



***nn* scattering length
from the
 ${}^6\text{He}(p, p\alpha)nn$ reaction**

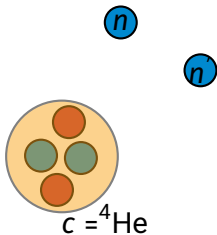
Matthias Göbel
Institut für Kernphysik
Technische Universität Darmstadt

March 24, 2021





Theory for obtaining the
nn scattering length
from the
 ${}^6\text{He}(p, p\alpha)nn$ reaction





project A05: Halos and clustering in nuclei

- explore cluster degrees of freedom of exotic nuclei & few-neutron systems
- use reactions with radioactive beams and (halo) effective field theory



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nn scattering length from ${}^6\text{He}(p, p\alpha)nn$

- exploit the experimental and theoretical tools in this field to measure the nn scattering length
- experiment: reaction in inverse kinematics
 - will be conducted by Aumann & SAMURAI collaboration @ RIKEN RIBF
 - Marco Knösel works on detector design & data analysis
- theory: describe reaction in Halo EFT



Motivation & approach

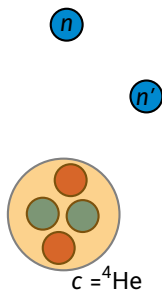
Halo EFT for ${}^6\text{He}$

Benchmarking the ground-state spectrum

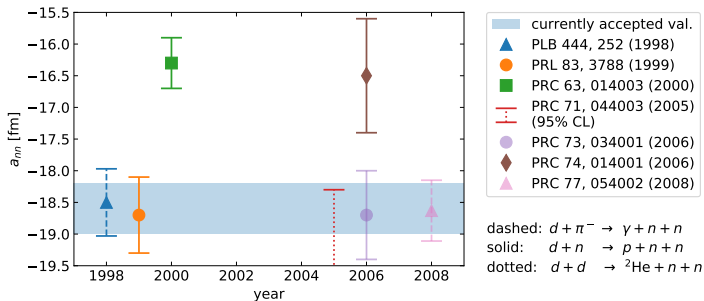
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Results for the E_{nn} spectrum

Conclusion & Outlook

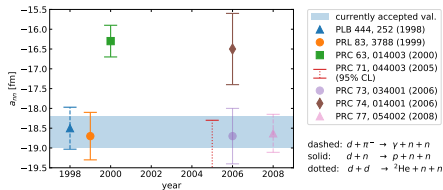


- **motivation:** no high-precision value for nn scattering length available





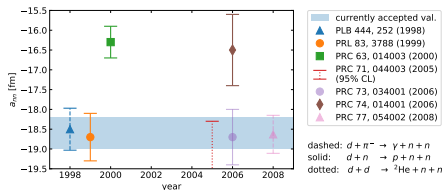
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- use the reaction ${}^6\text{He}(p, p\alpha)nn$ to determine the scattering length from the final E_{nn} spectrum

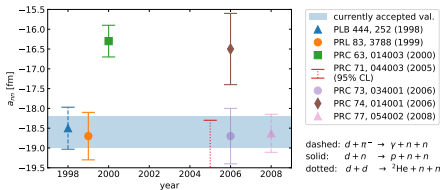


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- use the reaction ${}^6\text{He}(p, p\alpha)nn$ to determine the scattering length from the final E_{nn} spectrum
- advantages of this approach
 - different from the previous methods → not the same difficulties
 - final nn pair has high center-of-mass velocity in the lab system → avoids problems with detection efficiency

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→ use the reaction ${}^6\text{He}(p, p\alpha)nn$ to determine the scattering length from the final E_{nn} spectrum

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- experiment proposal from Aumann & SAMURAI collaboration approved by RIKEN RIBF [NP2012-SAMURAI55R1 \(2020\)](#)



■ **approach:** ${}^6\text{He}$ in Halo EFT

1. calculate wave function $\Psi_c(p, q)$ (& do comparisons with model calc.)
2. take final state interaction (FSI) into account
3. calculate the probability distribution for E_{nn}

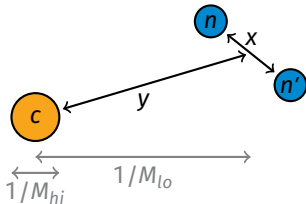


■ approach: ${}^6\text{He}$ in Halo EFT

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■ tool: Halo EFT

- $\not\equiv$ EFT
- core & valence nucleons as degrees of freedom
- results are expanded in k/M_{hi}
→ systematic improvement possible





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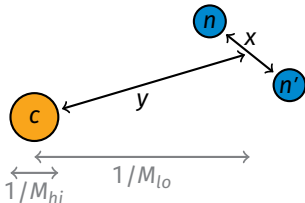
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■ properties of ${}^6\text{He}$

- Borromean $2n$ halo
- separation of scales: $S_{2n} = 0.975 \text{ MeV} < E_\alpha^* \approx 20 \text{ MeV}$
- quantum numbers: $J^\pi = 0^+$ (${}^4\text{He}$: $J^\pi = 0^+$)
- leading-order (LO) Halo EFT interaction channels:
 - nn : 1S_0
 - nc : ${}^2P_{3/2}$ (not at LO: ${}^2P_{1/2}, {}^2S_{1/2}$)



Halo EFT for ${}^6\text{He}$ formulated in Ji, Elster, Phillips, PRC 90 (2014)
review of Halo EFT in Hammer, Ji, Phillips, JPG 44 (2017)





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Halo EFT for ${}^6\text{He}$

Benchmarking the ground-state spectrum

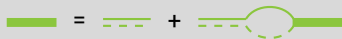
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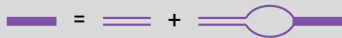
use EFT in dimer formalism

1. step: obtain dressed dimer propagators



& renormalize using **input values**

a_1, r_1



a_0

2. step: set up equations for Faddeev transition amplitudes



obtain wave functions by multiplying amplitudes with propagators

$$\Psi_n(p, q) = \begin{array}{c} p \\ | \\ q \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \mathcal{A}_n + \begin{array}{c} p \\ | \\ q \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \mathcal{A}_n + \begin{array}{c} p \\ | \\ q \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \mathcal{A}_c$$

$$\Psi_c(p, q) = \begin{array}{c} p \\ | \\ q \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \mathcal{A}_c + 2 \times \begin{array}{c} p \\ | \\ q \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \mathcal{A}_n$$

The diagrams illustrate the calculation of wave functions by multiplying amplitudes with propagators. The top row shows the calculation of $\Psi_n(p, q)$ as a sum of three terms. Each term consists of a vertical bracket on the left labeled with p at the top and q at the bottom, followed by two horizontal lines representing particle paths. The first term shows a green oval labeled \mathcal{A}_n with a green segment on the top path and a dashed orange segment on the bottom path. The second term shows a green oval labeled \mathcal{A}_n with a blue segment on the top path and a dashed orange segment on the bottom path. The third term shows a blue oval labeled \mathcal{A}_c with a blue segment on the top path and a dashed orange segment on the bottom path. The bottom row shows the calculation of $\Psi_c(p, q)$ as a sum of two terms. The first term shows a blue oval labeled \mathcal{A}_c with a purple segment on the top path and a dashed orange segment on the bottom path. The second term shows a green oval labeled \mathcal{A}_n with a green segment on the top path and a dashed orange segment on the bottom path, preceded by a coefficient of $+2 \times$.



obtain wave functions by multiplying amplitudes with propagators

$$\Psi_n(p, q) = \begin{array}{c} \text{diagram 1} \\ \text{diagram 2} \\ \text{diagram 3} \end{array} + \begin{array}{c} \text{diagram 4} \\ \text{diagram 5} \\ \text{diagram 6} \end{array} + \begin{array}{c} \text{diagram 7} \\ \text{diagram 8} \\ \text{diagram 9} \end{array}$$
$$\Psi_c(p, q) = \begin{array}{c} \text{diagram 10} \\ \text{diagram 11} \\ \text{diagram 12} \end{array} + 2 \times \begin{array}{c} \text{diagram 13} \\ \text{diagram 14} \\ \text{diagram 15} \end{array}$$

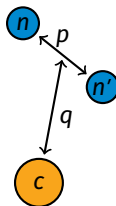
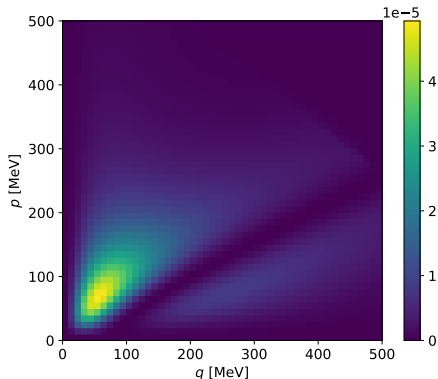
The diagrams illustrate the calculation of wave functions $\Psi_n(p, q)$ and $\Psi_c(p, q)$ using propagators \mathcal{A}_n and \mathcal{A}_c . Each diagram shows a set of lines representing particles, with vertical brackets on the left labeled p and q . The diagrams are arranged in two rows. The first row shows three diagrams for $\Psi_n(p, q)$, and the second row shows two diagrams for $\Psi_c(p, q)$. The diagrams use different colors (green, blue, purple) to represent different components or states.

→ use wave function with core as spectator (Ψ_c)

calculated ground-state wave functions and probability densities in Halo EFT

Göbel, Hammer, Ji, Phillips, FBS 60 (2019)

$$\Psi_c^2(p, q) p^2 q^2$$





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Compare ground-state distribution from EFT with model calculations



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- compare ground-state distributions $\rho(p_{nn}) \approx \int dq q^2 |\Psi_c(p_{nn}, q)|^2$
- \nexists published results for $\rho(p_{nn})$
- use FaCE [Thompson, Nunes, Danilin, Comput.Phys.Commun. 161 \(2004\)](#)



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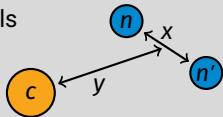


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computer code: Faddeev with Core Excitations (FaCE)

- solves the Schrödinger equation of three-body cluster models
- *input*:
 - local l -dependent two-body potentials (central or spin-orbit)
 - phenomenological three-body force

- *output*: hyperspherical wave function components $\chi_{K,l}^S(\rho)$ with $\rho^2 = x^2 + y^2$



Compare ground-state distribution from EFT with model calculations: defining the model



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the two-body potentials in use

use local, l -dependent **Gaussian** potentials

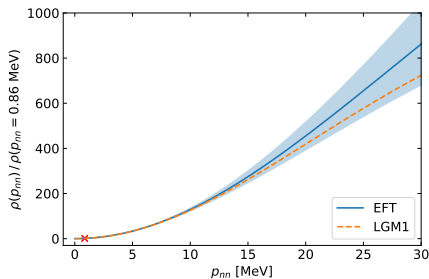
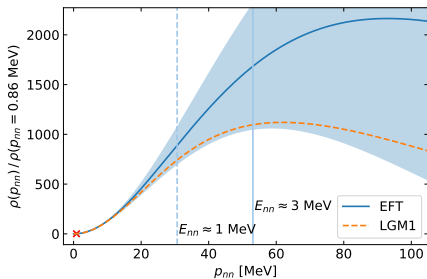
- central pot.: $\langle r; l, s | V_c^{(\tilde{l})} | r'; l', s' \rangle := \delta_{l,l'} \delta_{l,\tilde{l}} \frac{\delta(r'-r)}{r'^2} \bar{V}_c^{(l)} \exp(-r^2 / (a_{c;l}^2))$
- spin-orbit pot.

"standard setting": local Gaussian model 1 (LGM1)

- nn interaction: $V_c^{(0)}$
- nc interaction: $V_c^{(0)}, V_c^{(1)}, V_{SO}^{(1)}, V_c^{(2)}, V_{SO}^{(2)}$
- phenomenological three-body force



Compare ground-state distribution from EFT with model calculations: Results I



Göbel, Aumann, Bertulani, Frederico, Hammer, Phillips, arXiv:2103.03224 (2021)

- model calc. within uncertainty band of EFT ✓
- especially for small p_{nn} agreement is good
- investigate sources of discrepancies by doing additional model calculations

Compare ground-state distribution from EFT with model calculations: Results II



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additional model calculations

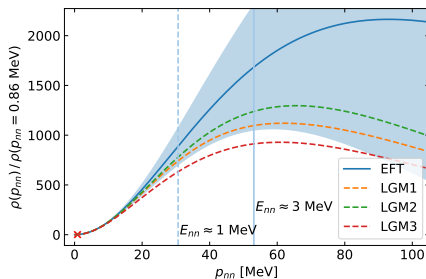
- LGM2: same as LGM1, but nc int. only in ${}^2P_{3/2}$
- LGM3: same as LGM2, but no three-body force

Compare ground-state distribution from EFT with model calculations: Results II

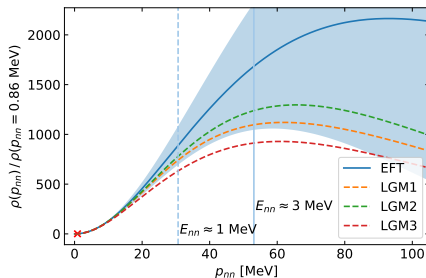


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Compare ground-state distribution from EFT with model calculations: Results II



conclusions so far

- additional nc int. channels are important
- three-body force important for obtaining the correct S_{2n}
- range of three-body force with tuned strength not so important

Compare ground-state distribution from EFT with model calculations: Results III



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additional model calculations II

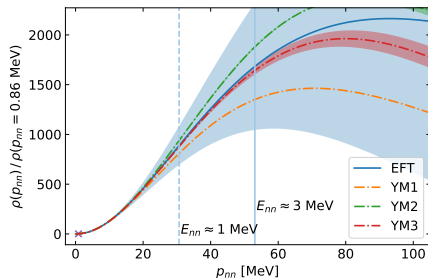
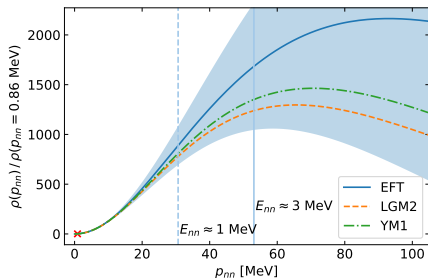
- turn three-body Halo EFT calc. into model calc. by using different t
- use t with Yamaguchi form factors → good phase shifts



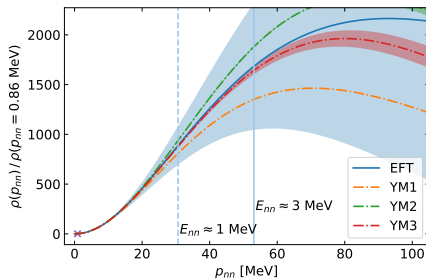
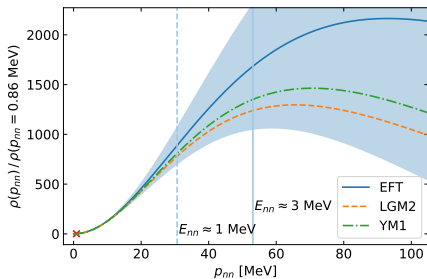
Compare ground-state distribution from EFT with model calculations: Results III

additional model calculations II

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Compare ground-state distribution from EFT with model calculations: Results III



conclusions

- *nc* int. accounts for most of the discrepancies between LO EFT and model calc.
- especially the unitarity term is important



Motivation & approach

Halo EFT for ${}^6\text{He}$

Benchmarking the ground-state spectrum

Reaction theory for ${}^6\text{He}(p, p\alpha)nn$

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Conclusion & Outlook





- ${}^6\text{He}(p, p\alpha)nn$ is a knock-out reaction, in which the α core of ${}^6\text{He}$ is removed by a p
- **initial state:** ${}^6\text{He}$ bound state $|\Psi\rangle$
$$(K_{nn} + K_{c(nn)} + V_{nn} + V_{nc} + V_{3B}) |\Psi\rangle = -B_3 |\Psi\rangle$$
- **final state:** $|p, q\rangle_c$ all particles are free (state of definite momentum!)
$$(K_{nn} + K_{c(nn)}) |p, q\rangle_c = (-B_3 + E_{KO}) |p, q\rangle_c$$



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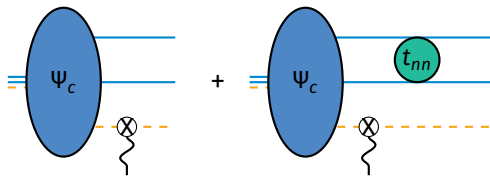


- treatments of FSI are based on two-potential scattering theory
Goldberger, Watson, "Collision Theory" (1964)
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 - exact calculation
- exact calculation is based on t_{nn}





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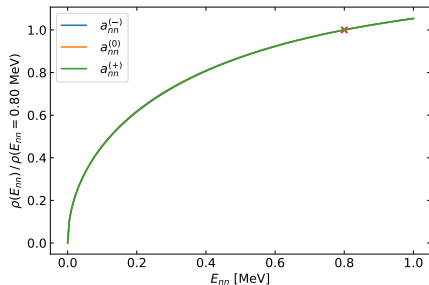
Conclusion & Outlook





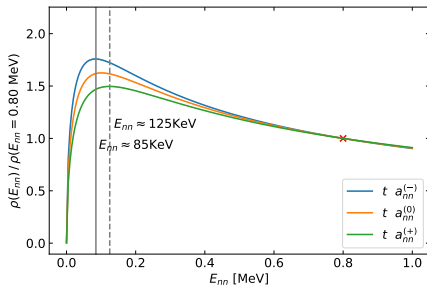
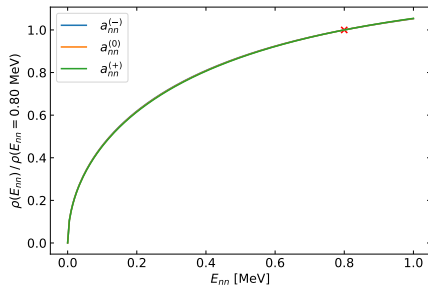
obtain distribution by using $\rho^{(t)}(p) = \int dq p^2 q^2 \left| \Psi_c^{(wFSI)}(p, q) \right|^2$

variation of a_{nn} : $a_{nn}^{(-)} = -20.7$ fm, $a_{nn}^{(0)} = -18.7$ fm, $a_{nn}^{(+)} = -16.7$ fm



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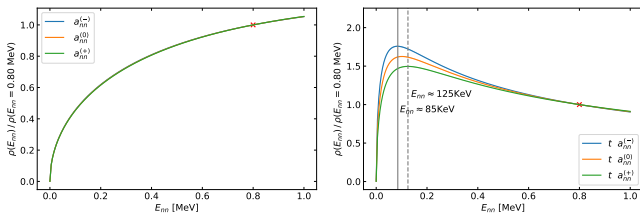
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Göbel, Aumann, Bertulani, Frederico, Hammer, Phillips, arXiv:2103.03224 (2021)

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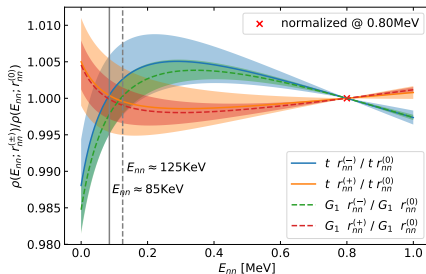


conclusions

- significant sensitivity on nn scattering length
- sensitivity almost entirely caused by FSI \rightarrow ${}^6\text{He}$ is simply a suitable neutron source (nevertheless, ${}^6\text{He}$ wave function is an important ingredient)

Sensitivity of the E_{nn} spectrum on effective range

variation of r_{nn} : $r_{nn}^{(-)} = 2.0$ fm, $r_{nn}^{(0)} = 2.73$ fm, $r_{nn}^{(+)} = 3.0$ fm



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- almost no sensitivity on nn effective range



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- go to NLO





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