

# Perspectives on SFB EOS theory

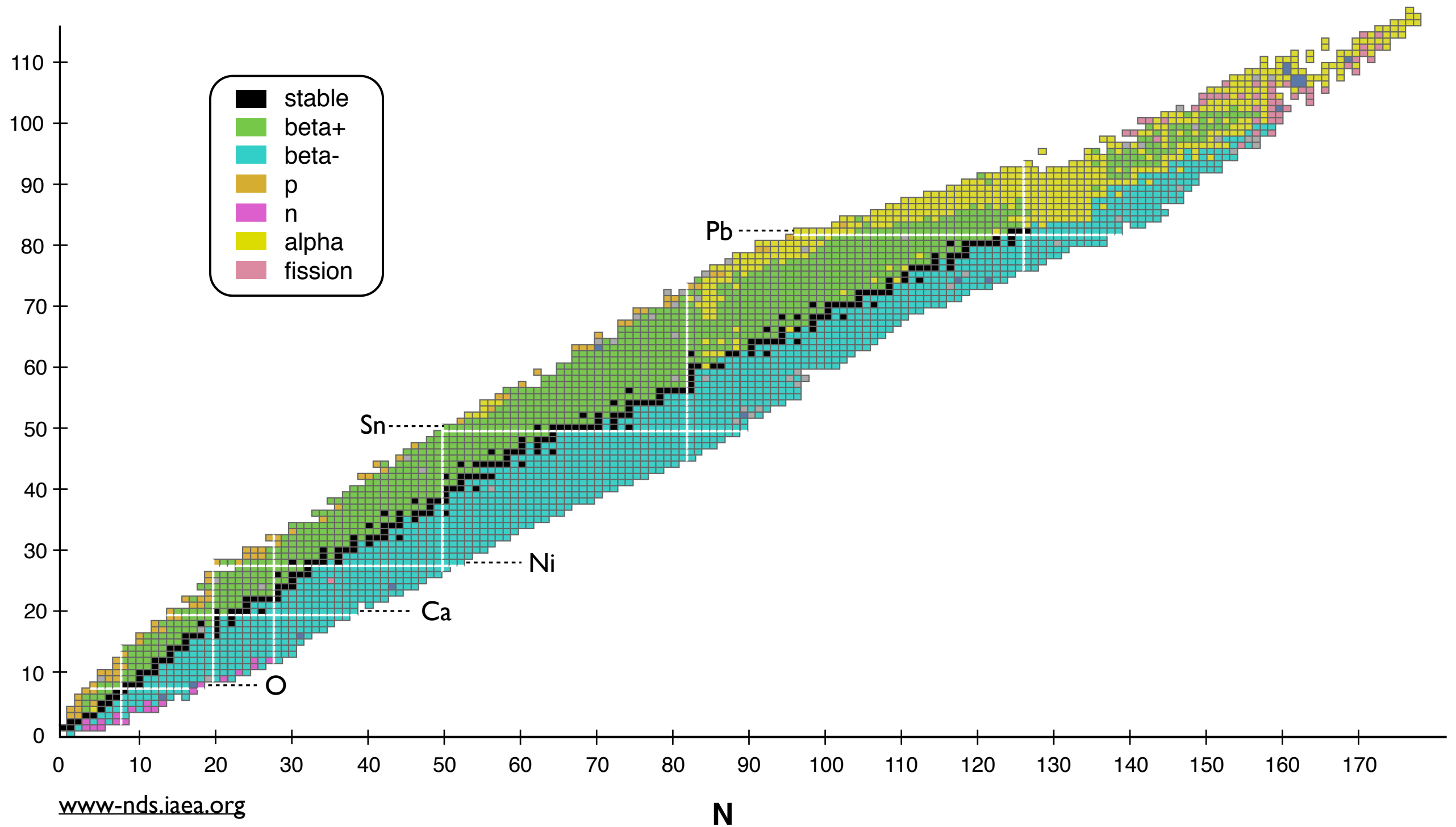


**3rd workshop of the SFB 1245**

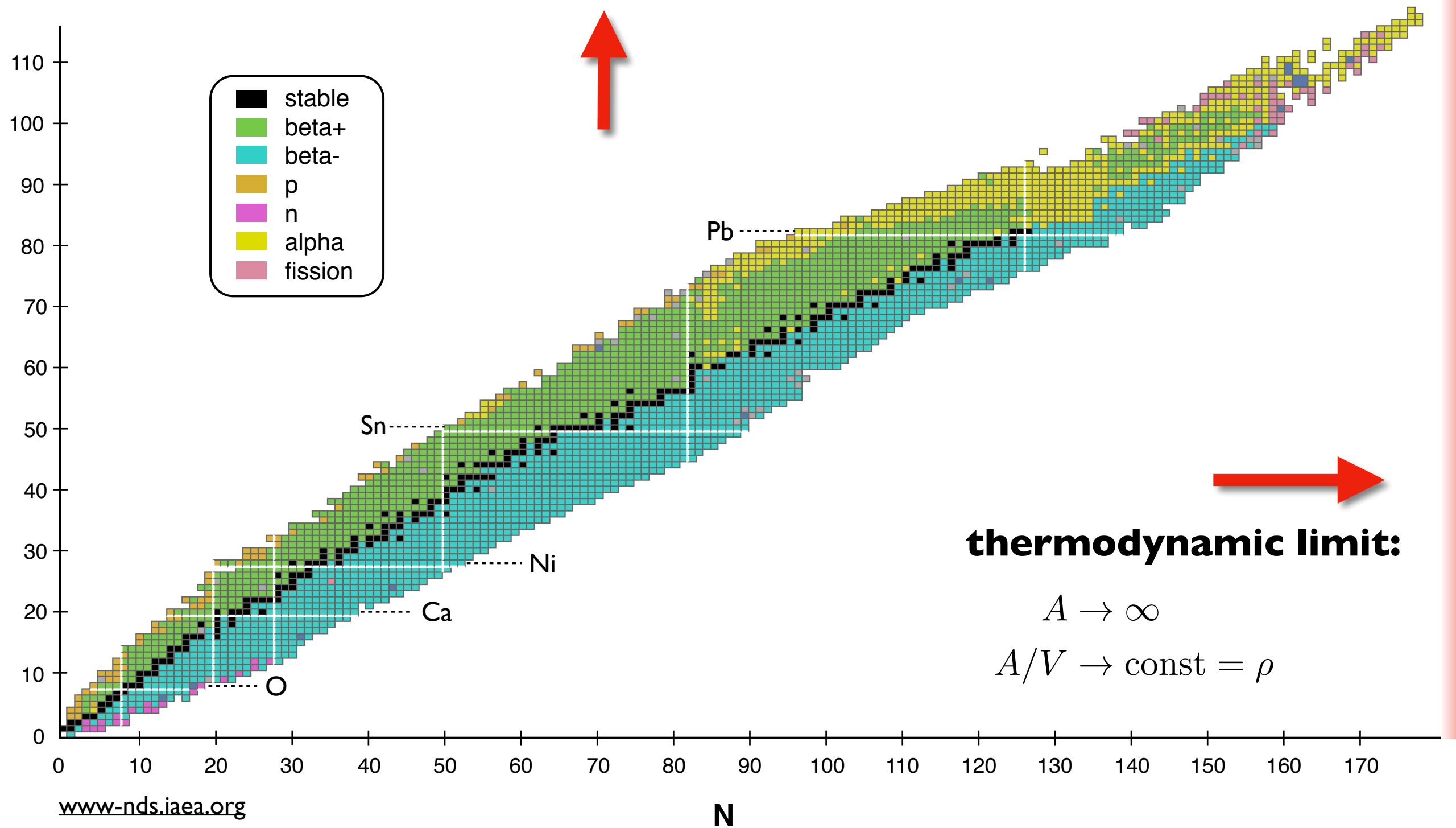
Kai Hebel

Mainz, July 5, 2018

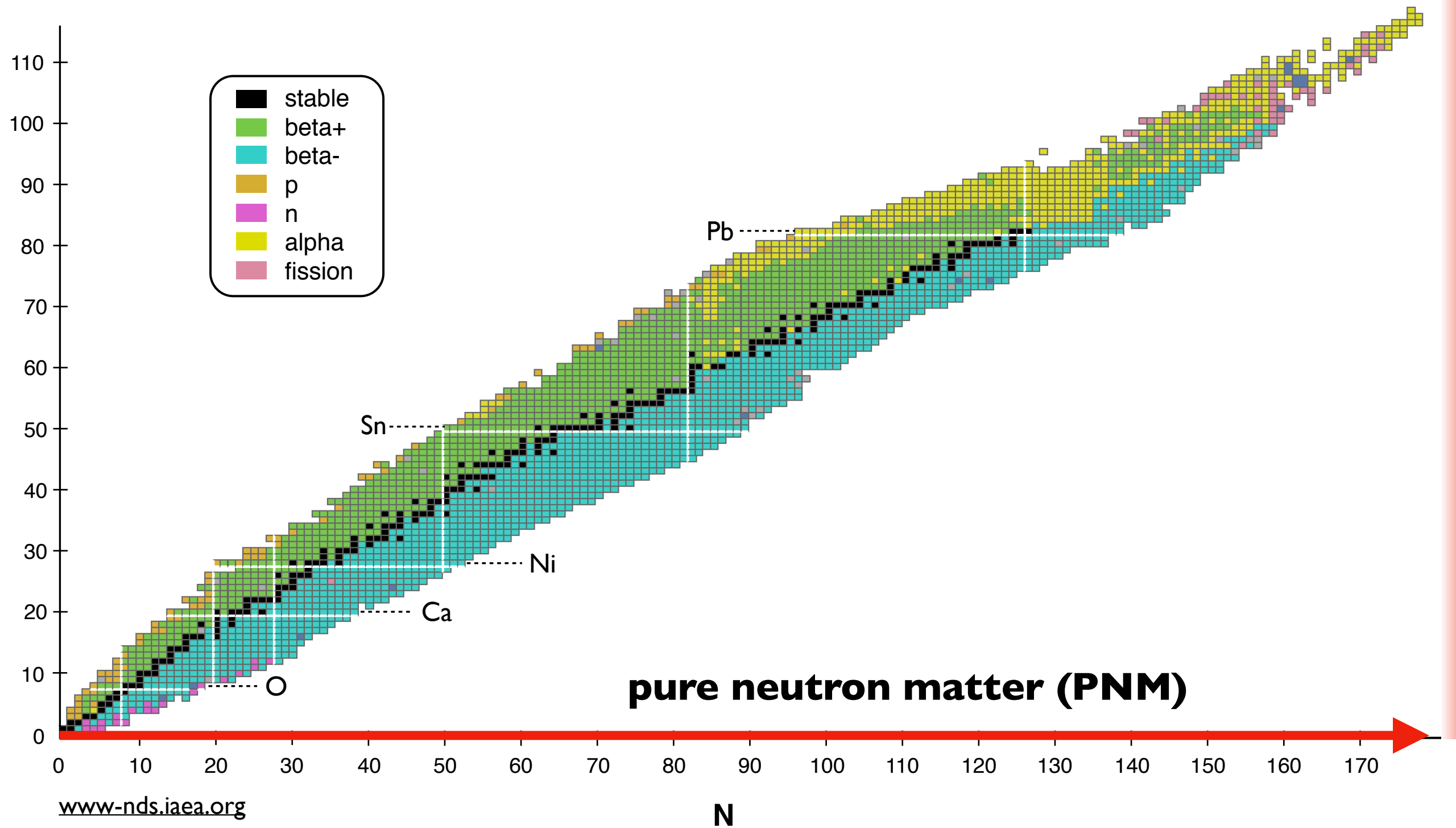
# The nuclear landscape and nuclear matter



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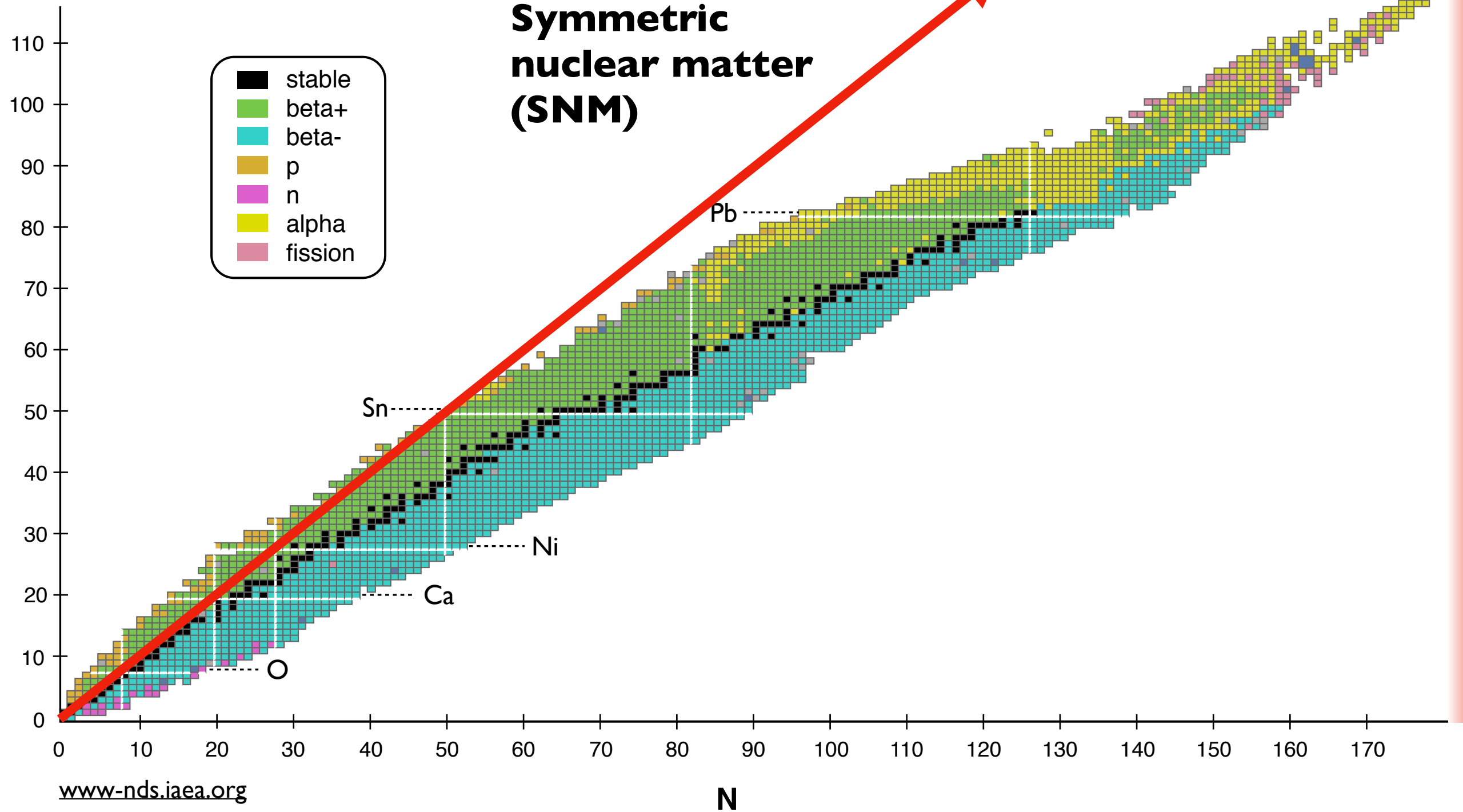


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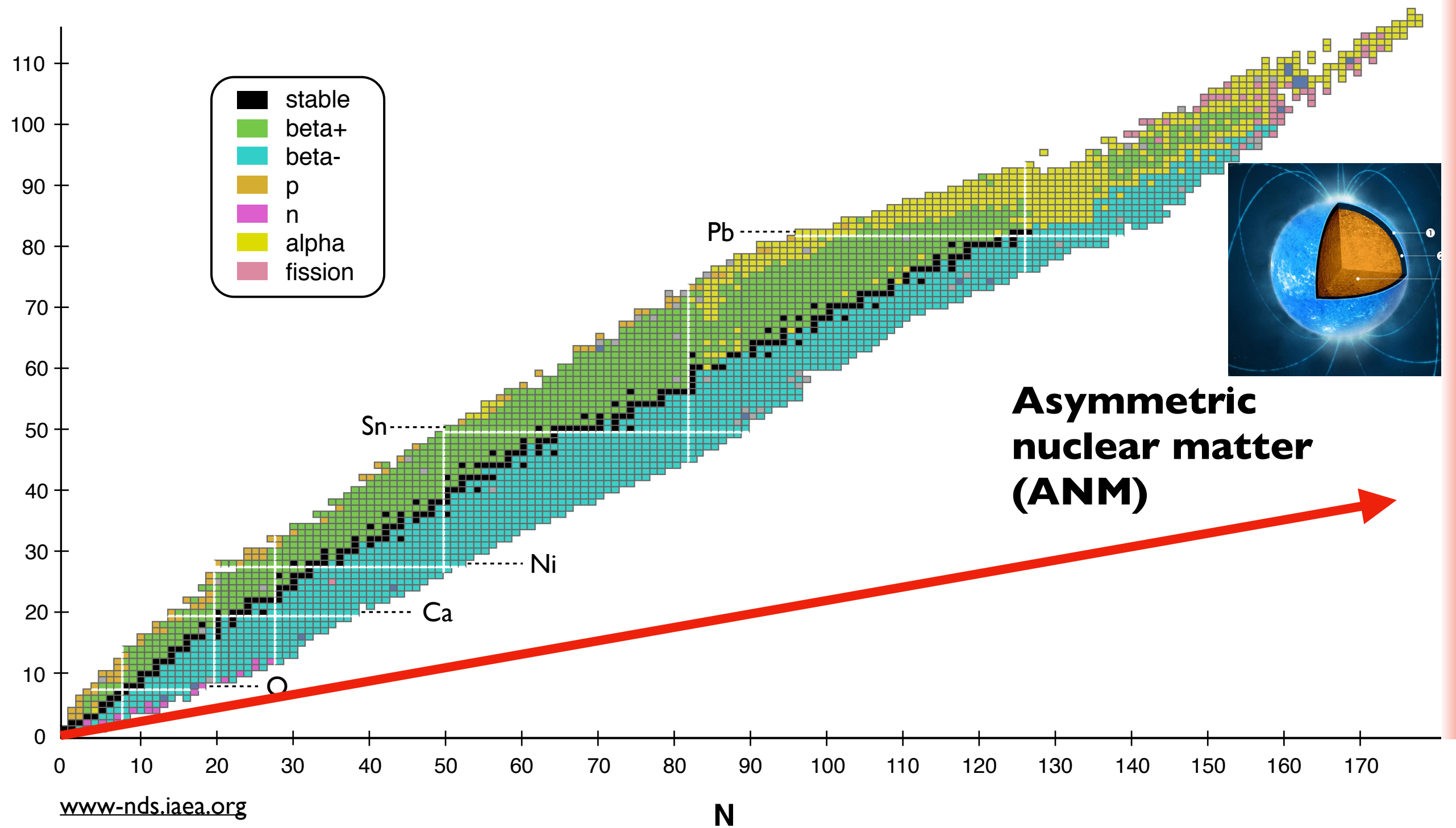


# The nuclear landscape and nuclear matter

## Symmetric nuclear matter (SNM)

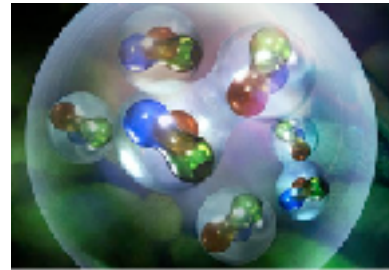


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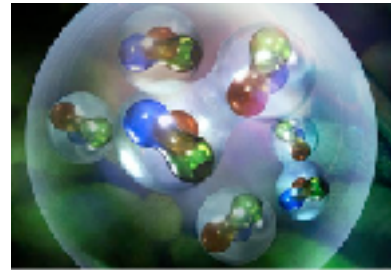
# From the strong interaction to nuclear matter

nuclear interactions



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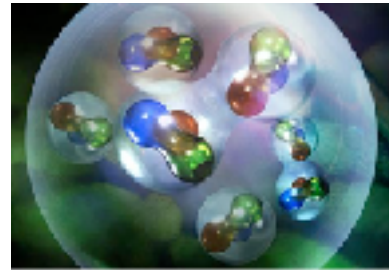


many-body frameworks



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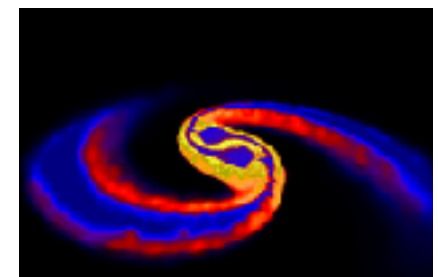
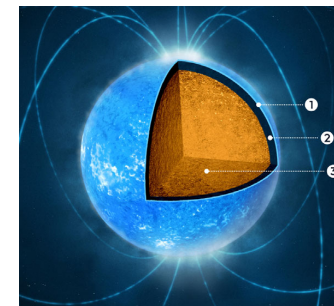


	00 Basis	01 Basis	02 Basis
$\langle 0   \sigma(\vec{r})$	X H	-	-
$\langle 1   \sigma(\vec{r})$	X H	-	-
$\langle 2   \sigma(\vec{r})$	H H	X X	-
$\langle 3   \sigma(\vec{r})$	X H	H H	X H

many-body frameworks

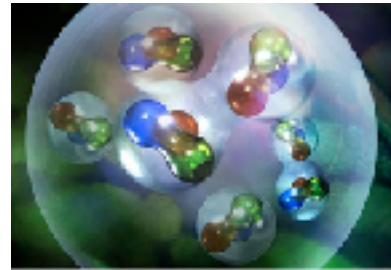
equation of state

$$E = E(\rho, T, Y_e, \dots)$$
$$P = P(\rho, T, Y_e, \dots)$$



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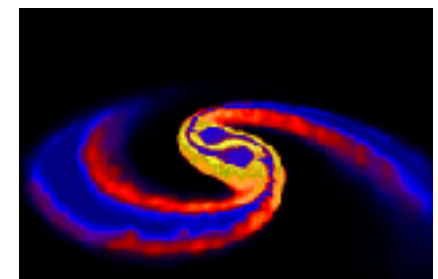
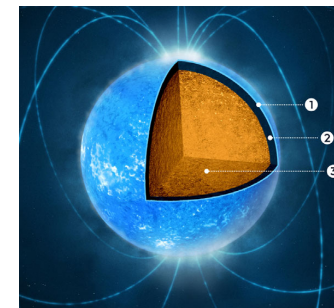
	00 Basis	01 Basis	02 Basis
$\chi_0 \in \mathcal{H}$	X H	-	-
$\chi_{10} \in \mathcal{H}$	X H	-	-
$\chi_{11} \in \mathcal{H}$	X H	X H	-
$\chi_{20} \in \mathcal{H}$	X H	X H	X H

many-body frameworks

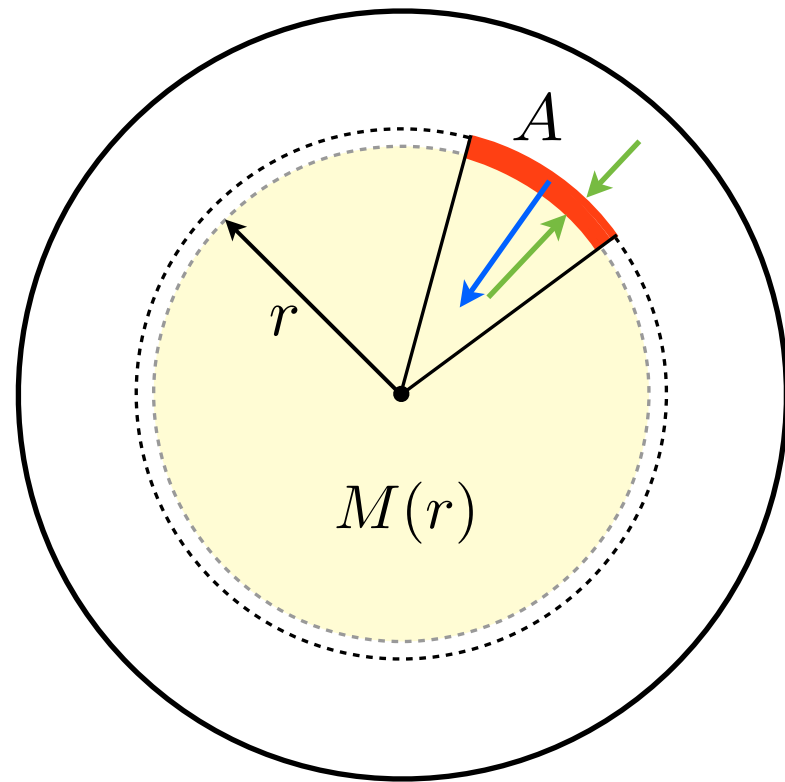
probing  
next-generation  
nuclear forces

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# The EOS of high-density matter: neutron stars

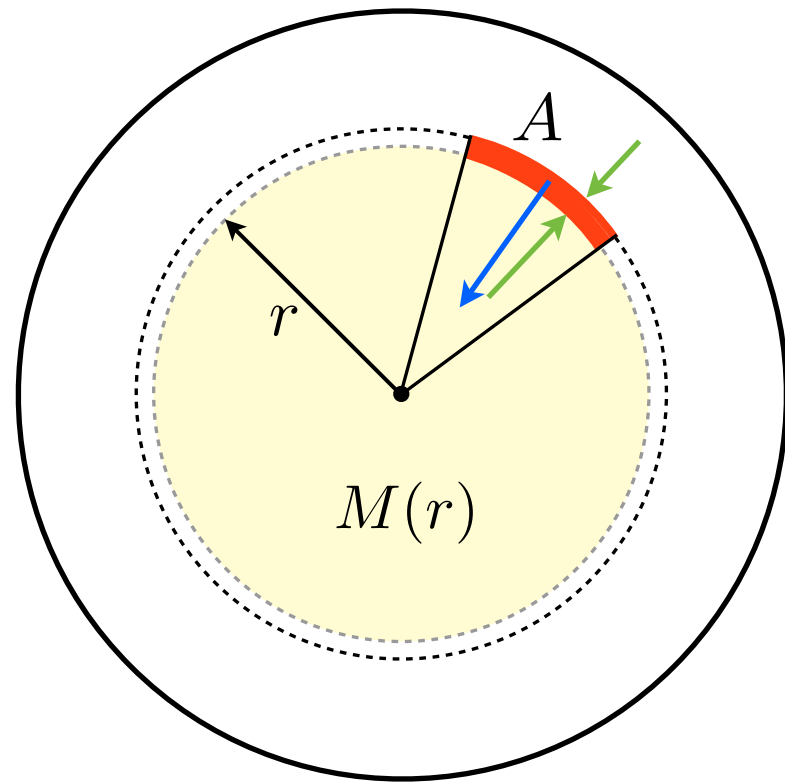


- consider forces on a mass element:

gravity: 
$$F_g = -\frac{GM(r)\rho(r)A dr}{r^2}$$

pressure difference: 
$$F_p = A (p_{\text{out}} - p_{\text{in}}) = A dp$$

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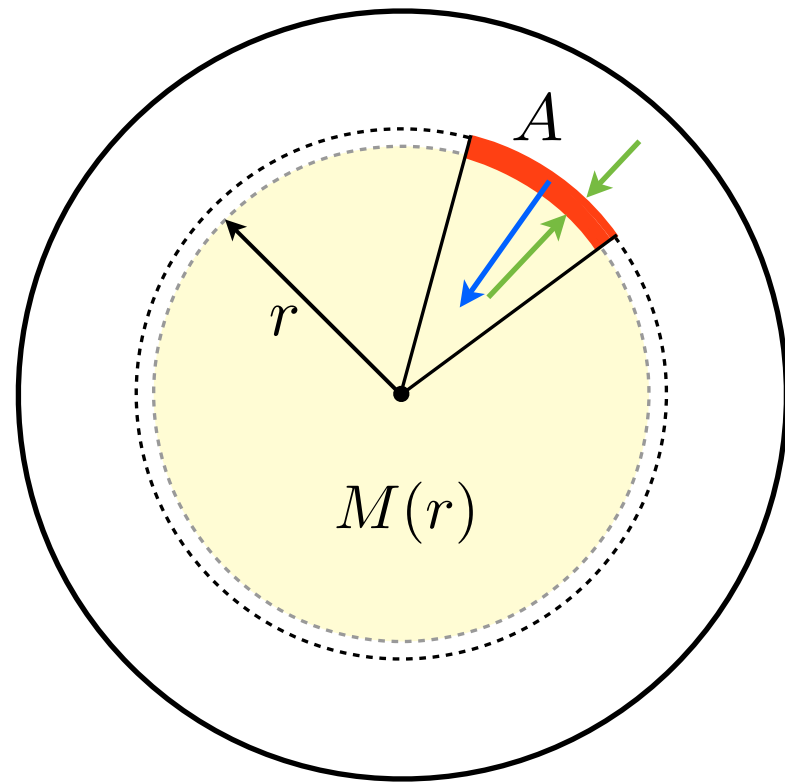
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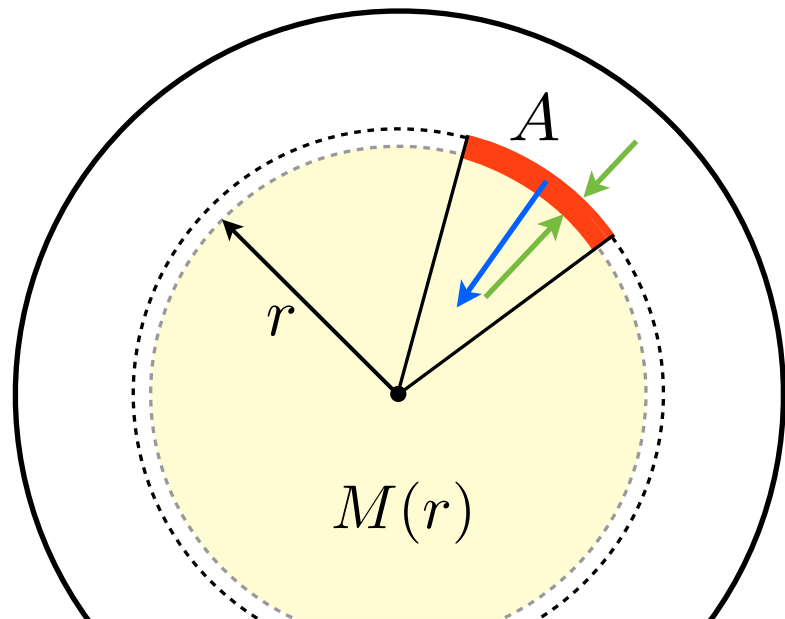
$$F_g = F_p \Rightarrow \frac{dp}{dr} = -\frac{GM(r)\rho(r)}{r^2}$$

- include general-relativistic corrections:

$$\frac{dp}{dr} = -\frac{GM(r)\varepsilon(r)}{r^2} \left[ 1 + \frac{p(r)}{\varepsilon(r)c^2} \right] \left[ 1 + \frac{4\pi r^3 p(r)}{M(r)c^2} \right] \left[ 1 - \frac{2GM(r)}{c^2 r} \right]$$

*'Tolman-Oppenheimer-Volkov'* equation

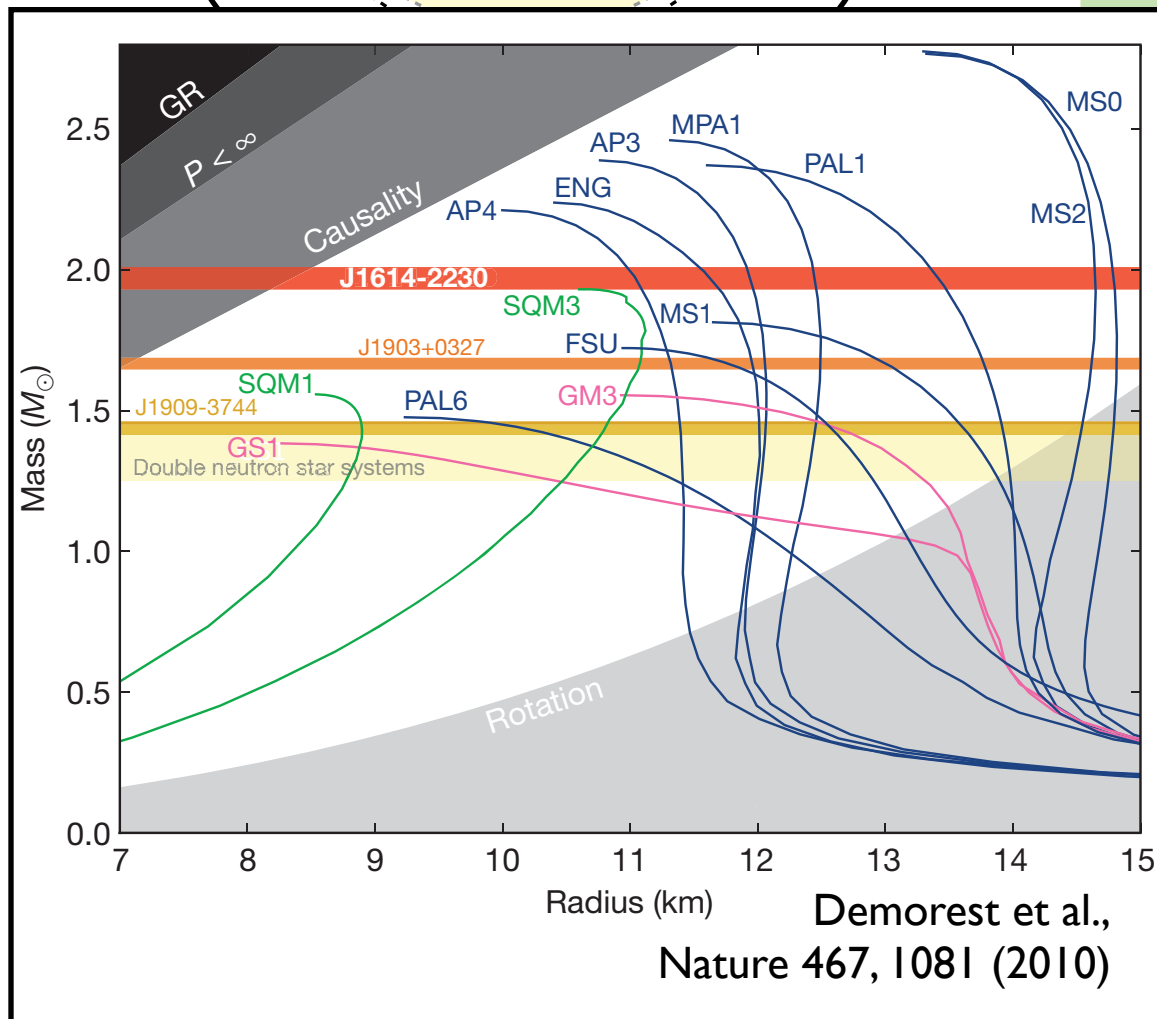
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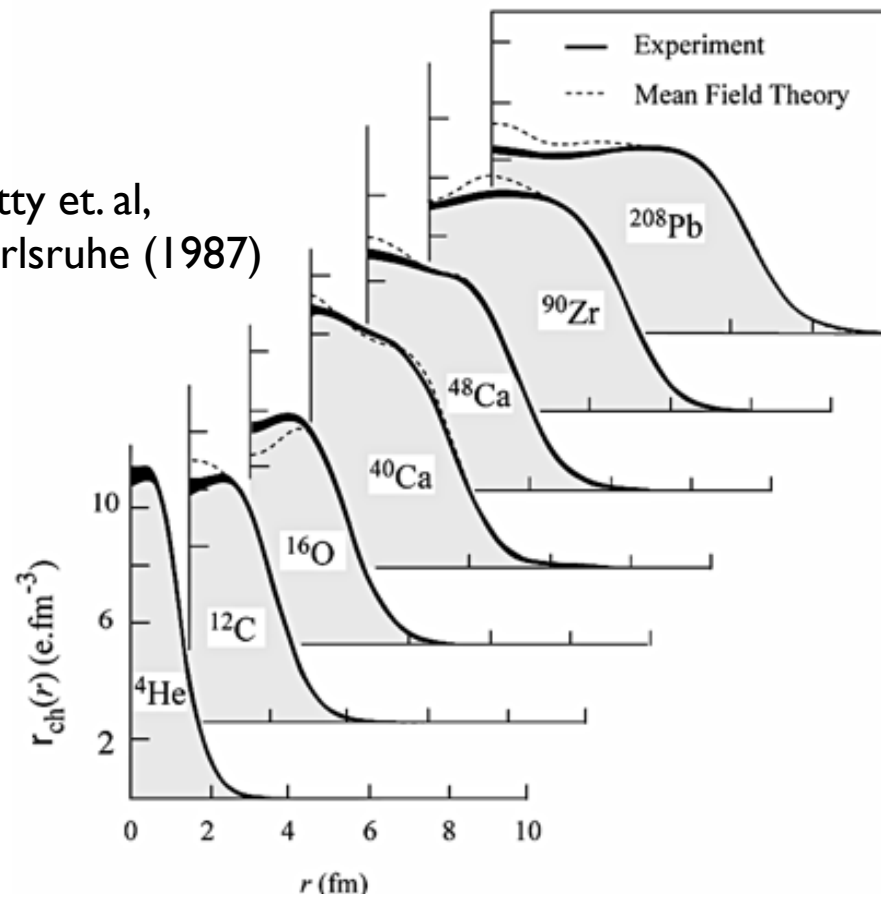
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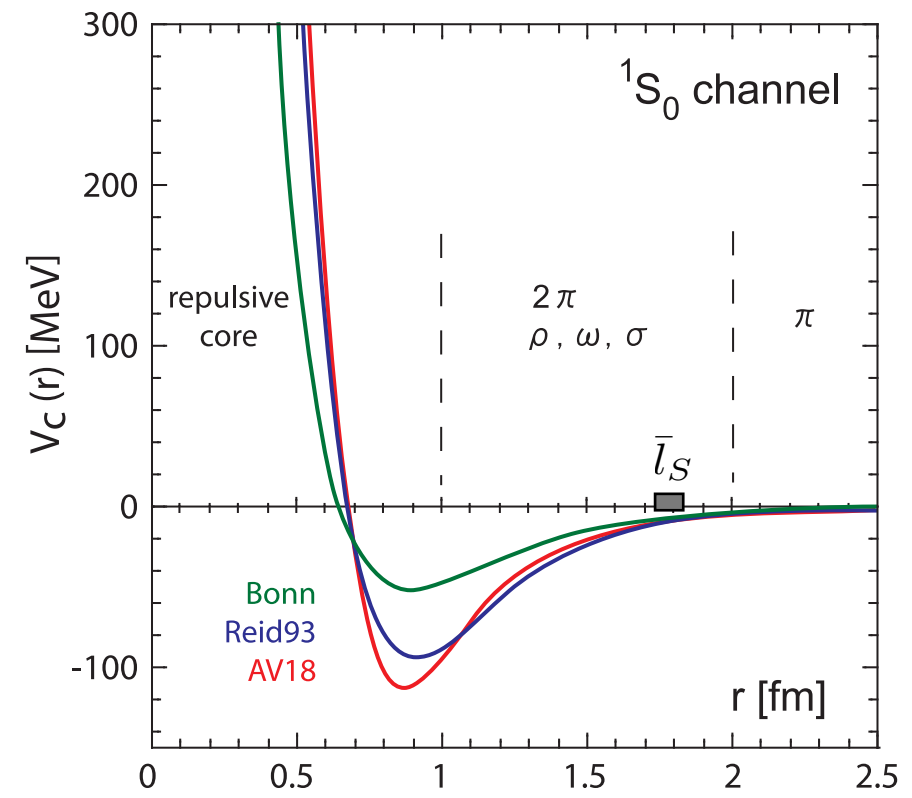
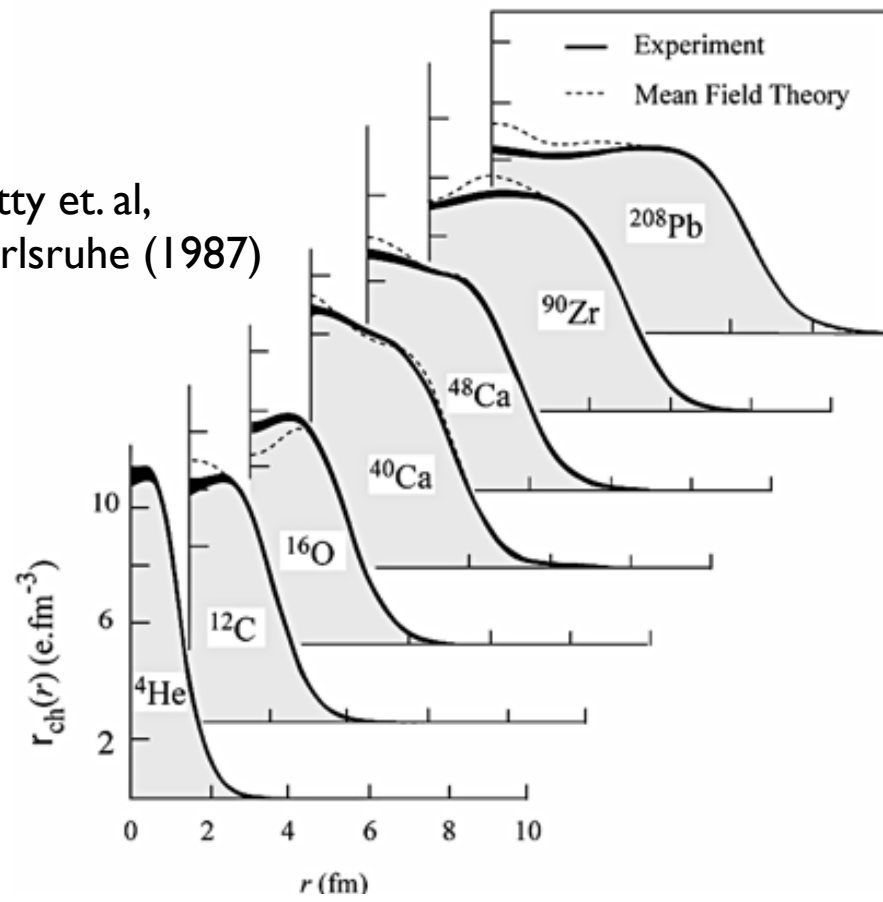
# Nuclear saturation

Batty et. al,  
Karlsruhe (1987)



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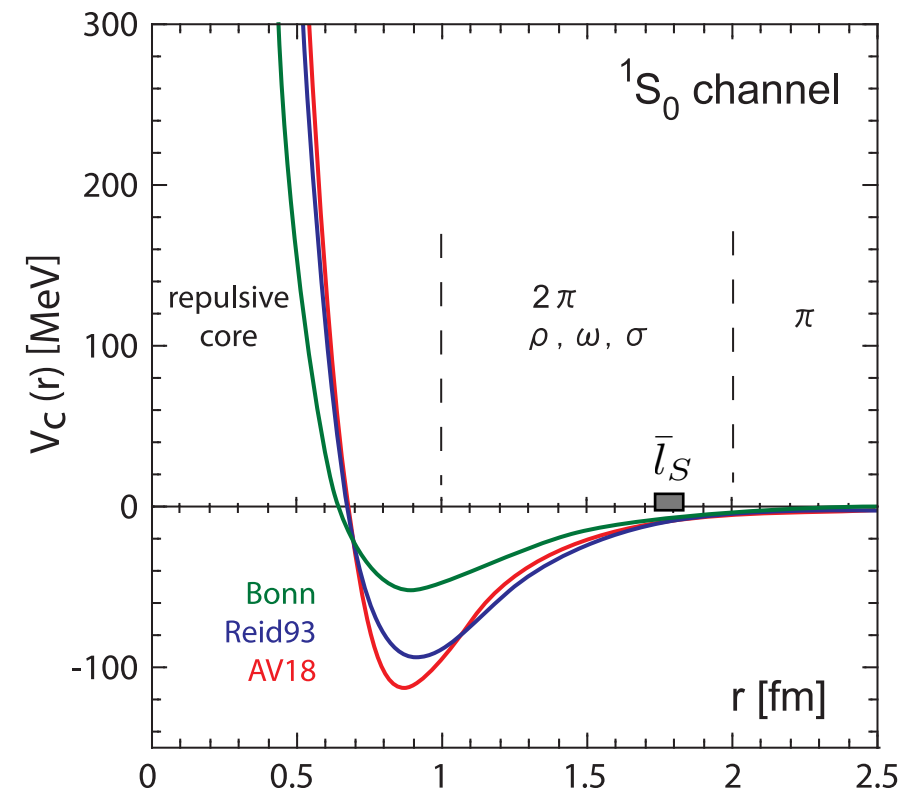
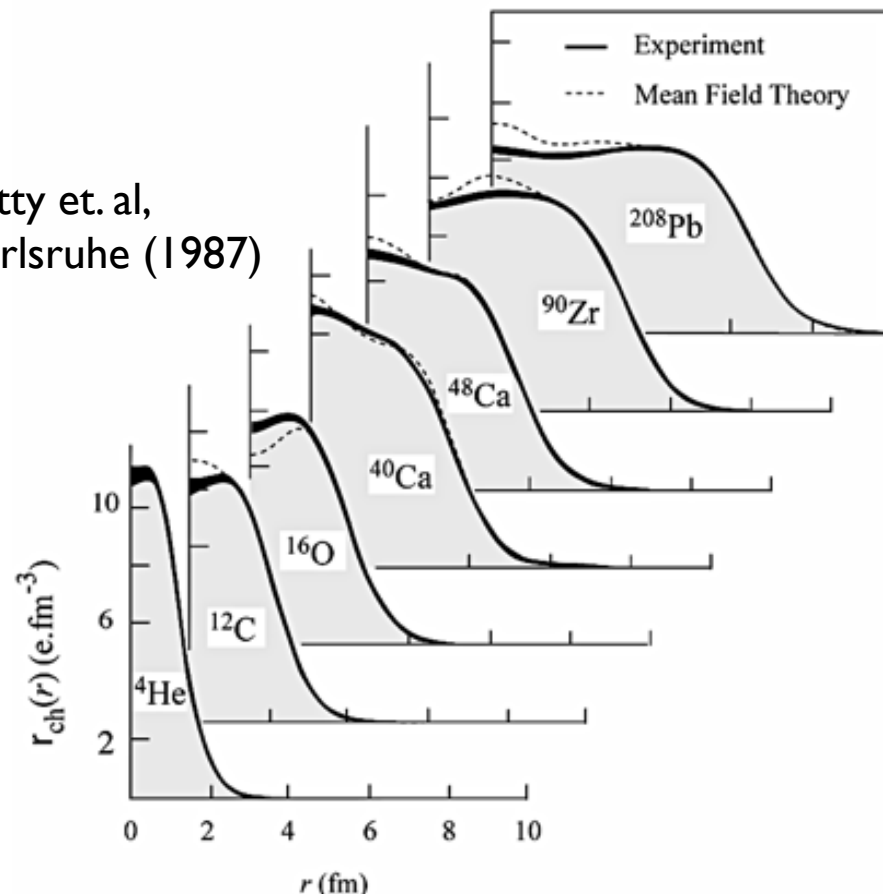
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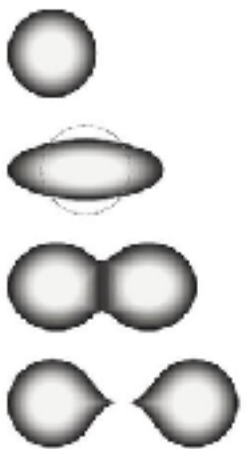
Batty et. al,  
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Semi-empirical mass formula:

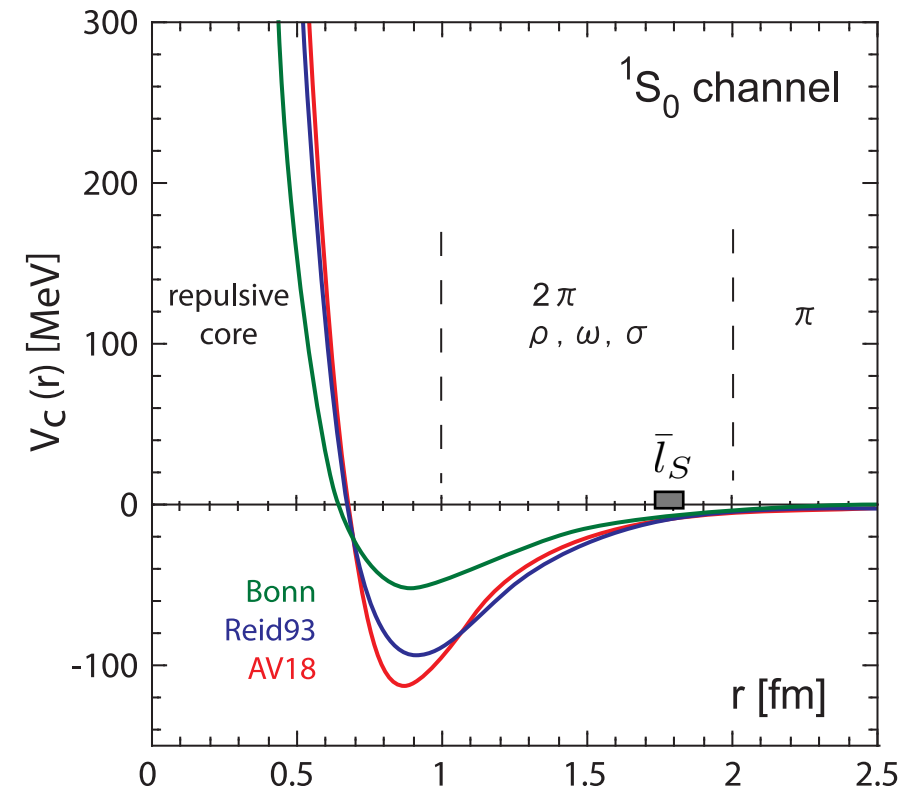
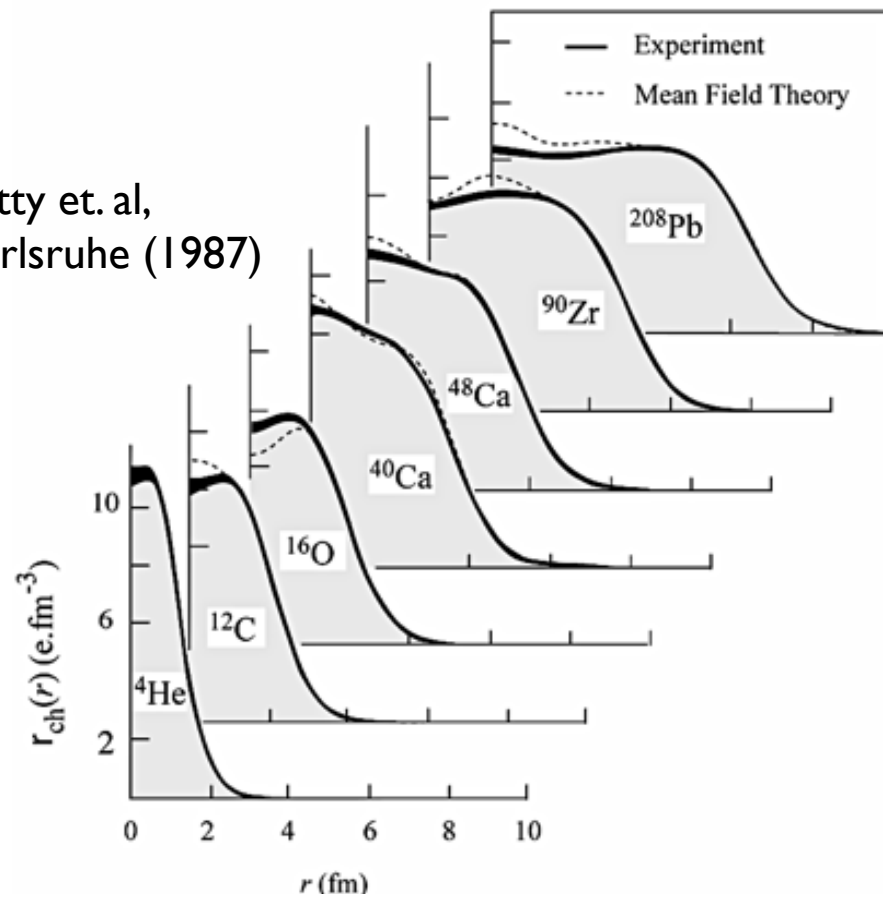
$$E_B = a_V A - a_S A^{2/3} - a_A \frac{(A - 2Z)^2}{A^{1/3}} - a_C \frac{Z(Z - 1)}{A^{1/3}} \pm \frac{a_p}{A^{1/2}}$$

$$a_V \sim 16 \text{ MeV}, a_S \sim 18 \text{ MeV}, a_A \sim 23 \text{ MeV}, a_C \sim 0.7 \text{ MeV}, a_p \sim 12 \text{ MeV}$$



# Nuclear saturation and the liquid drop model

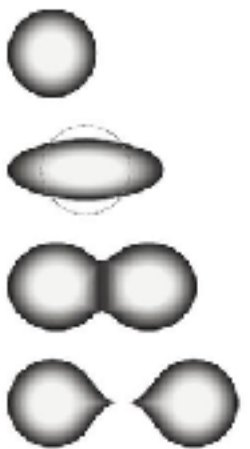
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Semi-empirical mass formula: **thermodynamic limit**  $A \rightarrow \infty$

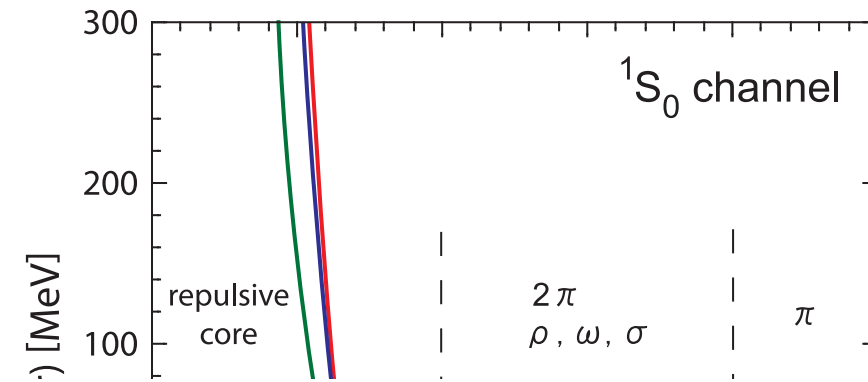
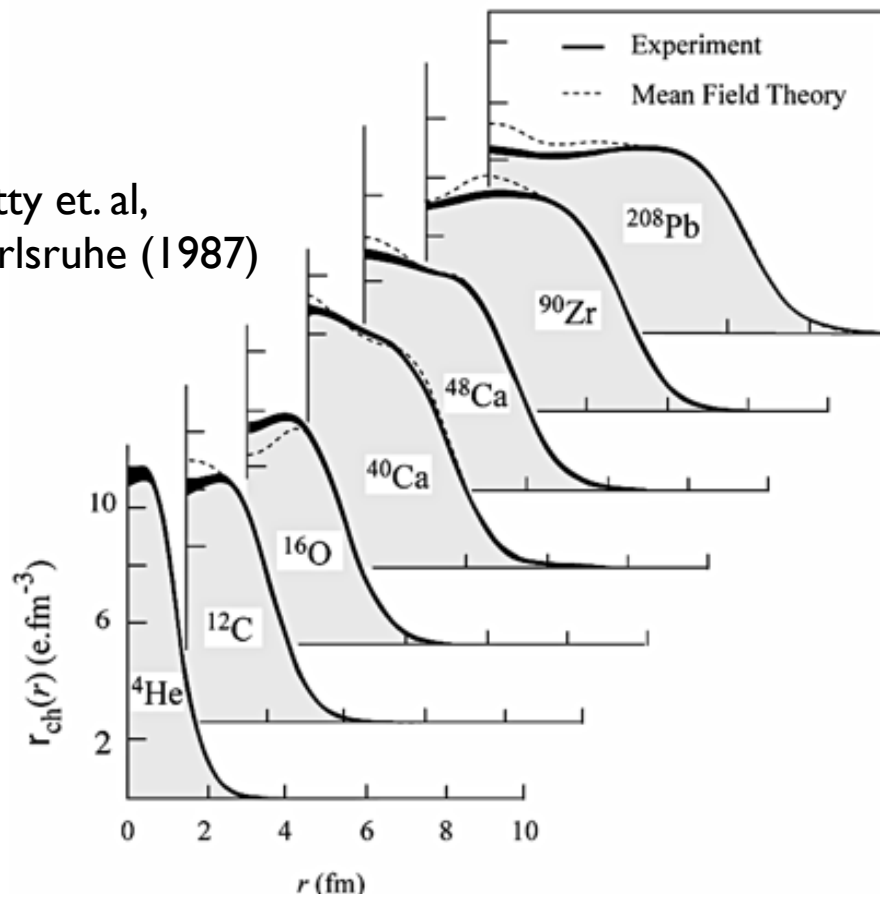
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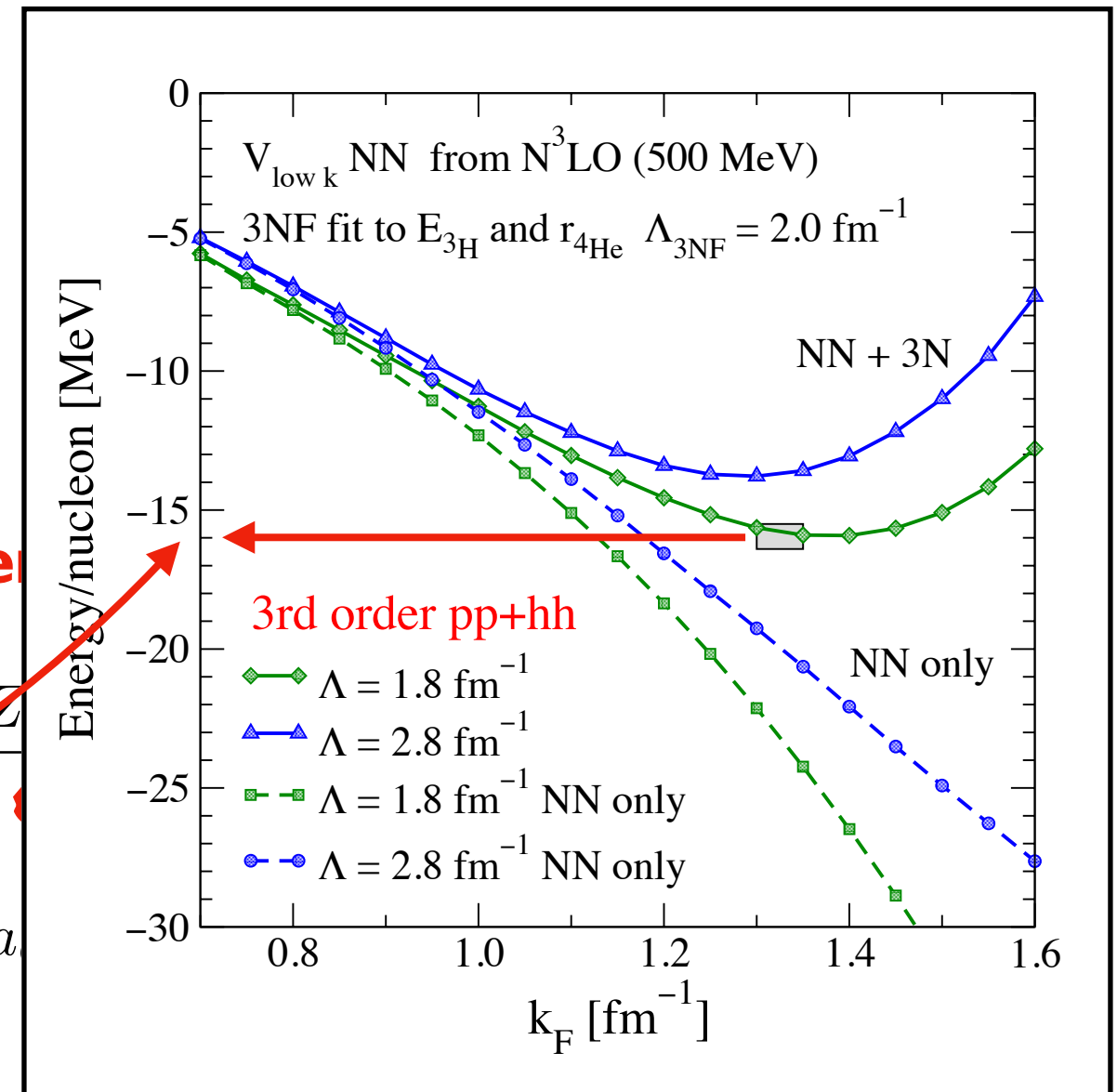
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Semi-empirical mass formula: **the**

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# Practical calculation of the EOS

Statistical mechanics reminder:

$$Z = \text{Tr} e^{-\beta H},$$

$$F = -k_B T \log Z = E - TS,$$

$$P = k_B T \frac{\partial \log Z}{\partial V}$$

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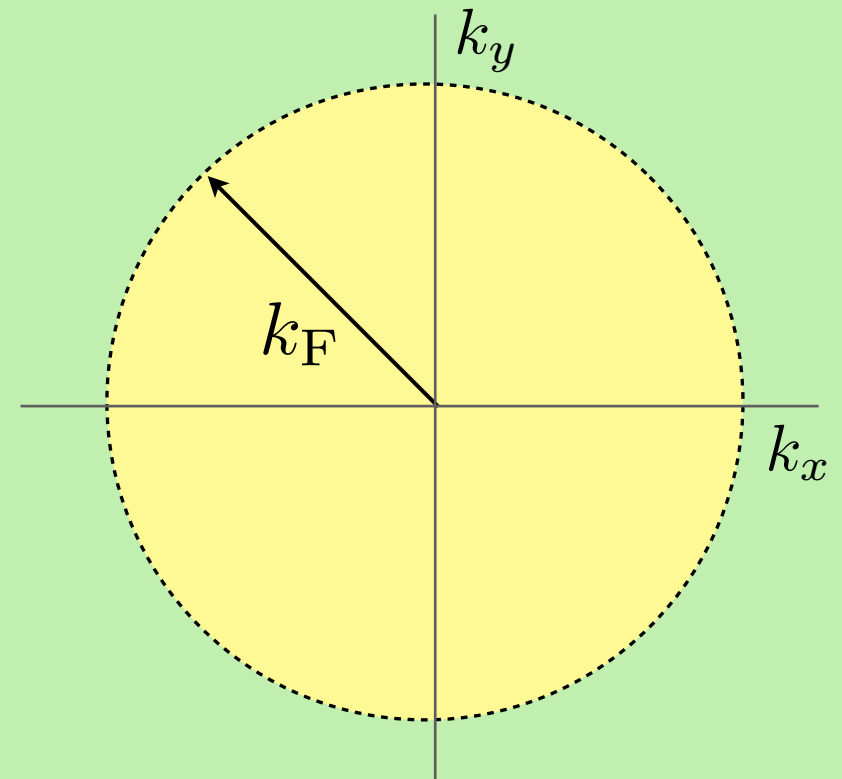
$$P = k_B T \frac{\partial \log Z}{\partial V}$$

Hamiltonian:

$$H = T + V_{NN} + V_{3N}$$

$$= H_0 + \underbrace{(-H_0 + V_{NN} + V_{3N})}_{H_1}$$

$H_0$  defines reference state:  
(e.g. free state or HF state)



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Evaluation of exact partition function in general highly nontrivial:

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One way to approximately evaluate  $Z$ : perturbation theory

$$\langle e^{-\beta H_1} \rangle_{H_0} = 1 - \beta \langle H_1 \rangle_{H_0} + \frac{\beta^2}{2!} \langle H_1^2 \rangle_{H_0} - \frac{\beta^3}{3!} \langle H_1^3 \rangle_{H_0} + \dots$$

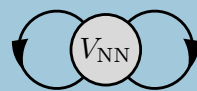
# Many-body perturbation theory: Diagrammatic representation

$E =$

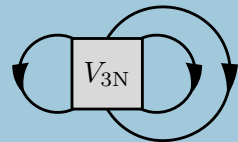


kinetic energy

+

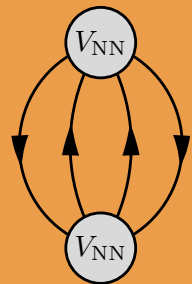


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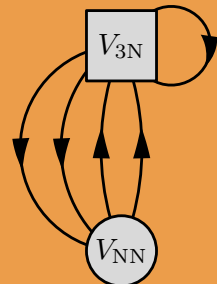


Hartree-Fock

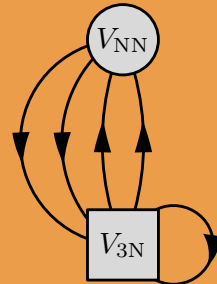
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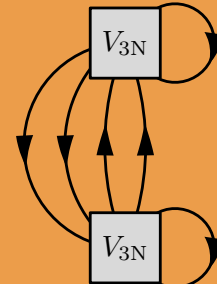
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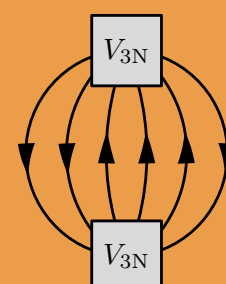
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2nd-order

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...

3rd-order  
and beyond