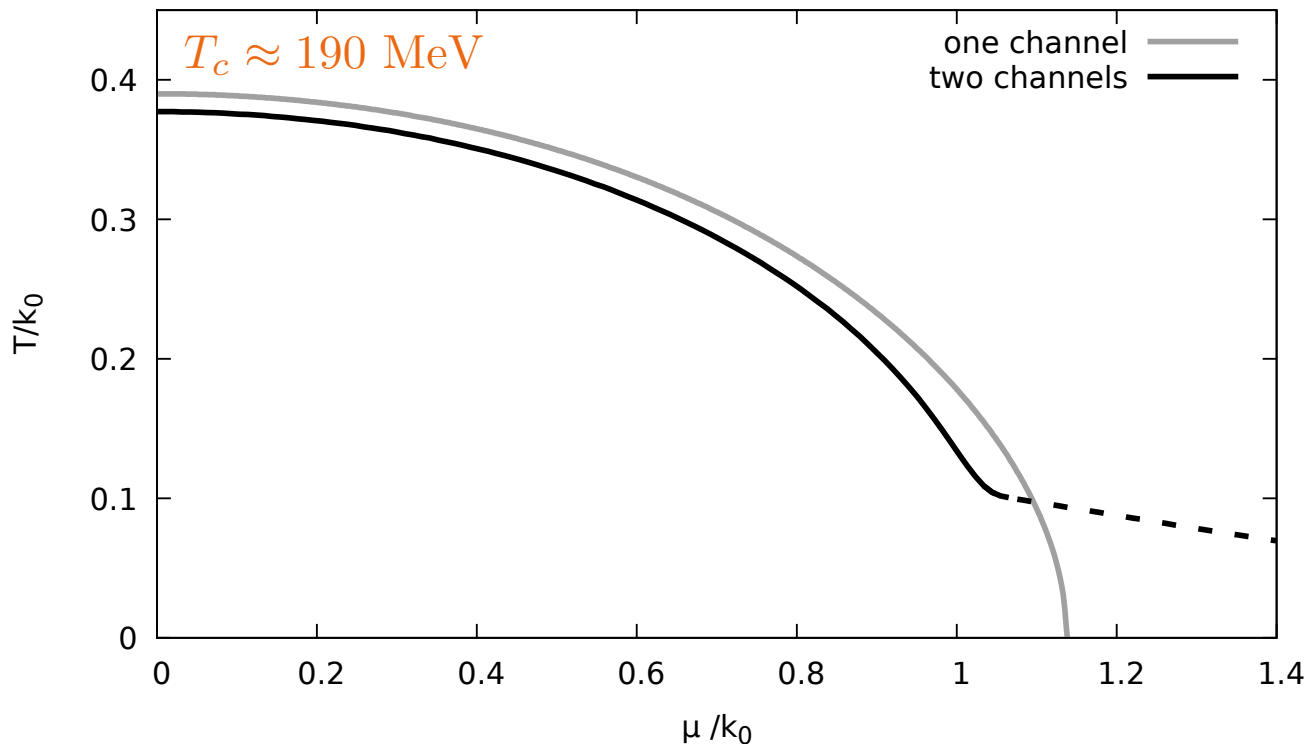
Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking



Ansatz two-channel approximation

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} (\sigma - \pi) + \frac{1}{2} \bar{\lambda}_{\text{csc}} (\text{csc}) \right\}$$



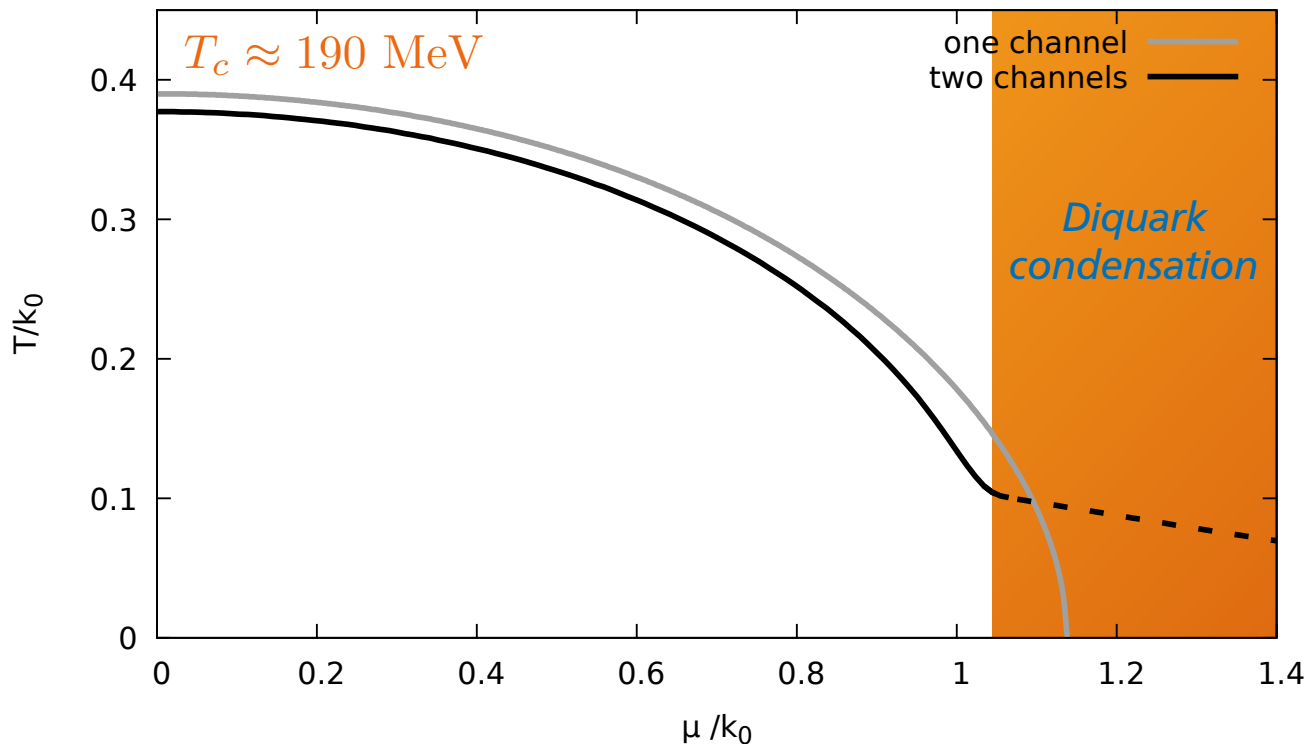
Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking



Ansatz two-channel approximation

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} (\sigma - \pi) + \frac{1}{2} \bar{\lambda}_{\text{CSC}} (\text{CSC}) \right\}$$

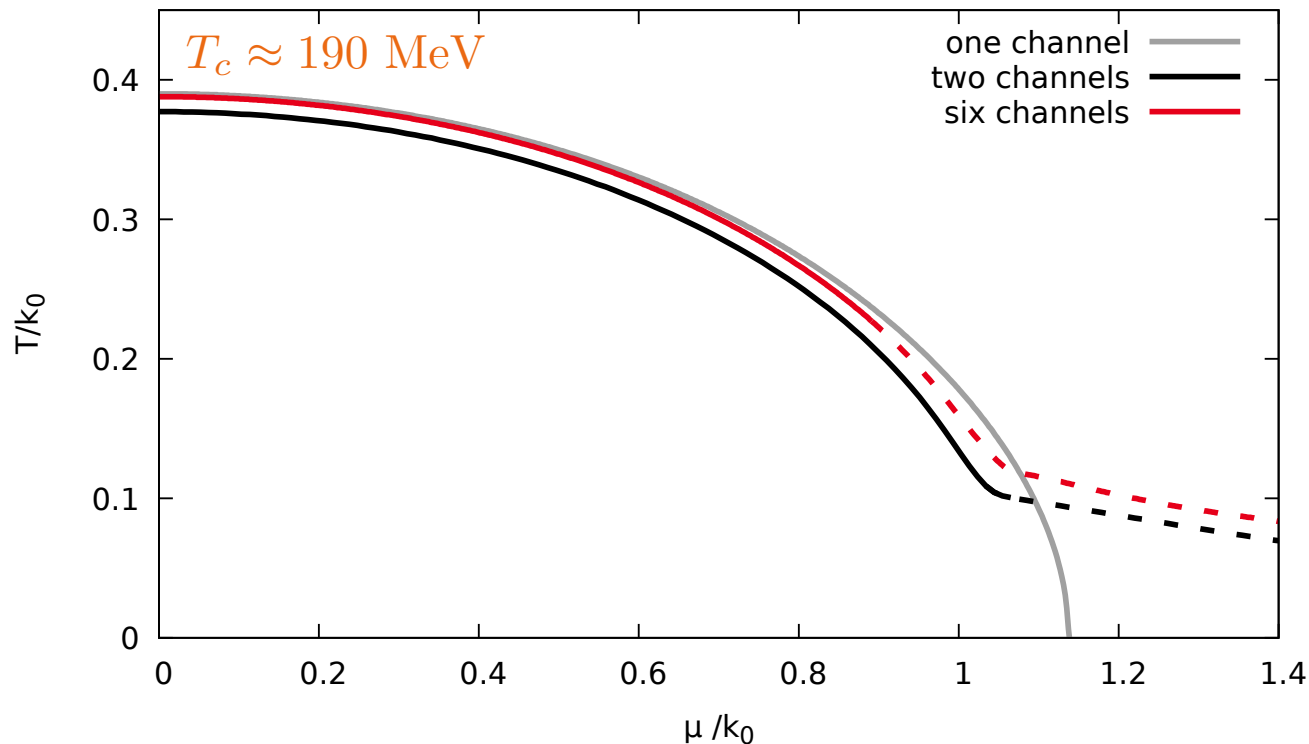


Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking

Ansatz Fierz-complete at $T = 0$ and $\mu = 0$

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \sum_{i=1}^6 \bar{\lambda}_i \mathcal{L}_{(\bar{\psi}\psi)^2}^i \right\} \quad \text{six channels}$$



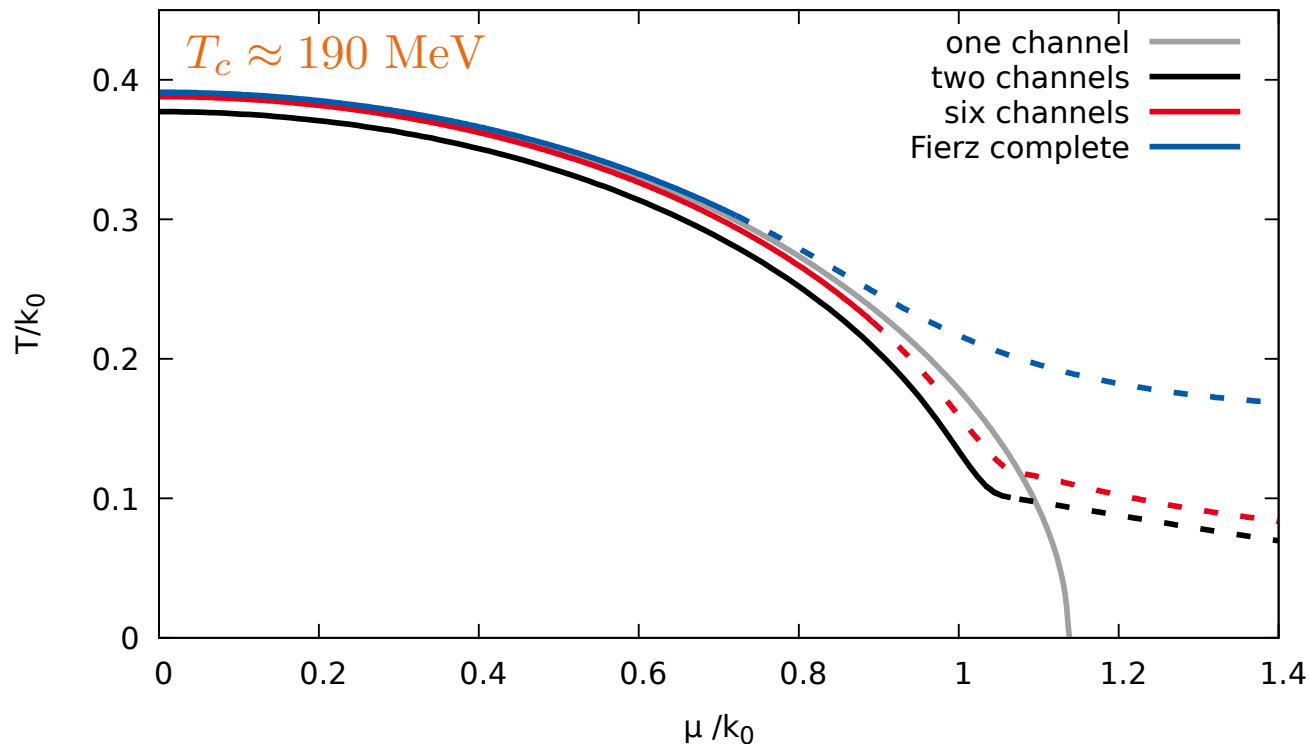
Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking

Fierz-complete ansatz

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \sum_i \bar{\lambda}_i \mathcal{L}_{(\bar{\psi}\psi)^2}^i \right\}$$

comprises all 10 channels

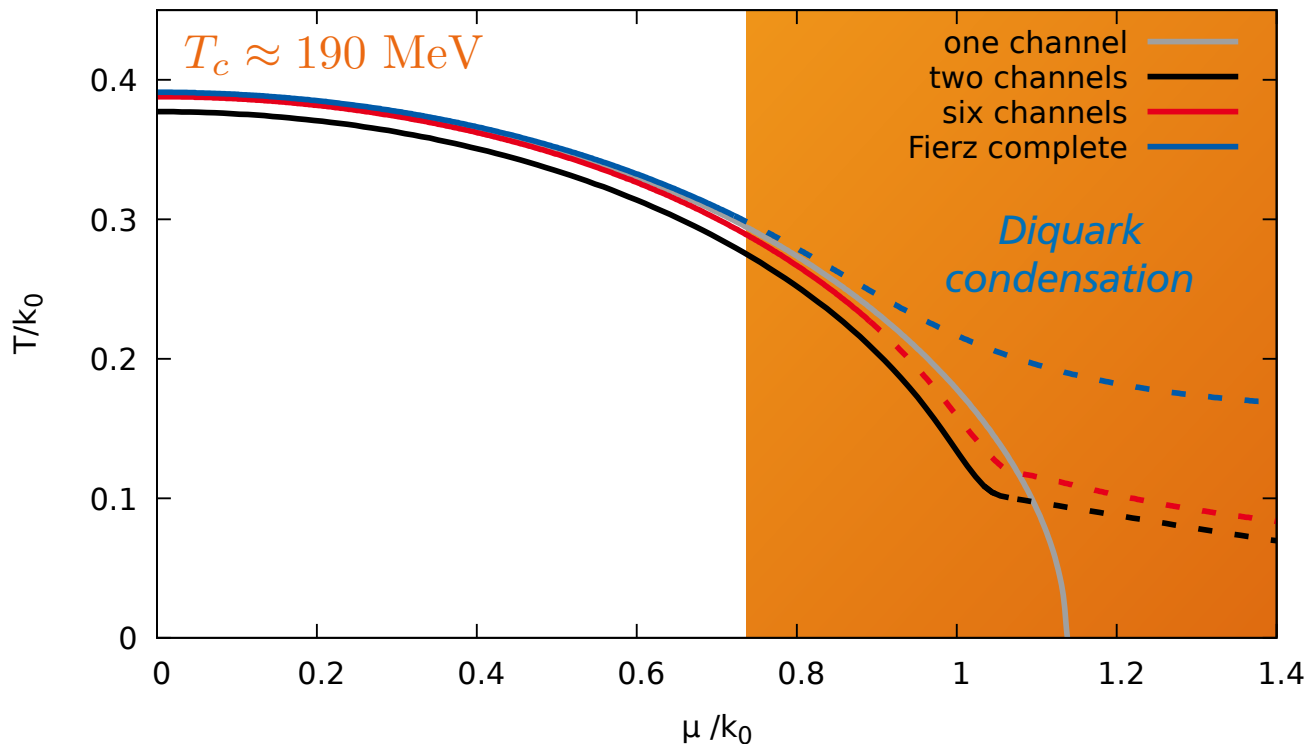


Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking

Fierz-complete ansatz

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \sum_i \bar{\lambda}_i \mathcal{L}_{(\bar{\psi}\psi)^2}^i \right\} \text{ comprises all 10 channels}$$



$T_c \sim \Delta_{\text{CSC}}(T=0)$
[BCS theory]
Implications
for EoS!

Conclusions and outlook

Chiral effective field theory at lower densities

- Efficient Monte-Carlo framework for MBPT (**automatic** code generation; **4th order**)
- Improve fits of LECs by guiding in terms of nuclear saturation
- Nuclear thermodynamics from χ EFT interactions: T, ρ, δ

Outlook Apply saturation guided fitting to next-generation interactions, extract single-particle properties from nuclear thermodynamics

Functional renormalization group at higher densities

- **First Fierz-complete study** of effective action
- **Importance** of Fierz-completeness to probe the regime at **high quark chemical potential** and low temperature
- Forming of diquark condensate (color superconducting phase)

Outlook Inclusion of dynamic gauge fields (equations worked out) and first estimate of EoS, work in progress.

Backup

Difermion-type degrees of freedom

$$\begin{Bmatrix} (S - P) \\ (V_{\parallel}) \\ (V_{\perp}) \end{Bmatrix}$$

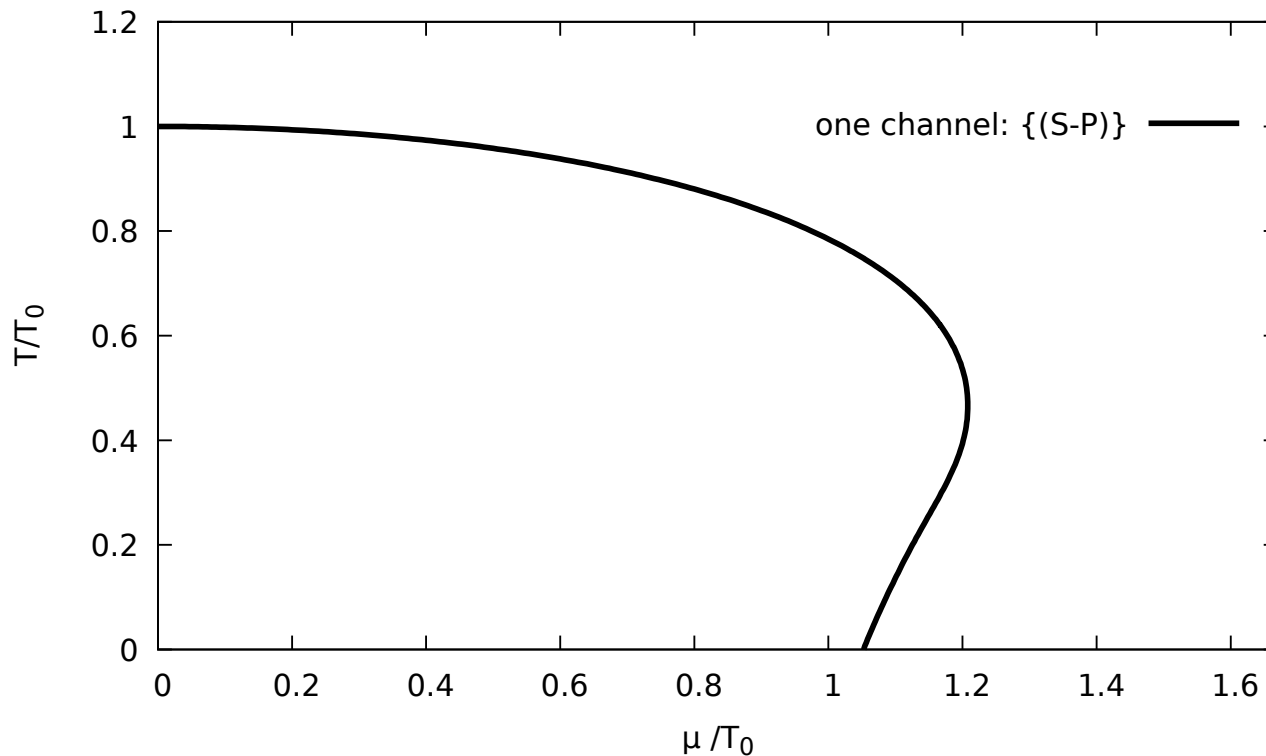


$$\begin{Bmatrix} (S - P) \\ (SC - PC) \\ (A_{\parallel}C) \end{Bmatrix}$$

$$C = i\gamma_2\gamma_0$$

$$(SC - PC) \equiv (\bar{\psi}C\bar{\psi}^T)(\psi^T C\psi) - (\bar{\psi}\gamma_5 C\bar{\psi}^T)(\psi^T C\gamma_5\psi)$$

$$(A_{\parallel}C) \equiv (\bar{\psi}\gamma_0\gamma_5 C\bar{\psi}^T)(\psi^T C\gamma_0\gamma_5\psi)$$



Difermion-type degrees of freedom

$$\begin{pmatrix} (S - P) \\ (V_{\parallel}) \\ (V_{\perp}) \end{pmatrix}$$

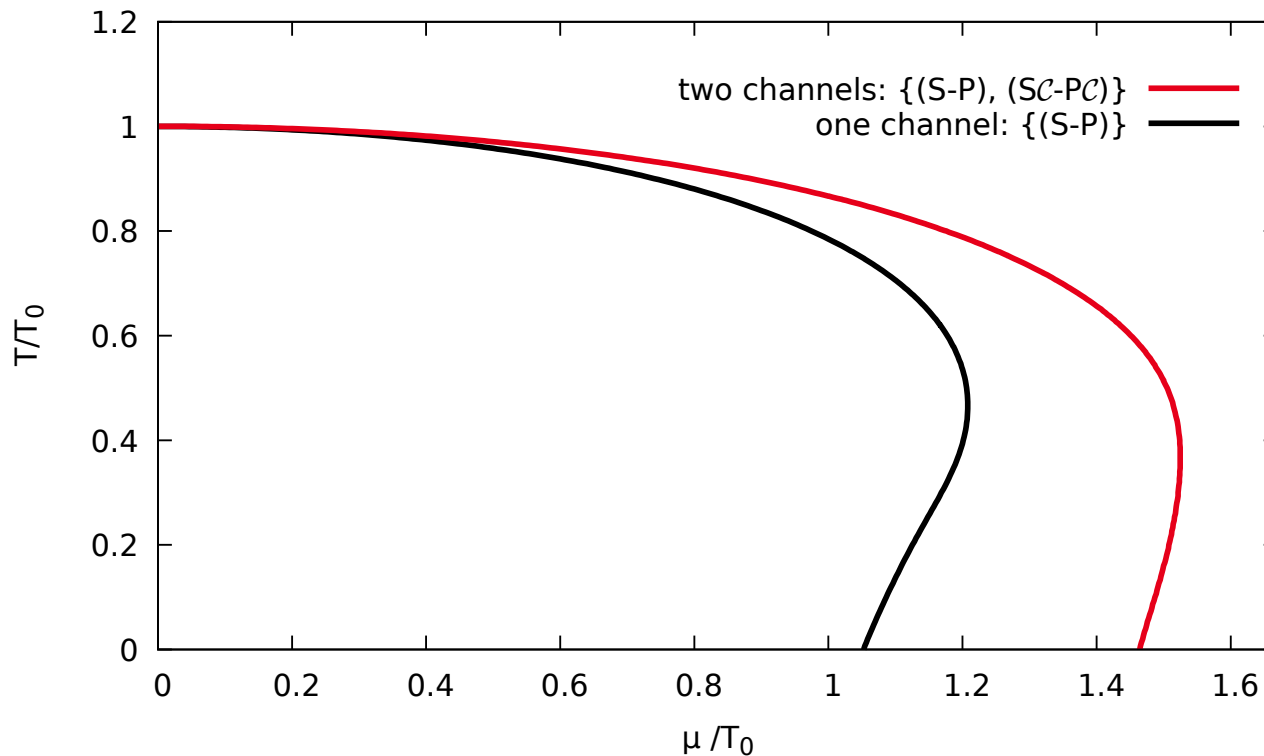


$$\begin{pmatrix} (S - P) \\ (SC - PC) \\ (A_{\parallel}C) \end{pmatrix}$$

$$C = i\gamma_2\gamma_0$$

$$(SC - PC) \equiv (\bar{\psi}C\bar{\psi}^T)(\psi^T C\psi) - (\bar{\psi}\gamma_5 C\bar{\psi}^T)(\psi^T C\gamma_5\psi)$$

$$(A_{\parallel}C) \equiv (\bar{\psi}\gamma_0\gamma_5 C\bar{\psi}^T)(\psi^T C\gamma_0\gamma_5\psi)$$



Difermion-type degrees of freedom

$$\begin{Bmatrix} (S - P) \\ (V_{\parallel}) \\ (V_{\perp}) \end{Bmatrix}$$

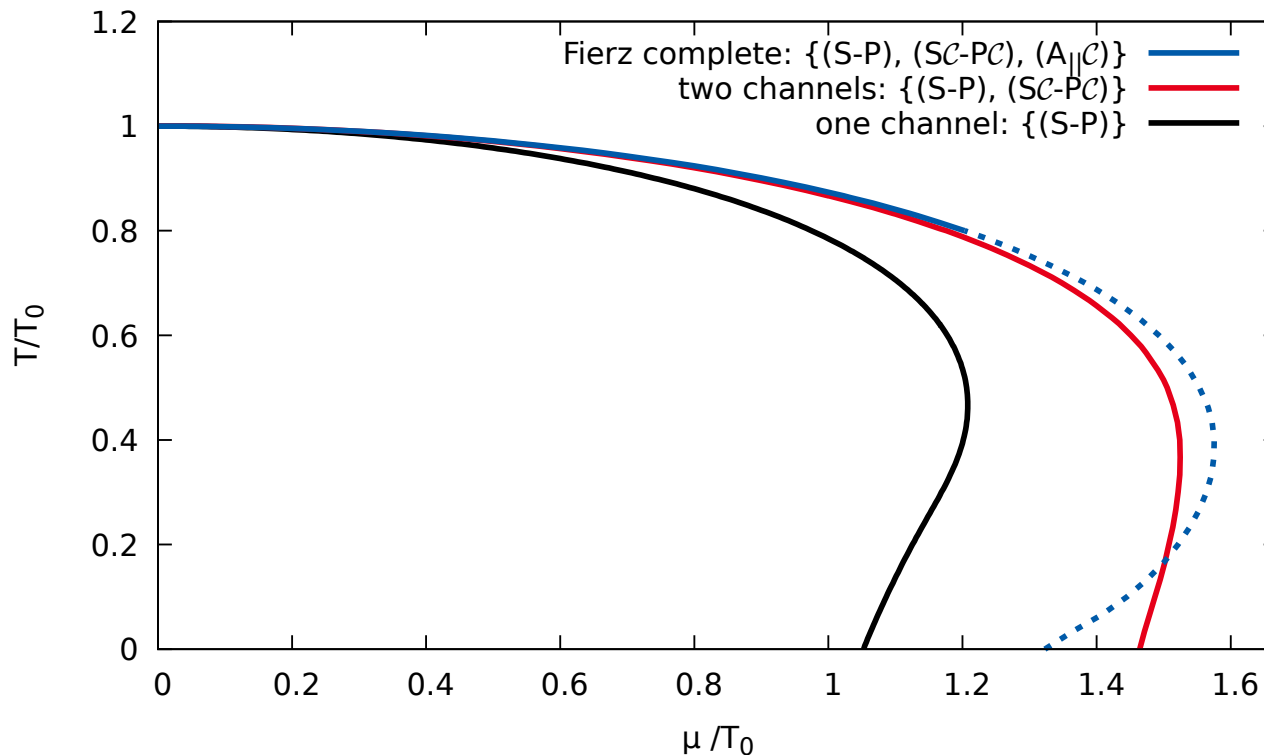


$$\begin{Bmatrix} (S - P) \\ (SC - PC) \\ (A_{\parallel}C) \end{Bmatrix}$$

$$C = i\gamma_2\gamma_0$$

$$(SC - PC) \equiv (\bar{\psi}C\bar{\psi}^T)(\psi^T C\psi) - (\bar{\psi}\gamma_5 C\bar{\psi}^T)(\psi^T C\gamma_5\psi)$$

$$(A_{\parallel}C) \equiv (\bar{\psi}\gamma_0\gamma_5 C\bar{\psi}^T)(\psi^T C\gamma_0\gamma_5\psi)$$



Backup

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{(S+P)_-^{\text{adj}}} \quad \mathcal{L}_{(\bar{\psi}\psi)^2}^{(V+A)_\parallel} \quad \mathcal{L}_{(\bar{\psi}\psi)^2}^{(V+A)_\perp} \quad \mathcal{L}_{(\bar{\psi}\psi)^2}^{(V-A)_\parallel} \quad \mathcal{L}_{(\bar{\psi}\psi)^2}^{(V-A)_\perp} \quad \mathcal{L}_{(\bar{\psi}\psi)^2}^{(V+A)_\parallel^{\text{adj}}} \quad \mathcal{L}_{(\bar{\psi}\psi)^2}^{(V-A)_\perp^{\text{adj}}}$$

Four-quark interactions and symmetries

Fierz-complete basis of interactions



$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ \bar{\psi} (iZ_{\parallel} \gamma_0 \partial_0 + iZ_{\perp} \gamma_i \partial_i - iZ_{\mu} \mu \gamma_0) \psi + \frac{1}{2} \sum_i \bar{\lambda}_i \mathcal{L}_{(\bar{\psi}\psi)^2}^i \right\}$$

In total **20 channels** meet
symmetry constraints

Fierz identities →

Fierz-complete basis:
10 channels

$U_A(1)$ breaking channels:

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{(\sigma-\pi)} = (\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2$$

↔ formation of
chiral condensate

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{\text{csc}} \sim (i\bar{\psi}\gamma_5\tau_A t_c^{A'} \mathcal{C}\bar{\psi}^T)(i\psi^T \mathcal{C}\gamma_5\tau_A t_c^{A'} \psi) \quad J^P = 0^+$$

[Rapp, Schäfer, Shuryak, Velkovsky, 1998]

↔ formation of
diquark condensate

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{\text{det}} = (\bar{\psi}\psi)^2 + (\bar{\psi}\gamma_5\psi)^2 - (\bar{\psi}\vec{\tau}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2$$

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{(\text{S+P})\text{-adj}} = (\bar{\psi}T^a\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}T^a\psi)^2 + (\bar{\psi}\gamma_5T^a)^2 - (\bar{\psi}\vec{\tau}T^a\psi)^2$$

+ six $U_A(1)$ symmetric channels

[Braun, ML, Pospiech '17]



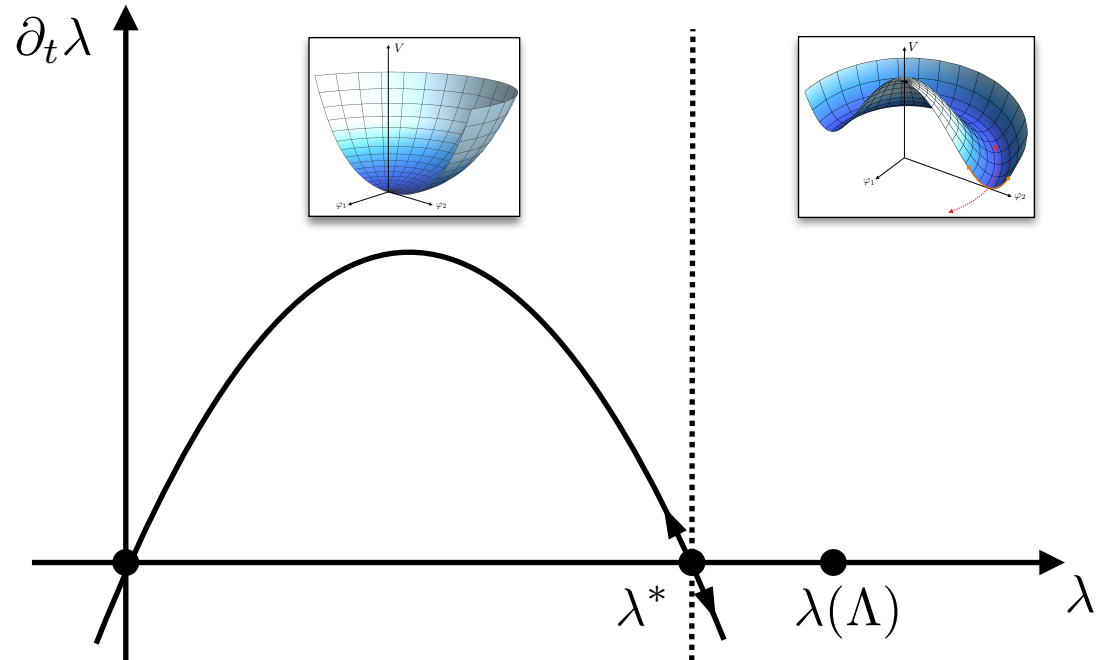
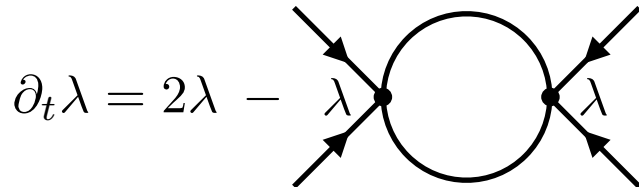
RG flow of four-quark interactions

Qualitative behavior and the effect of external parameters



TECHNISCHE
UNIVERSITÄT
DARMSTADT

RG flow equation:



Scale fixing procedure:

- Only $\lambda_{(\sigma-\pi)}(\Lambda)$ assumes **finite** value as inspired by gluon-induced four-quark flows [Braun, 2011]
- UV value tuned so that specific **scale k_0 of symmetry-breakdown** is obtained
(defined by $1/\lambda(k_0) = 0$, sets the scale for low-energy observables)
- One-channel approximation can be mapped onto mean-field gap-equation
to access deep infrared: $m_q(k_0) \approx 300$ MeV, $m_\sigma(k_0) \approx 800$ MeV

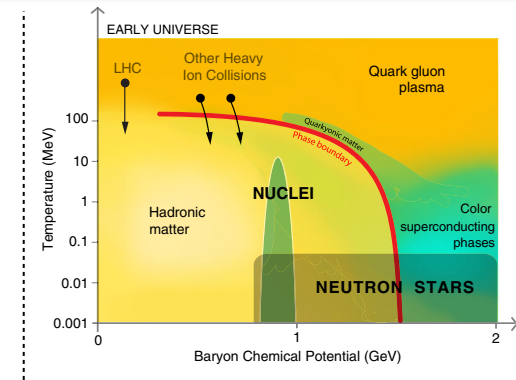
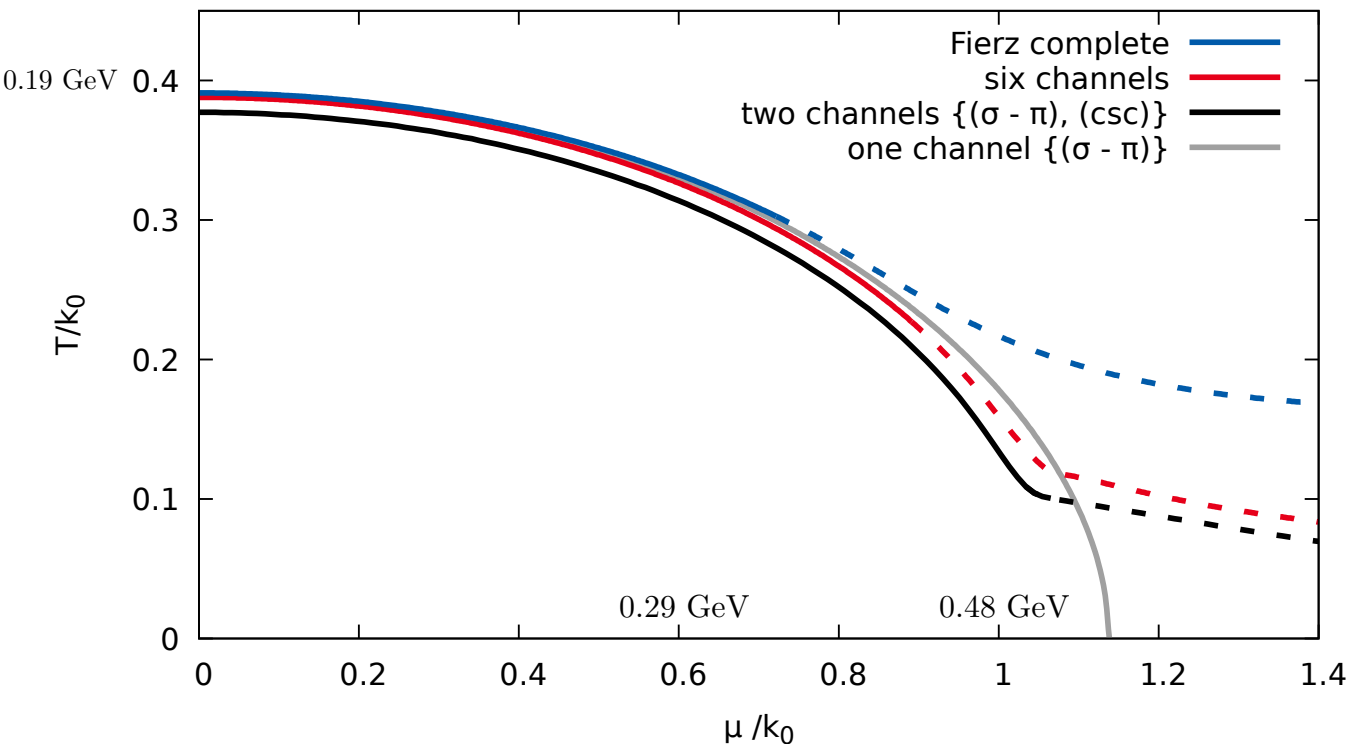


Structure of the phase boundary

$$\lambda_{UV}^{\sigma-\pi} \approx 7.317, \quad \lambda_{UV}^{(i)} = 0 \text{ for } i \neq \sigma-\pi \quad \longrightarrow \quad k_0/\Lambda \approx 0.483 \quad \longleftrightarrow \quad m_\psi \approx 0.3 \text{ GeV}$$

[Braun, 2006]

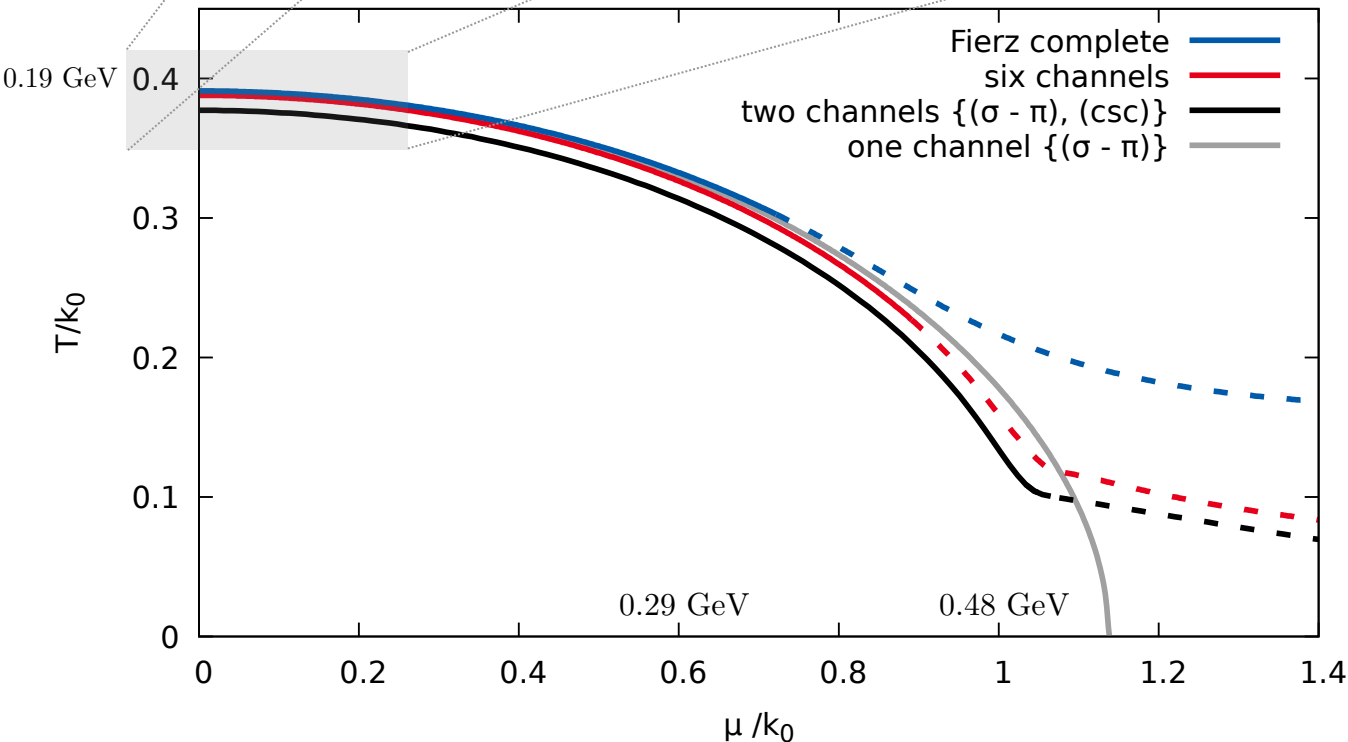
[Mitter, Pawłowski, Strodthoff, 2014]



Structure of the phase boundary

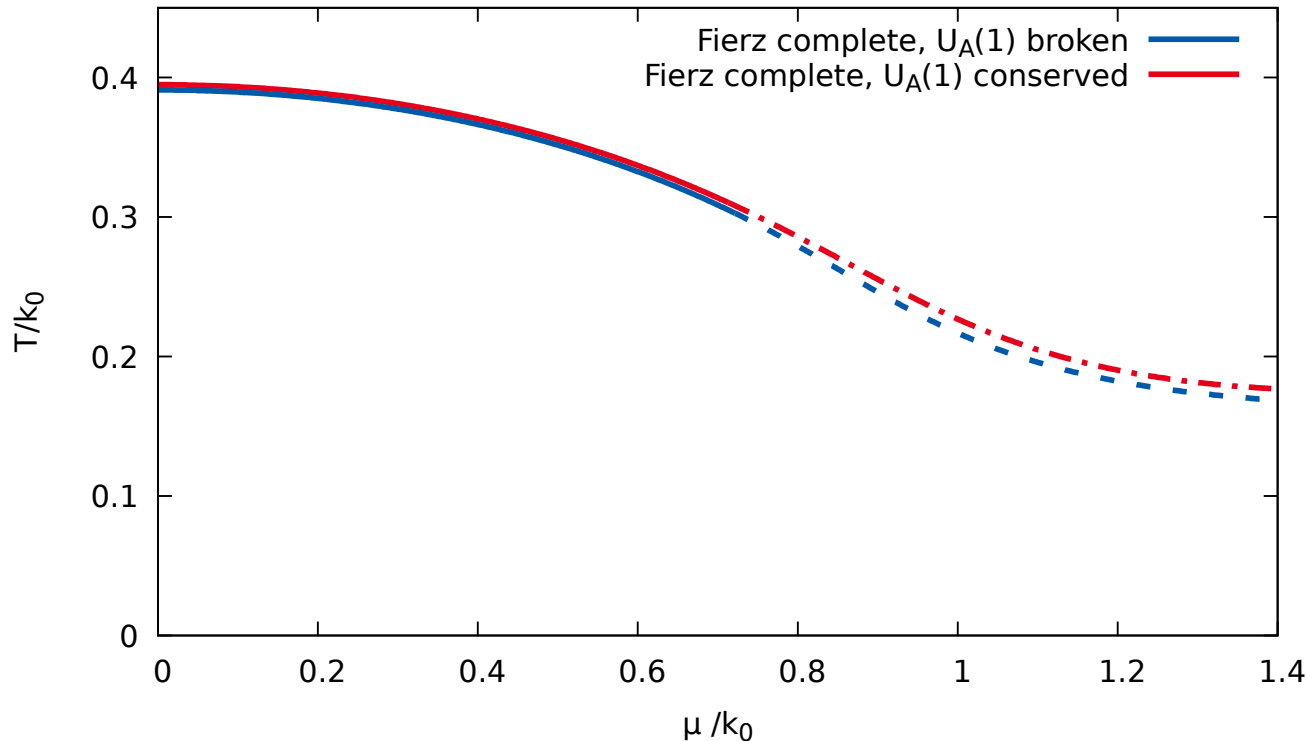
$$\text{Curvature } \kappa := -T_\chi(0) \left. \frac{dT_\chi(\mu^2)}{d\mu^2} \right|_{\mu=0}$$

$$\kappa \approx 0.06$$



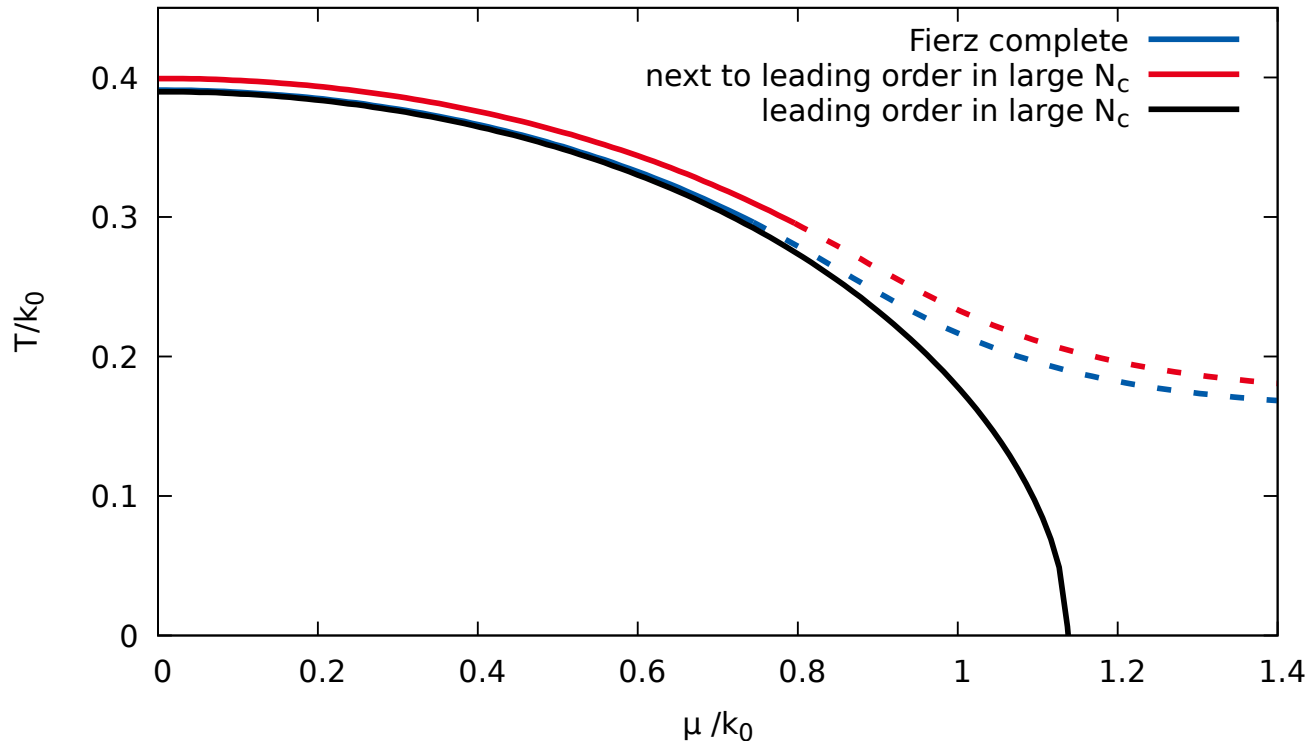
Structure of the phase boundary

$$\lambda_{UV}^{\sigma-\pi} \approx 7.317, \quad \lambda_{UV}^{(i)} = 0 \text{ for } i \neq \sigma-\pi \quad \longrightarrow \quad k_{cr}/\Lambda \approx 0.483 \quad \longleftrightarrow \quad m_\psi/\Lambda \approx 0.3$$



Structure of the phase boundary

$$\lambda_{UV}^{\sigma-\pi} \approx 7.317, \quad \lambda_{UV}^{(i)} = 0 \text{ for } i \neq \sigma-\pi \quad \longrightarrow \quad k_{cr}/\Lambda \approx 0.483 \quad \longleftrightarrow \quad m_\psi/\Lambda \approx 0.3$$

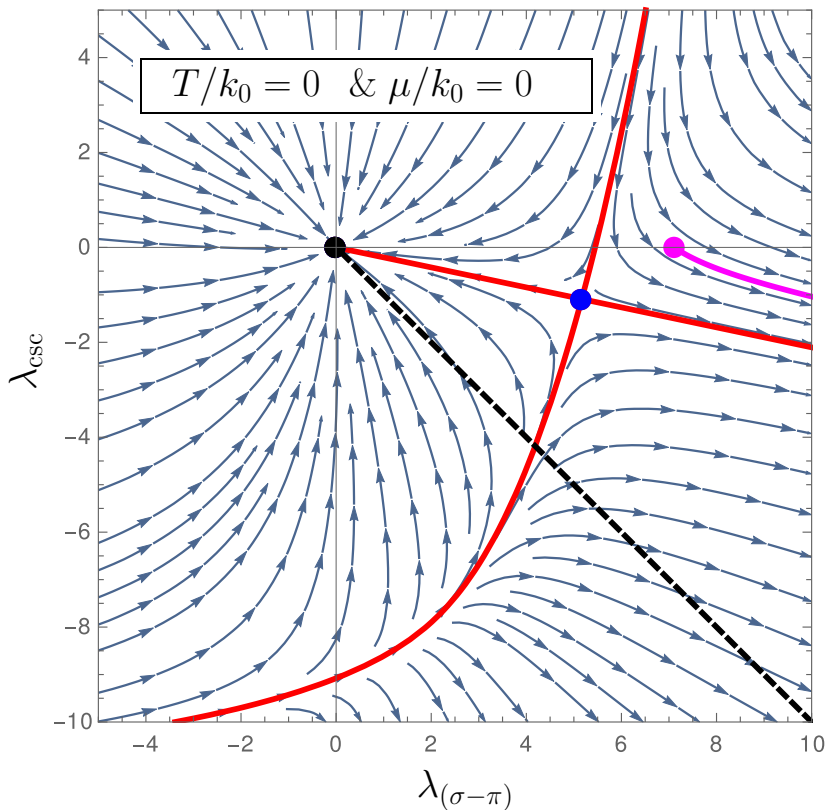


Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking

Ansatz two-channel approximation

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} (\sigma - \pi) + \frac{1}{2} \bar{\lambda}_{\text{csc}} (\text{csc}) \right\}$$

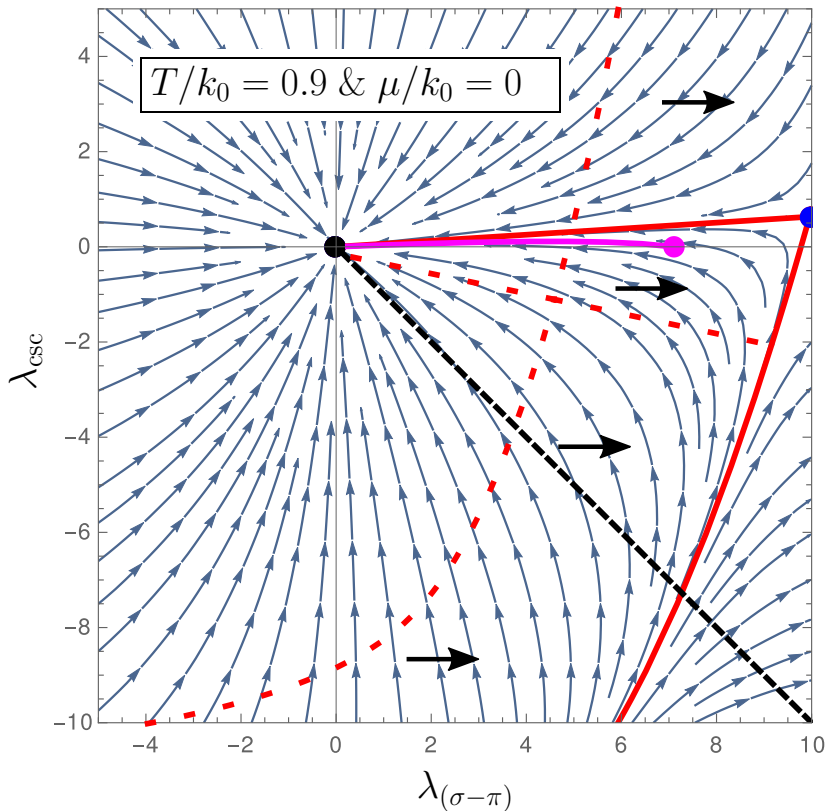


Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking

Ansatz two-channel approximation

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} (\sigma - \pi) + \frac{1}{2} \bar{\lambda}_{\text{csc}} (\text{csc}) \right\}$$

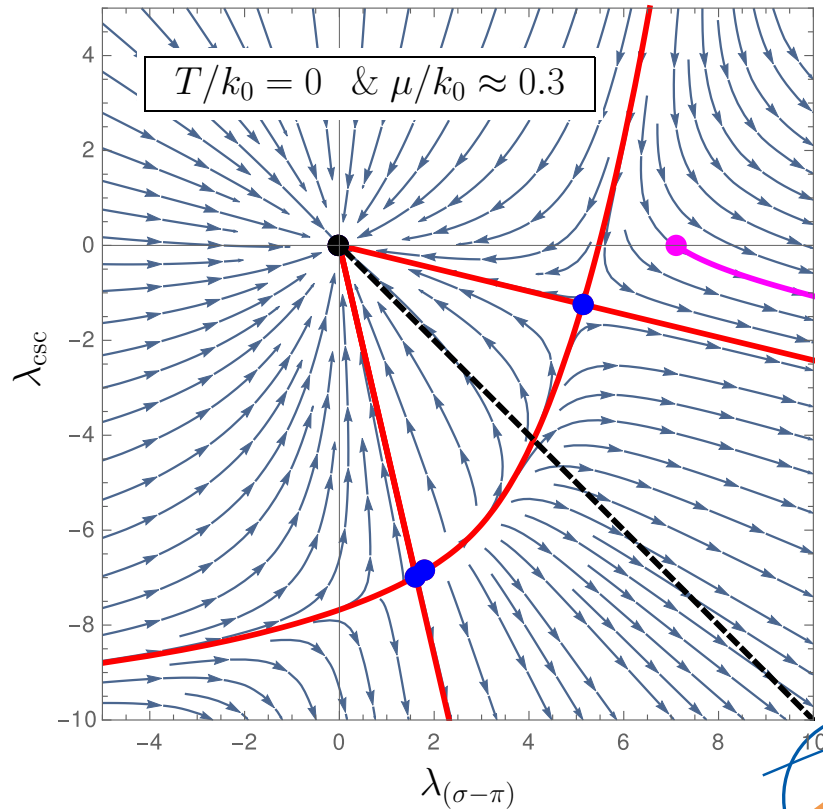
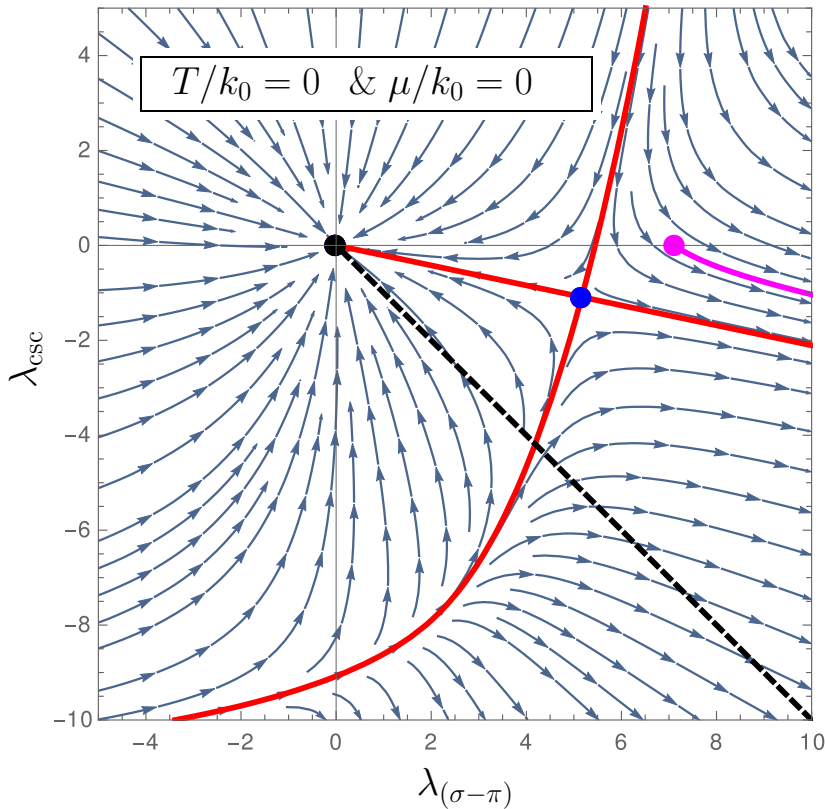


Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking

Ansatz two-channel approximation

$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} (\sigma - \pi) + \frac{1}{2} \bar{\lambda}_{\text{csc}} (\text{csc}) \right\}$$

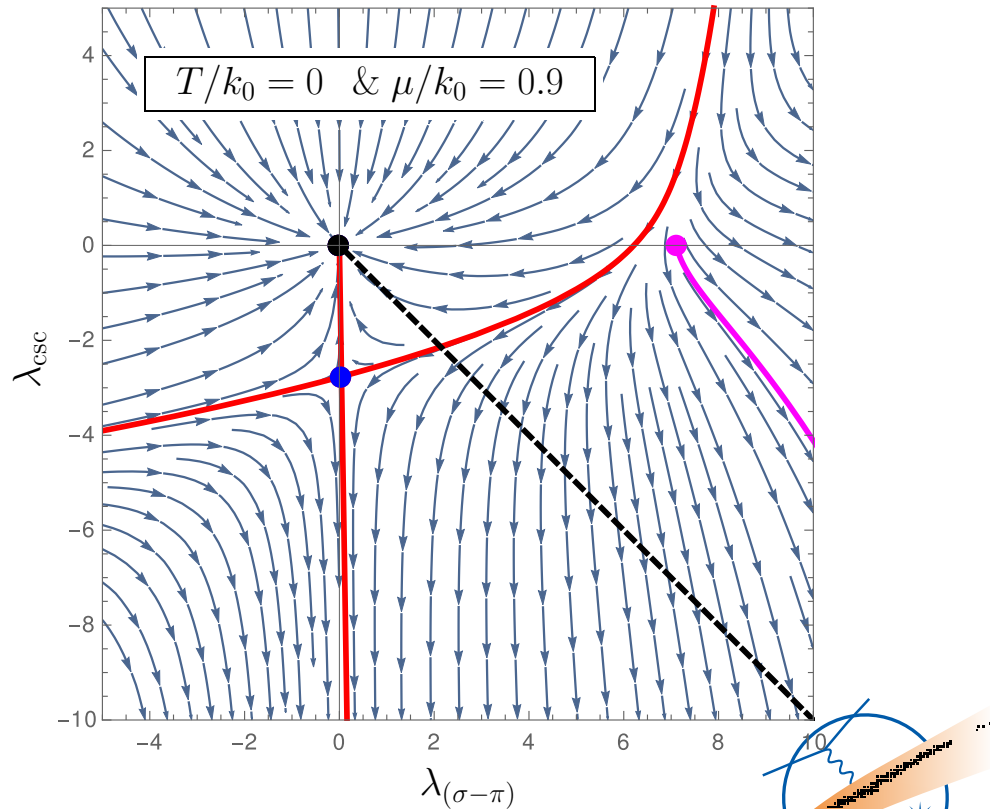
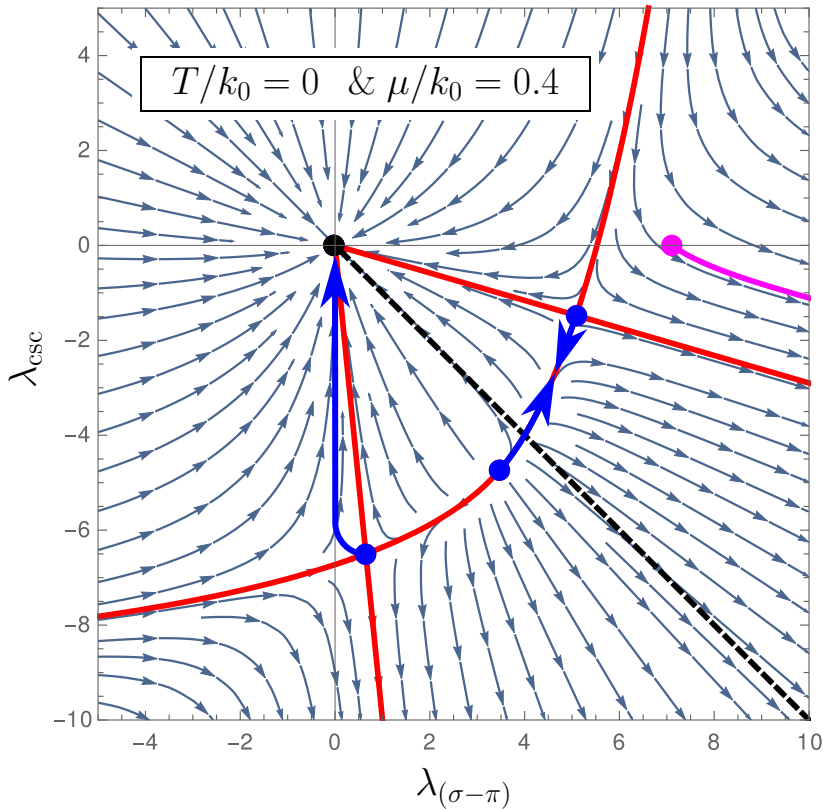


Exploring the phase diagram

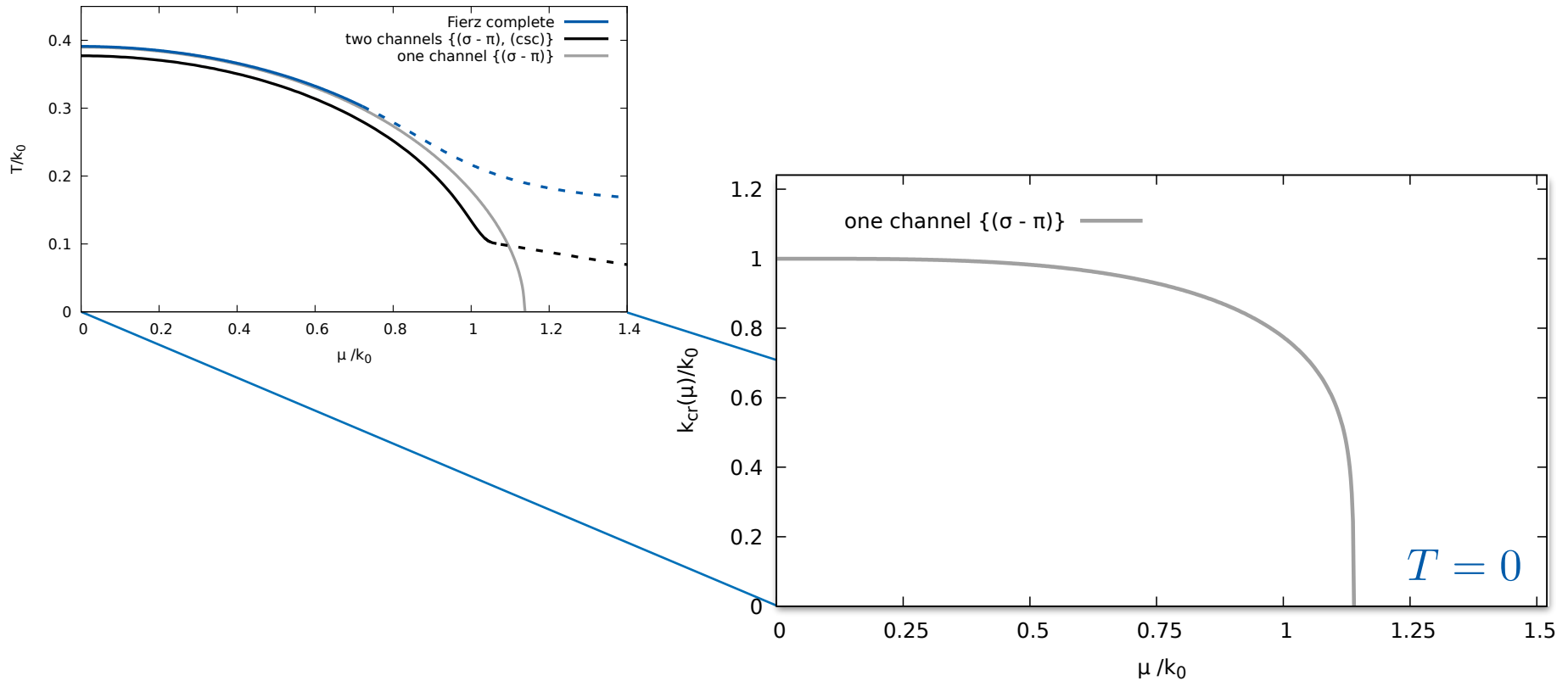
Fixed-point structure and patterns of symmetry breaking

Ansatz two-channel approximation

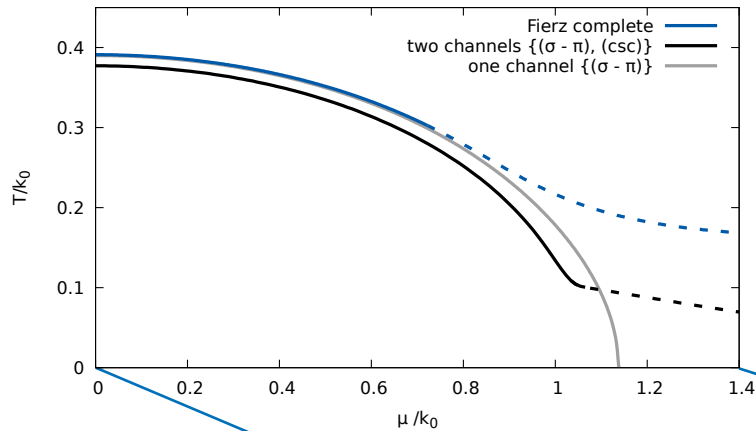
$$\Gamma_k[\bar{\psi}, \psi] = \int_x \left\{ (\text{kinetic term}) + \frac{1}{2} \bar{\lambda}_{(\sigma-\pi)} (\sigma - \pi) + \frac{1}{2} \bar{\lambda}_{\text{csc}} (\text{csc}) \right\}$$



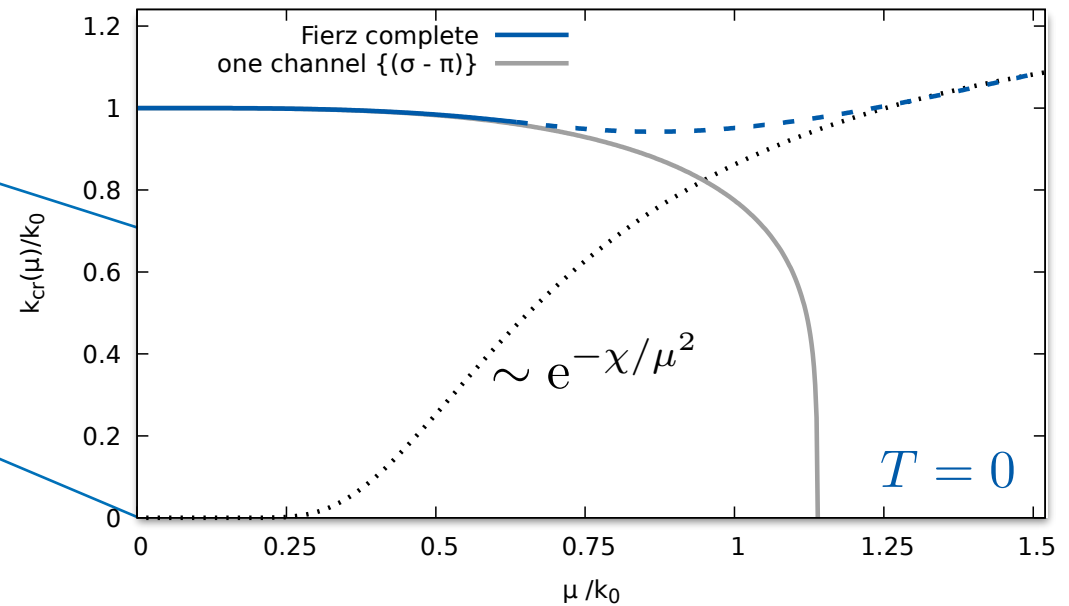
Structure of the phase boundary



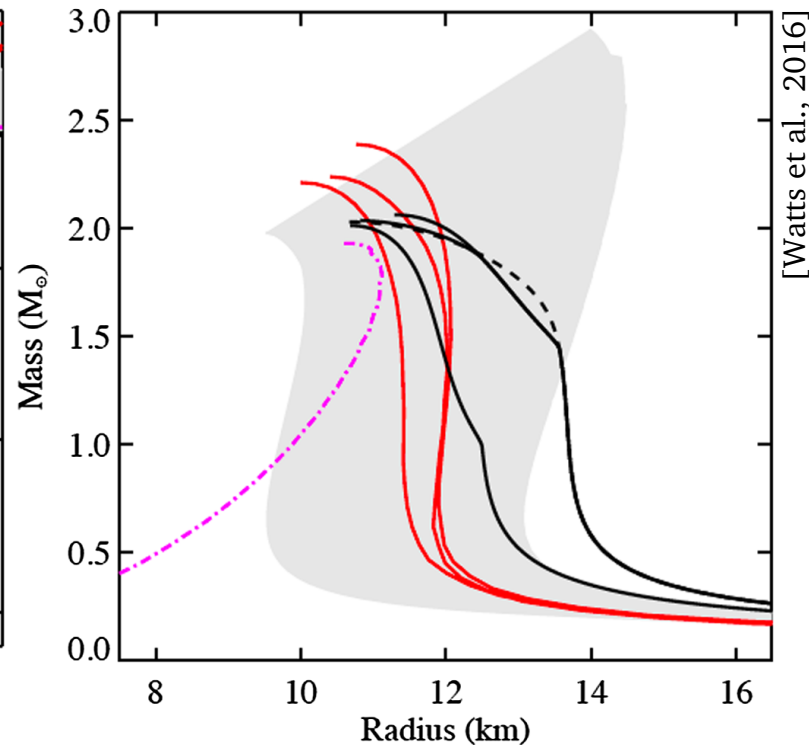
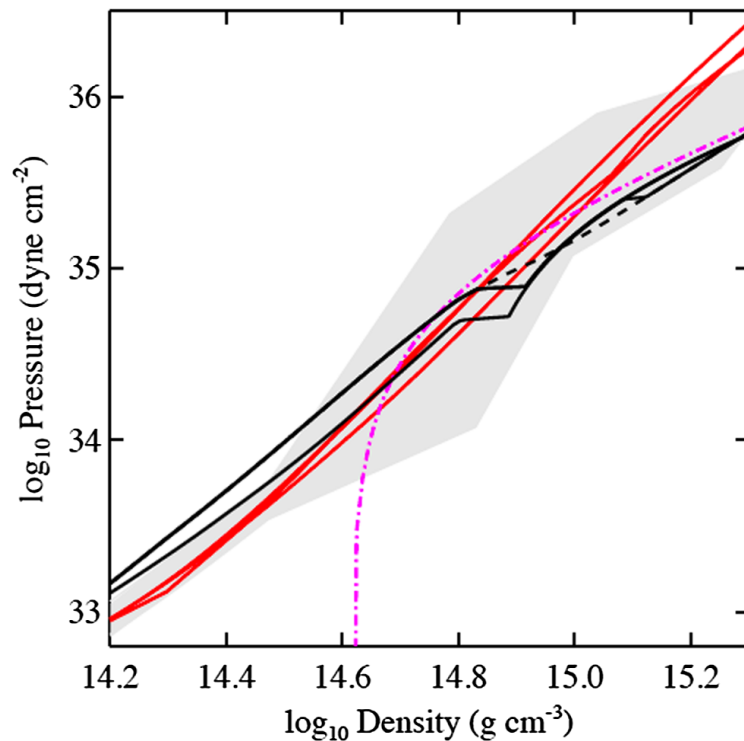
Structure of the phase boundary



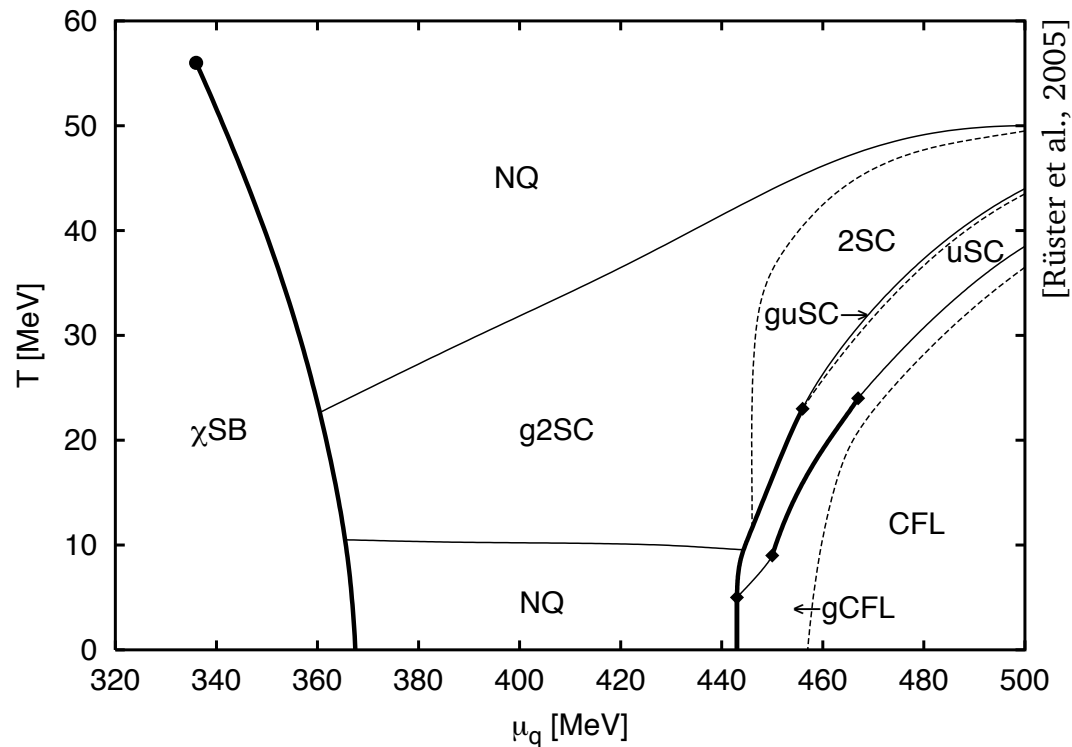
BCS-type exponential scaling



QCD phase diagram: Nuclear matter EOS and neutron stars



QCD phase diagram: Nuclear matter EOS and neutron stars



Symmetrie breaking and four-quark interactions in QCD

Structure of the phase boundary at finite temperature and density



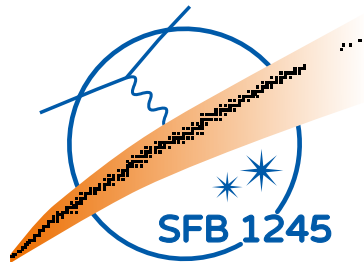
TECHNISCHE
UNIVERSITÄT
DARMSTADT

Marc Leonhardt

Institut für Kernphysik, Technische Universität Darmstadt

with Martin Pospiech and Jens Braun

CRC 1245 Seminar
Integrated Research Training Group



HIC | **FAIR**
for
Helmholtz International Center

QCD phase diagram ... and project B05

Nuclear matter equation of state for astrophysical applications

*“[...] complementary approaches, chiral effective field theory at lower densities and the **functional renormalization group starting from quark-gluon dynamics at higher densities**, to obtain a quantitative determination of the **nuclear matter equation of state** over a wide range of densities, temperatures, and proton fractions.”*

Recap: QFT concepts

All physical information is stored in **correlation functions/n-point functions**.

$$\langle \phi(x_1) \dots \phi(x_n) \rangle := \mathcal{N} \int \mathcal{D}\phi \phi(x_1) \dots \phi(x_n) e^{-S[\phi]}$$

e.g. scattering amplitude (S-matrix elements) via LSZ reduction formula

Statistical Physics

QFT

$$Z[J] = \int \mathcal{D}\phi e^{-S[\phi] + \int J\phi}$$

Partition function

Generating functional

$$W[J] = \log Z[J]$$

Helmholtz free energy

Generating functional of connected diagrams

Effective action $\Gamma[\Phi]$

Gibbs free energy

Generating functional of 1PI diagrams

Functional renormalization group (FRG)

Flow from high to low energies in QCD



TECHNISCHE
UNIVERSITÄT
DARMSTADT



[adapted from
H. Gies, 2006]

Flow equation
[Wetterich, 1993]

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

“RG time”
 $t = \ln(k/\Lambda)$



Functional renormalization group (FRG)

Flow from high to low energies in QCD



TECHNISCHE
UNIVERSITÄT
DARMSTADT



$$\text{UV: } \Gamma_k \xrightarrow{k \rightarrow \Lambda} S$$

$$\text{IR: } \Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$$

[adapted from
H. Gies, 2006]

Flow equation
[Wetterich, 1993]

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

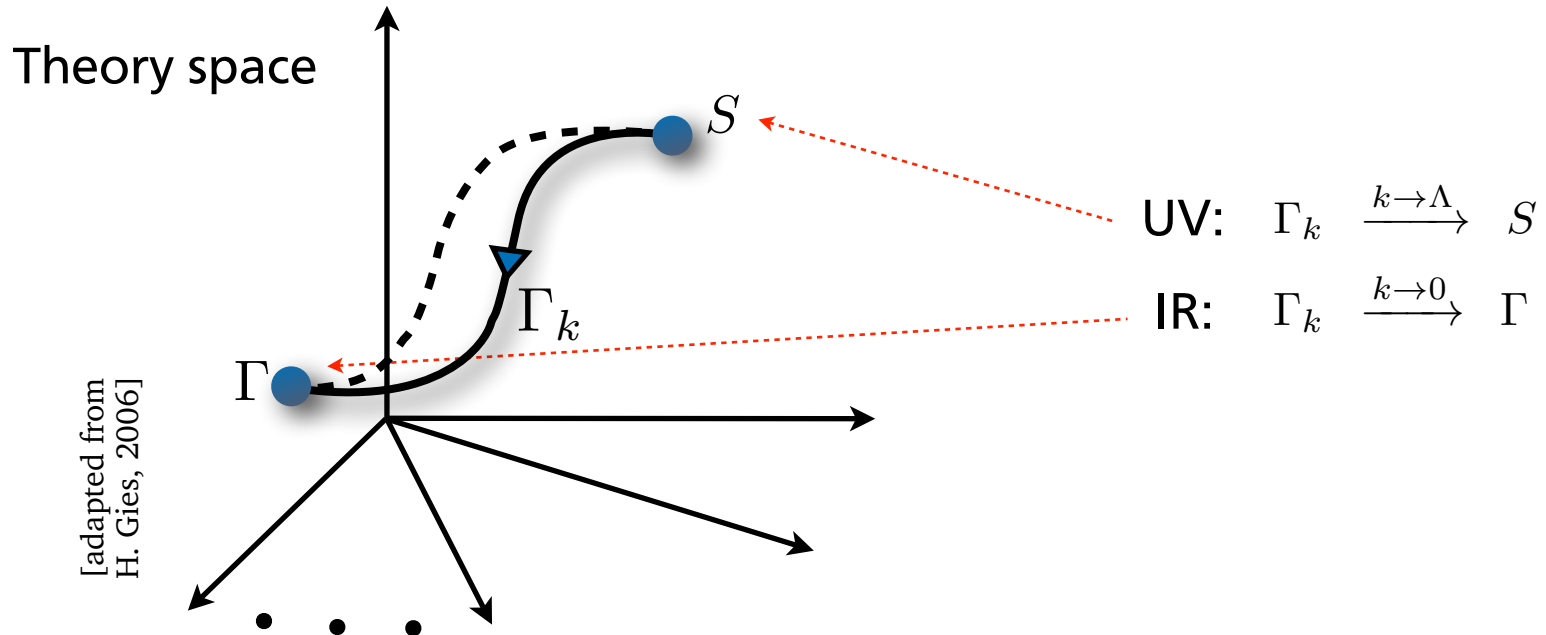
“RG time”
 $t = \ln(k/\Lambda)$



Functional renormalization group (FRG)

Flow from high to low energies in QCD

$$\Gamma \xrightarrow{R_k} \text{Effective average action } \Gamma_k$$



Flow equation
[Wetterich, 1993]

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

“RG time”
 $t = \ln(k/\Lambda)$



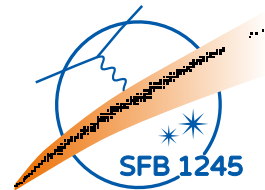
Functional renormalization group (FRG)

Flow from high to low energies in QCD



TECHNISCHE
UNIVERSITÄT
DARMSTADT

$$\Gamma \xrightarrow{R_k} \text{Effective average action } \Gamma_k$$



Functional renormalization group (FRG)

Flow from high to low energies in QCD

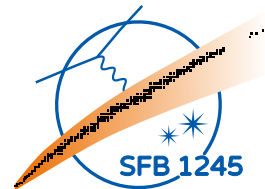


TECHNISCHE
UNIVERSITÄT
DARMSTADT

$\Gamma \xrightarrow{R_k}$ Effective average action Γ_k

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$

IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$

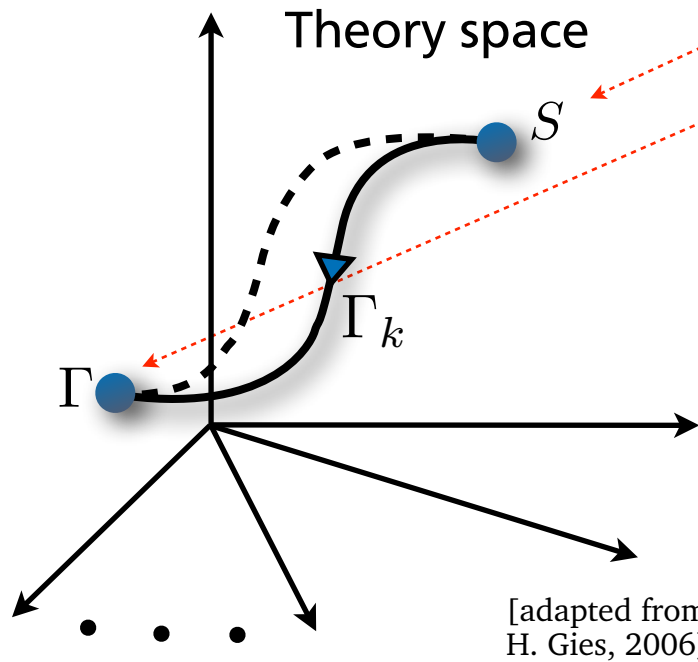


Functional renormalization group (FRG)

Flow from high to low energies in QCD

$$\Gamma \xrightarrow{R_k} \text{Effective average action } \Gamma_k$$

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$
IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$



[adapted from
H. Gies, 2006]

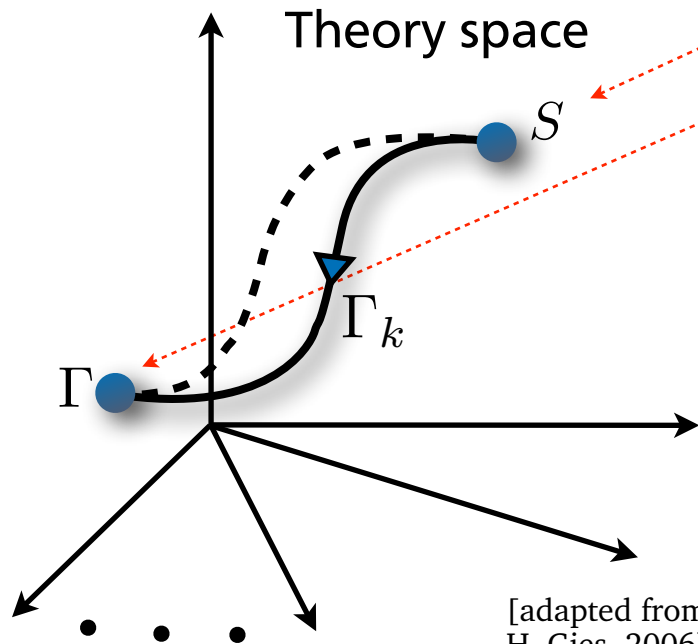
Functional renormalization group (FRG)

Flow from high to low energies in QCD

$\Gamma \xrightarrow{R_k}$ Effective average action Γ_k

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$

IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$



[adapted from
H. Gies, 2006]

Flow
equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

$$t = \ln(k/\Lambda)$$

[C.Wetterich, *Phys. Lett. B*, 301, 1993]



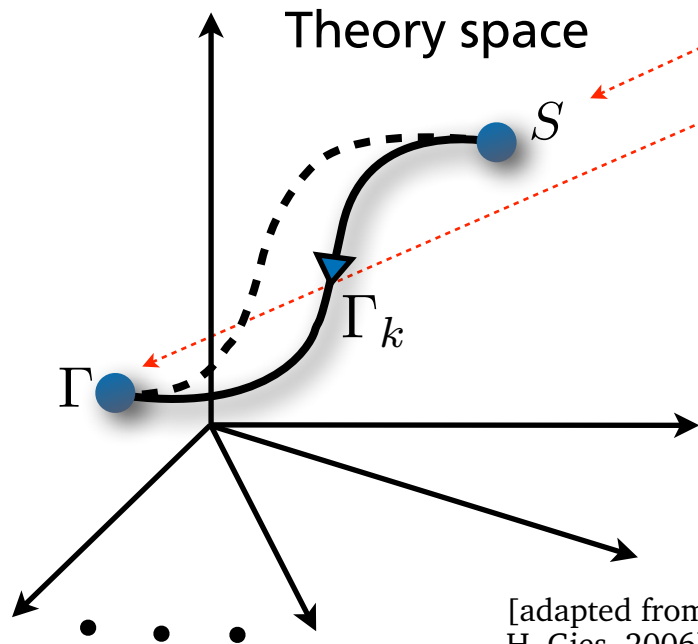
Functional renormalization group (FRG)

Flow from high to low energies in QCD

$\Gamma \xrightarrow{R_k}$ Effective average action Γ_k

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$

IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$



[adapted from
H. Gies, 2006]

*Everything that is not
forbidden is allowed.*

Flow
equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

$$t = \ln(k/\Lambda)$$

[C.Wetterich, *Phys. Lett. B*, 301, 1993]



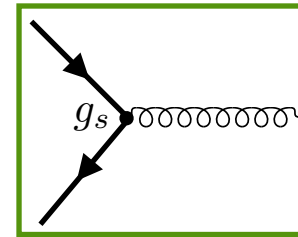
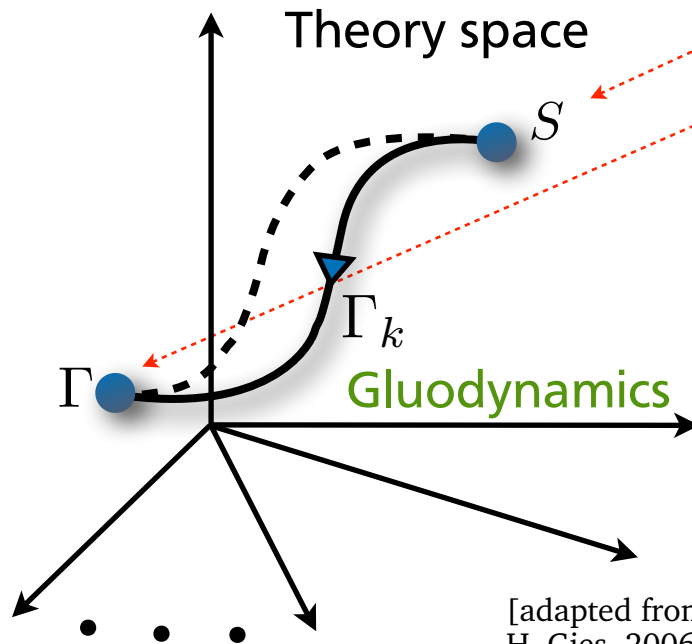
Functional renormalization group (FRG)

Flow from high to low energies in QCD

$\Gamma \xrightarrow{R_k}$ Effective average action Γ_k

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$

IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$



Everything that is not forbidden is allowed.

[adapted from
H. Gies, 2006]

Flow
equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

$$t = \ln(k/\Lambda)$$

[C.Wetterich, *Phys. Lett. B*, 301, 1993]

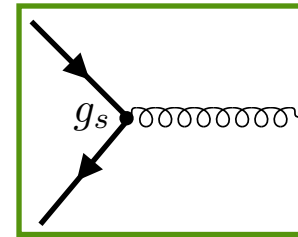
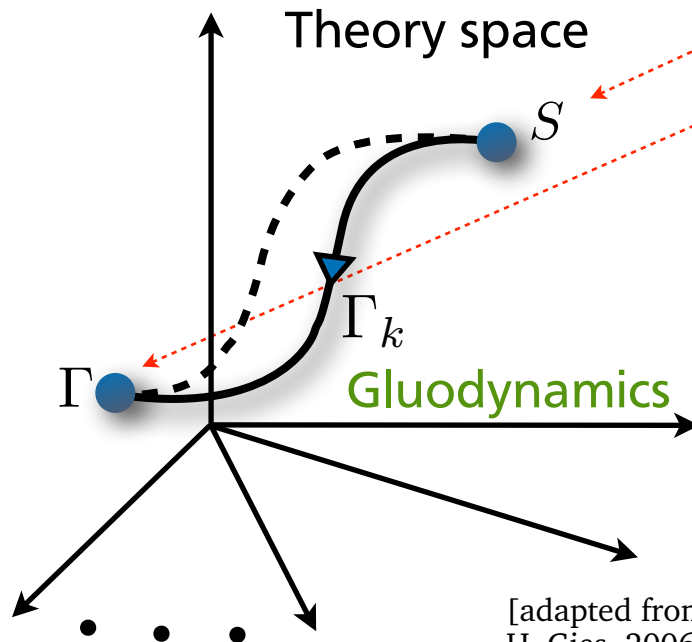
Functional renormalization group (FRG)

Flow from high to low energies in QCD

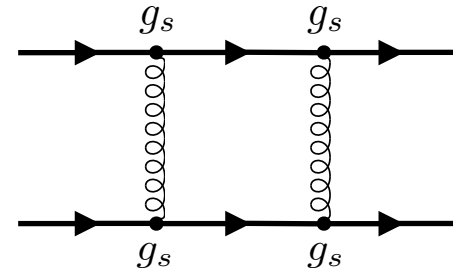
$\Gamma \xrightarrow{R_k}$ Effective average action Γ_k

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$

IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$



Everything that is not forbidden is allowed.



[adapted from
H. Gies, 2006]

Flow
equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

$$t = \ln(k/\Lambda)$$

[C.Wetterich, *Phys. Lett. B*, 301, 1993]



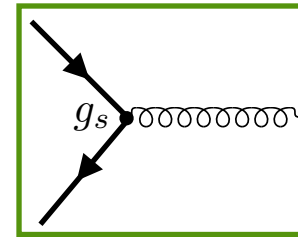
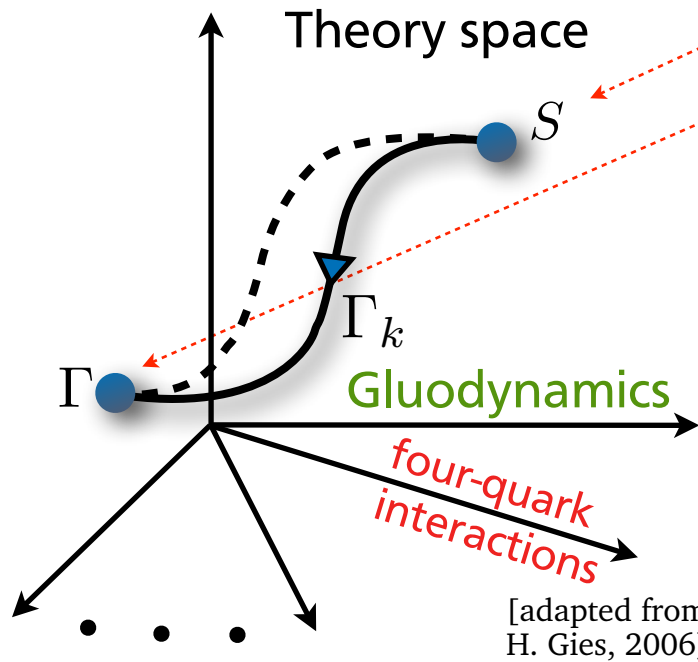
Functional renormalization group (FRG)

Flow from high to low energies in QCD

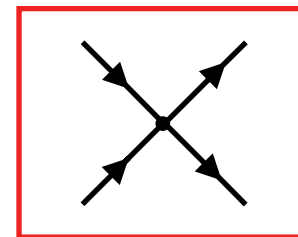
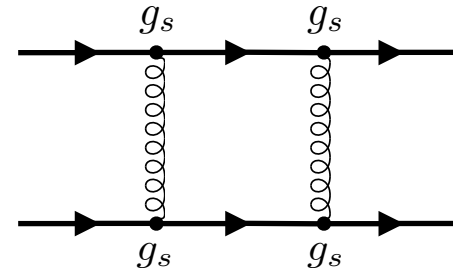
$\Gamma \xrightarrow{R_k}$ Effective average action Γ_k

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$

IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$



Everything that is not forbidden is allowed.



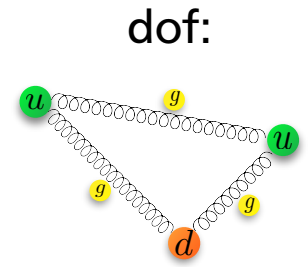
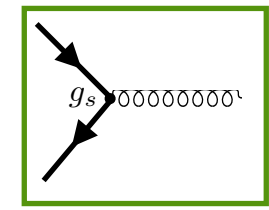
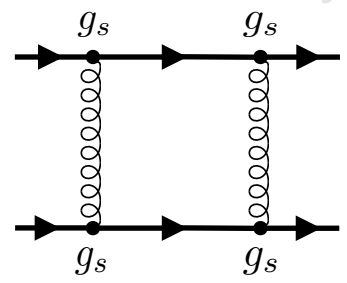
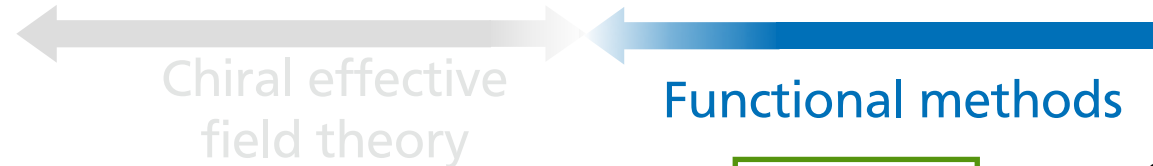
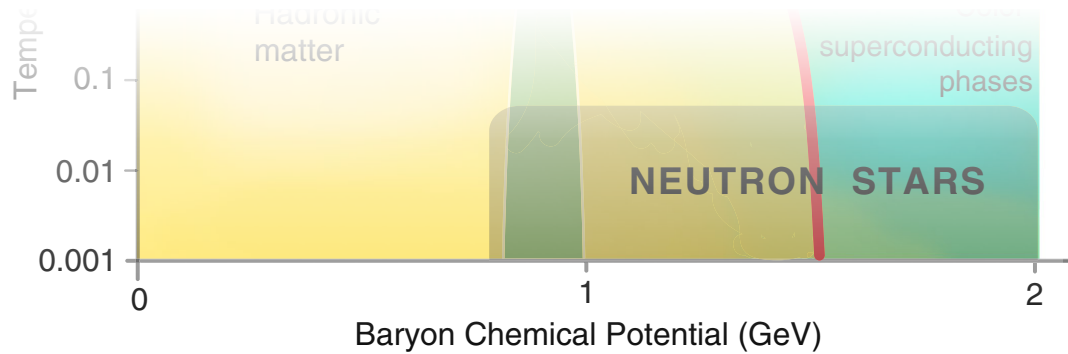
Flow equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

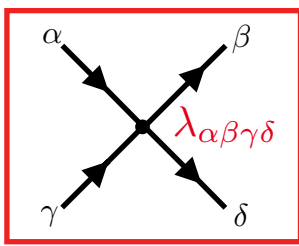
$$t = \ln(k/\Lambda)$$

[C.Wetterich, *Phys. Lett. B*, 301, 1993]

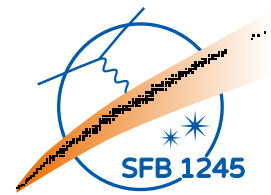
QCD phase diagram: Neutron stars and the cold dense EoS



Four-quark interactions



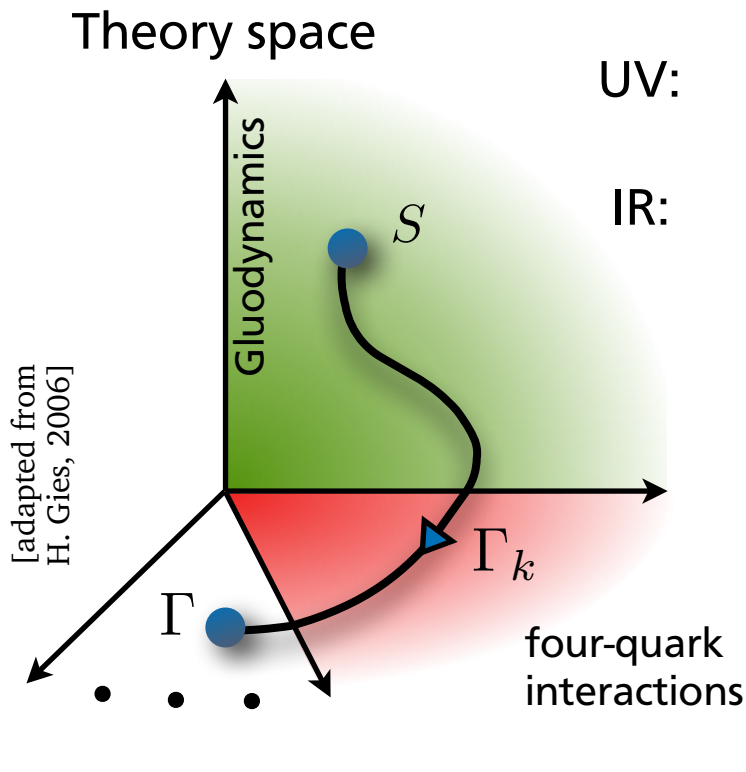
1. Important effective interactions
2. Encode information on the ground state and SB
3. Full treatment possible and crucial



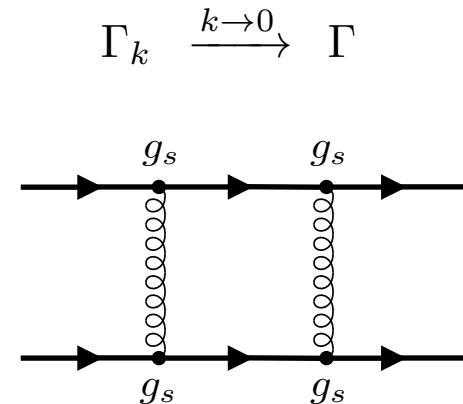
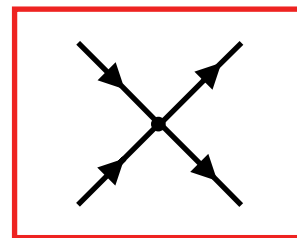
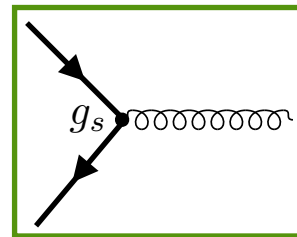
Functional renormalization group (FRG)

Flow from high to low energies in QCD

Effective average action Γ_k



$$\Gamma_{k \rightarrow \Lambda} = \int_x \left\{ \frac{1}{4} F_{\mu\nu} F_{\mu\nu} + \bar{\psi} (i \not{\partial} + i g_s \not{A}) \psi \right\}$$



Functional renormalization group (FRG)

Flow from high to low energies in QCD



TECHNISCHE
UNIVERSITÄT
DARMSTADT

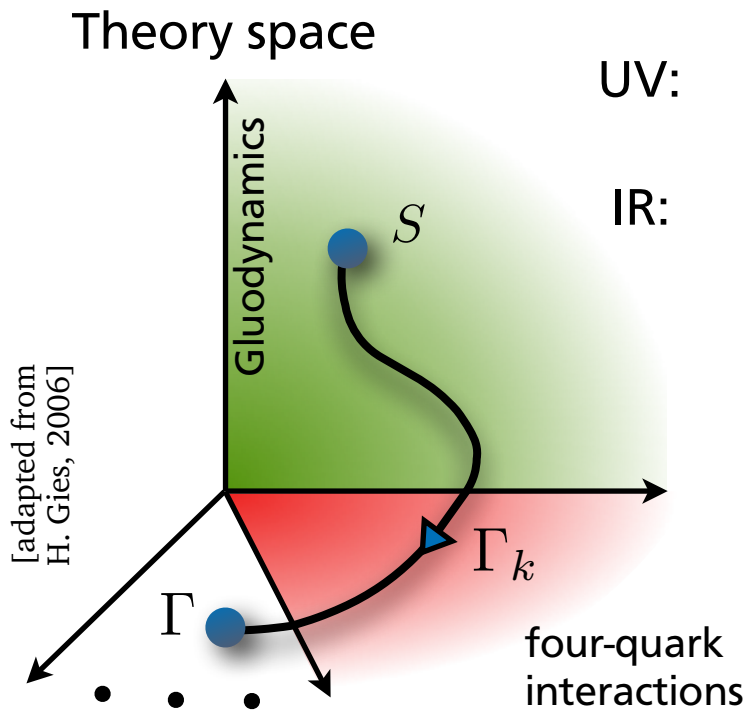
Effective average action Γ_k

RG flow equation

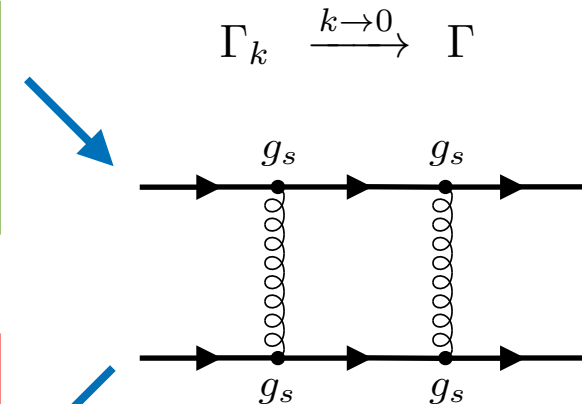
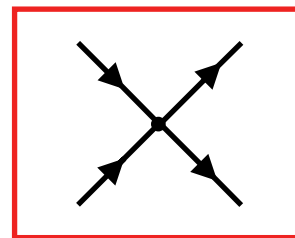
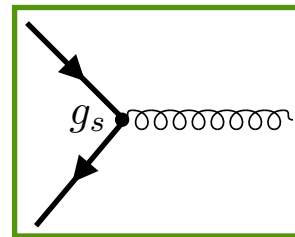
[C. Wetterich, 1993]

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

$$t = \ln(k/\Lambda)$$



$$\Gamma_{k \rightarrow \Lambda} = \int_x \left\{ \frac{1}{4} F_{\mu\nu} F_{\mu\nu} + \bar{\psi} (i \not{\partial} + i g_s \not{A}) \psi \right\}$$



Functional renormalization group (FRG)

Partition function/
generating functional

$$Z = \text{tr} \left[e^{-\beta \hat{H}} \right] = \int \mathcal{D}\varphi e^{-S[\varphi]}$$

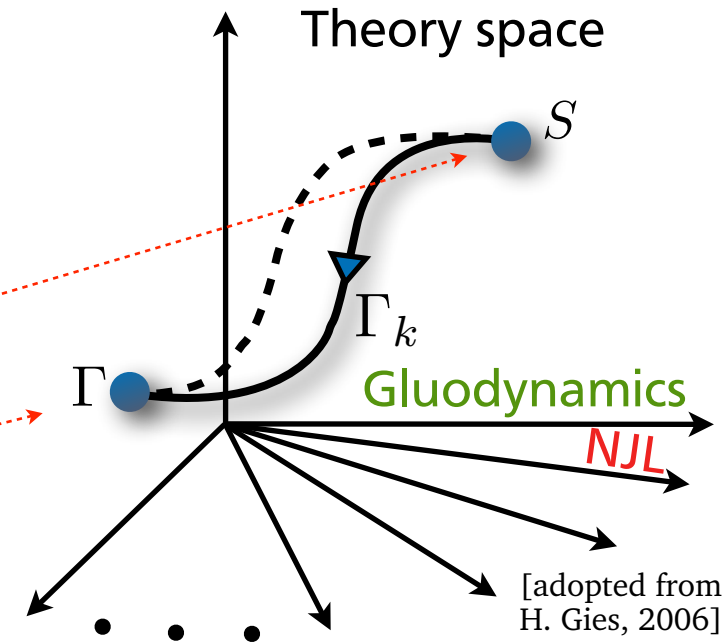
$$\beta = 1/T$$

→ Effective action Γ

→ R_k Effective average action Γ_k

UV: $\Gamma_k \xrightarrow{k \rightarrow \Lambda} S$

IR: $\Gamma_k \xrightarrow{k \rightarrow 0} \Gamma$



Flow
equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ \left[\Gamma_k^{(2)} + R_k \right]^{-1} \cdot (\partial_t R_k) \right\}$$

$$t = \ln(k/\Lambda)$$

[C.Wetterich, *Phys. Lett. B*, 301, 1993]