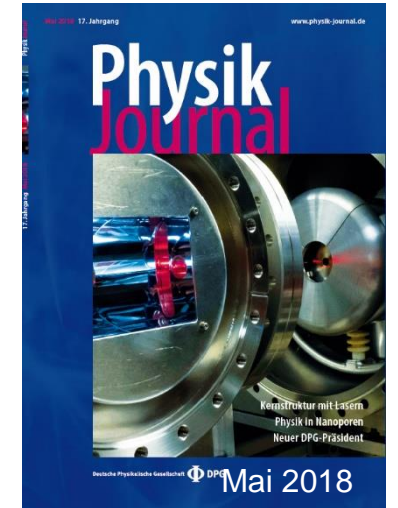
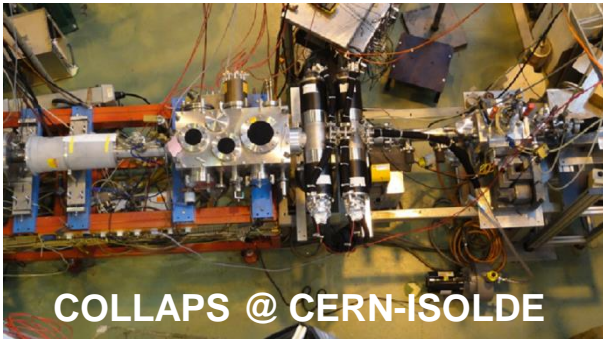


Meet 'n Greet: About "absolute" and differential charge radii from laser spectroscopy



GEFÖRDERT VOM



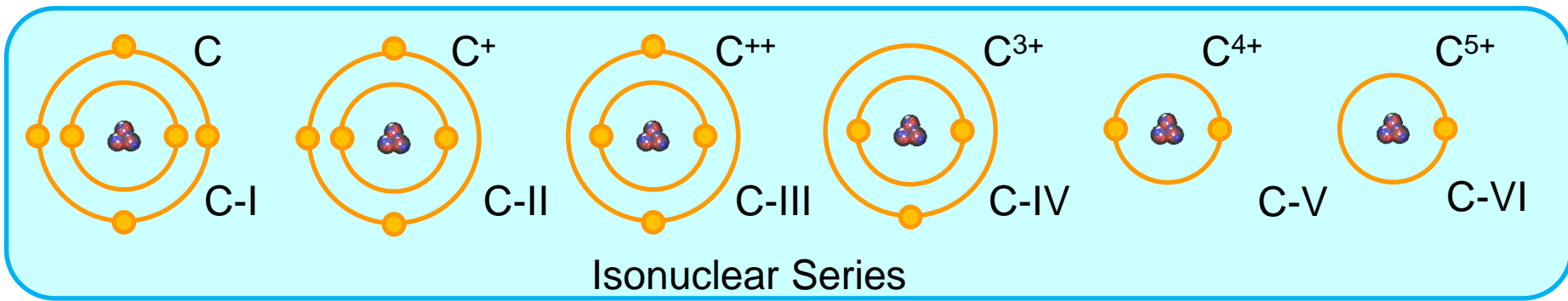
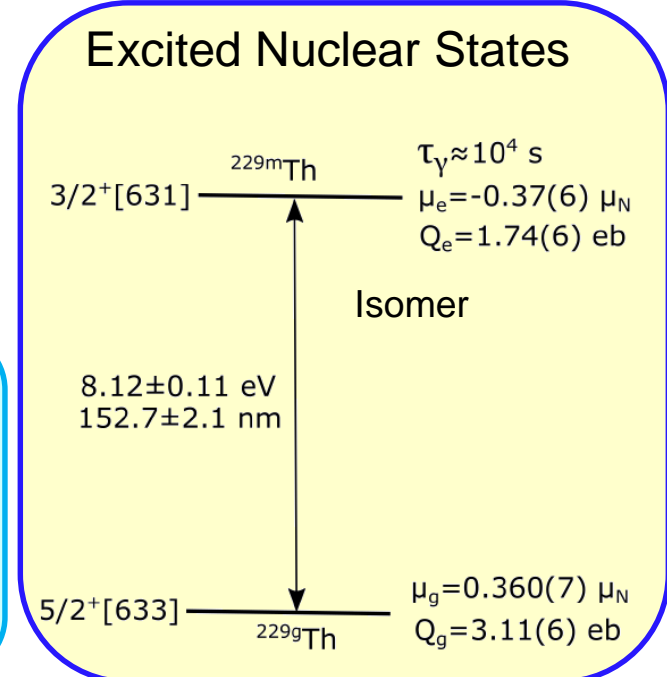
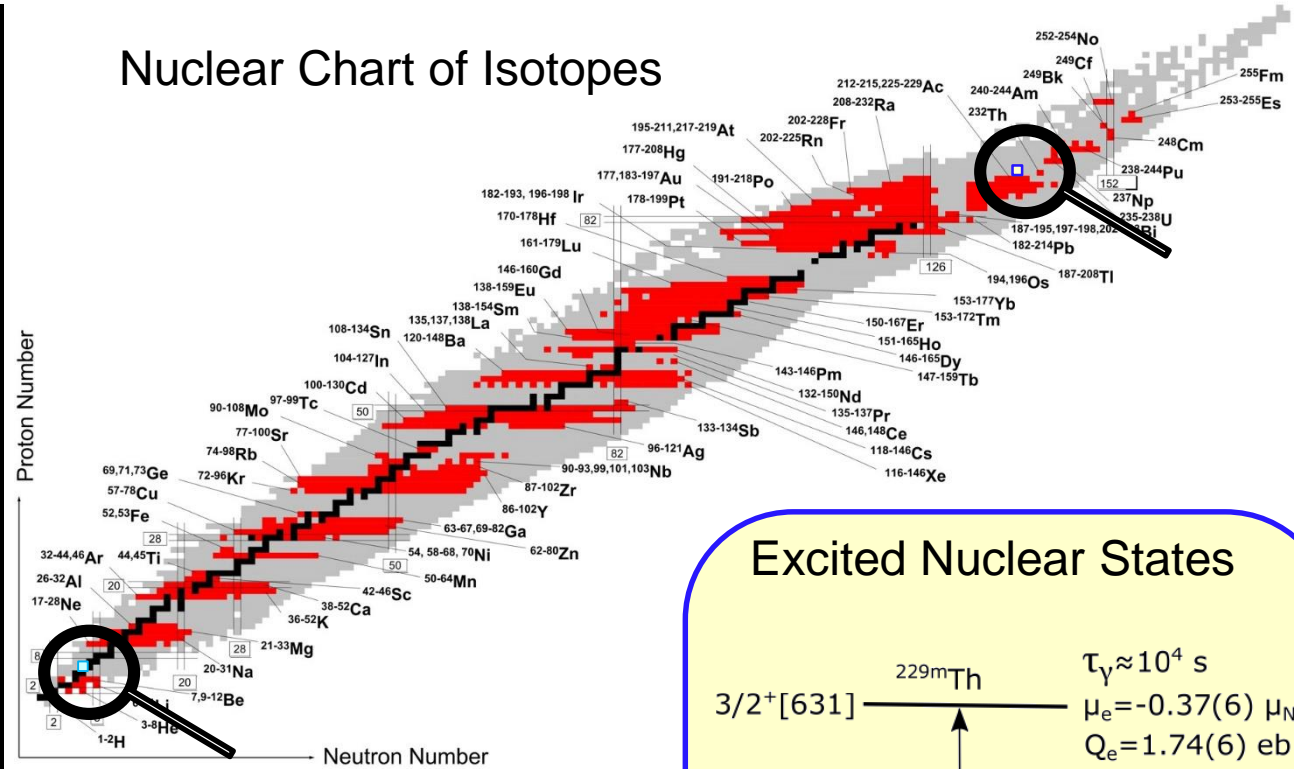
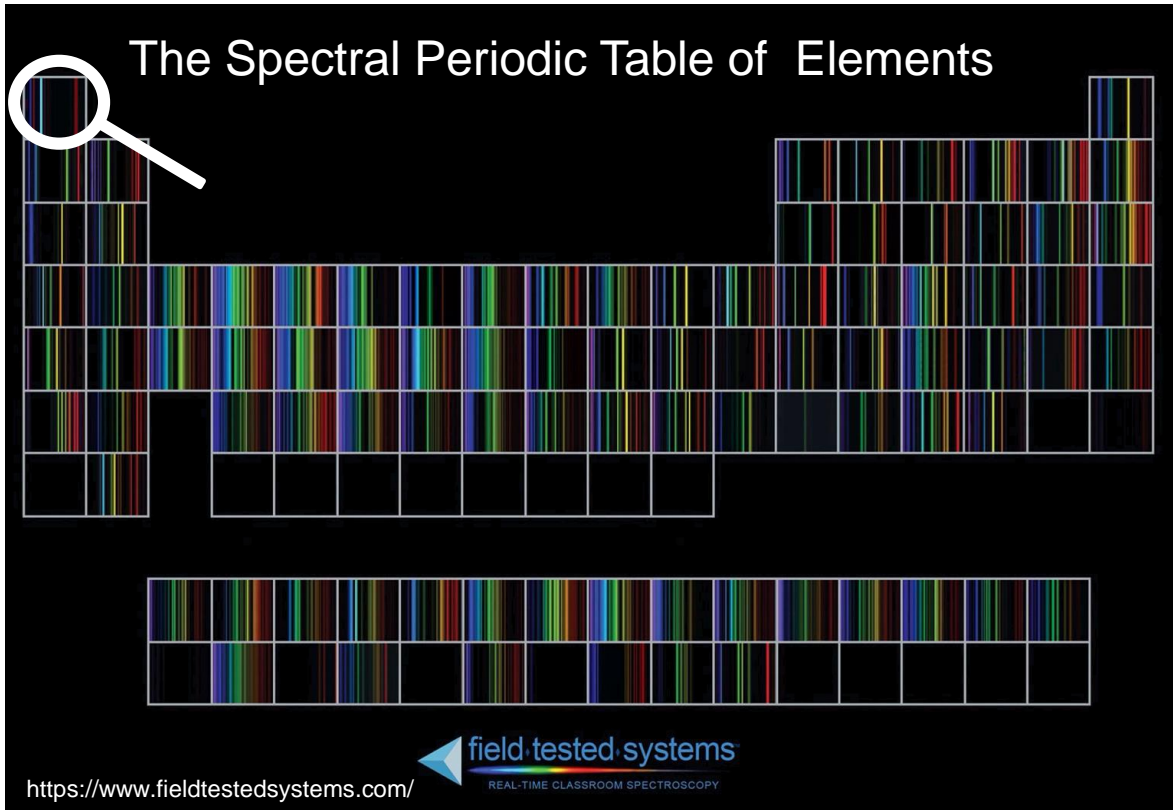
Gefördert durch
DFG Deutsche
Forschungsgemeinschaft



SFB 1245 "Nuclei: From Fundamental
Interactions to Structure and Stars



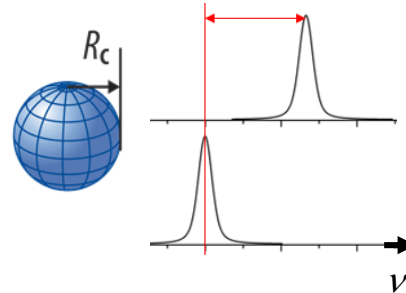
Emission Spectra of the Elements



History of Nuclear Effects in Atomic Spectra

N. Bohr, *The spectra of helium and hydrogen*
Nature, 92, 231 (1913)

H. Urey et al., *A Hydrogen Isotope of Mass 2*
Physical Review 39, 164 (1932).



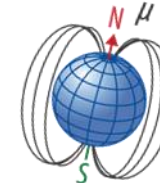
$$\delta\nu_{IS} = \delta\nu_{MS} + F \delta \langle r_c^2 \rangle^{AA'}$$

Nagaoka & Sugiura, *Spectroscopic evidence of isotopy*,
Jpn. J. Phys. 2, 167 (1923)

W. Pauli, *Zur Frage der theoretischen Deutung der Satelliten einiger Spektrallinien und ihrer Beeinflussung durch magnetische Felder*,
Naturwissenschaften 12, 741 (1924).



Spin I



Magnetic
Dipole
Moment μ

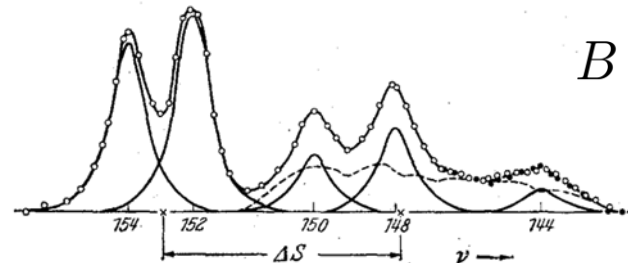
$$A = \frac{\mu_I B_e(0)}{IJ}$$

H. Schüler and T. Schmidt, *Über Abweichungen des Atomkerns von der Kugelsymmetrie*,
Zeitschrift für Physik A 94, 457 (1935).

Electric Quadrupole
Moment Q_s



P. Brix and H. Kopfermann, *Zur Isotopieverschiebung im Spektrum des Samariums*,
Zeitschrift für Physik 126, 344 (1949).

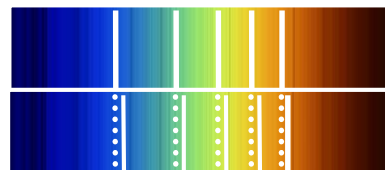


$$B = eQ_s V_{zz}$$

Mass Shift



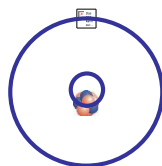
Isotope 1



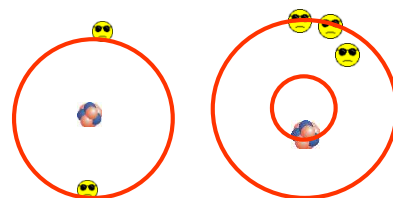
Isotope 2



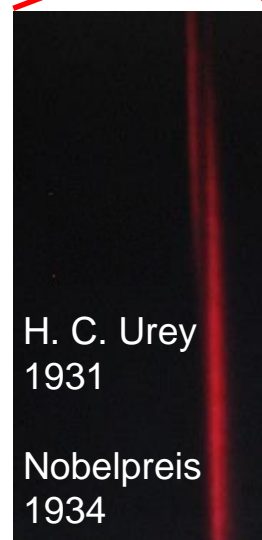
$$\delta\nu_{IS}^{AA'} = \nu^{A'} - \nu^A$$



Normal mass effect (NMS)

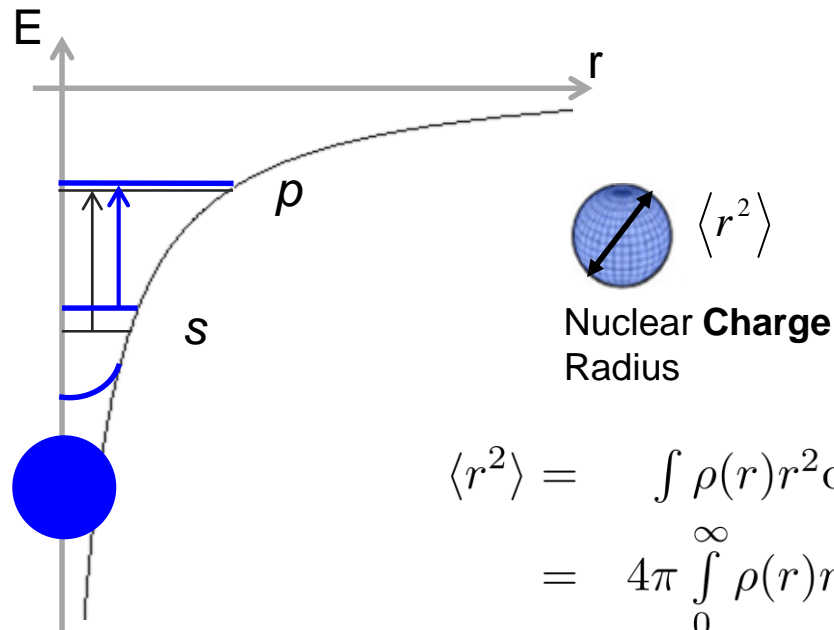


Specific mass effect (SMS)

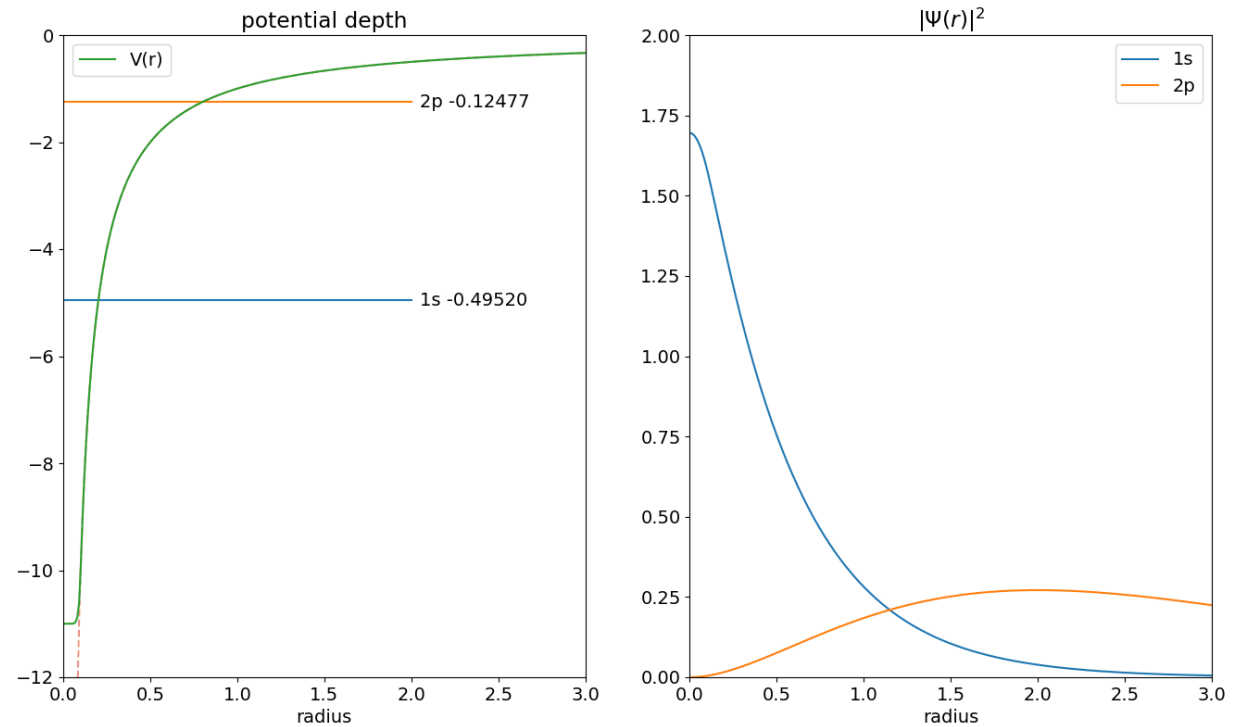


$$\begin{aligned} \delta\nu_{MS}^{AA'} &= \delta\nu_{NMS}^{AA'} + \delta\nu_{SMS}^{AA'} && \text{Mass shift (MS)} \\ &= (K_{NMS} + K_{SMS}) \cdot \frac{M_{A'} - M_A}{M_A M_{A'}}. && K_{NMS} = \nu_0 \cdot m_e \end{aligned}$$

Field Shift (Volume Shift, Finite Nuclear-Size Effect)



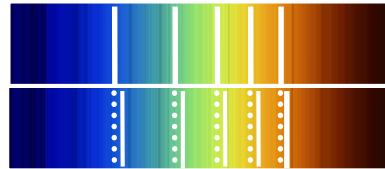
$$\begin{aligned}
 \langle r^2 \rangle &= \int \rho(r) r^2 dV \\
 &= 4\pi \int_0^\infty \rho(r) r^4 dr
 \end{aligned}$$



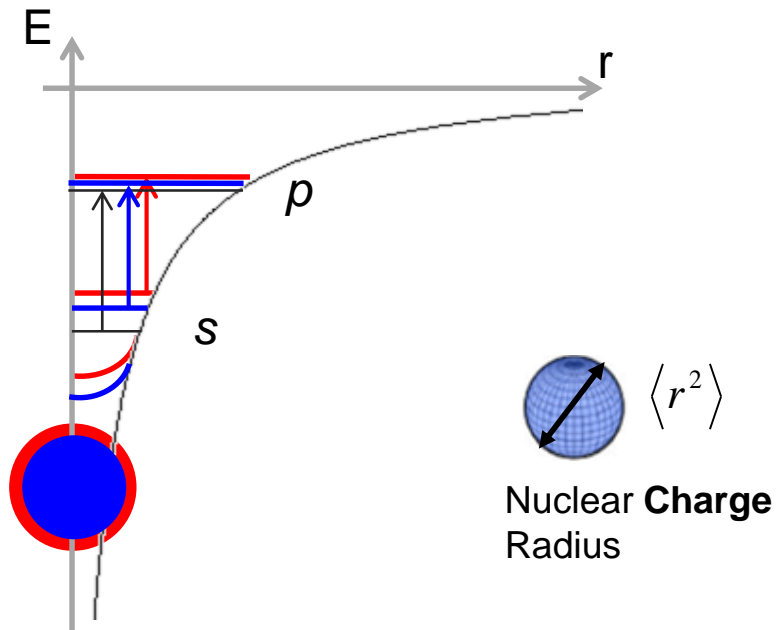
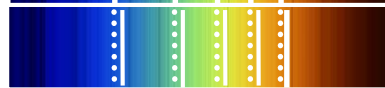
Field Shift (Volume Shift, Finite Nuclear-Size Effect)



Isotope 1



Isotope 2



$$\begin{aligned} \delta V_{\text{FS}}^{AA'} &= \frac{2\pi Ze}{3} \Delta |\Psi(0)|^2 \left(\langle r^2 \rangle^{A'} - \langle r^2 \rangle^A \right) \\ &= \underbrace{\frac{2\pi Ze}{3} \Delta |\Psi(0)|^2}_{\text{Electronic Factor}} \underbrace{\delta \langle r^2 \rangle^{AA'}}_{\text{Nuclear Size}} \end{aligned}$$

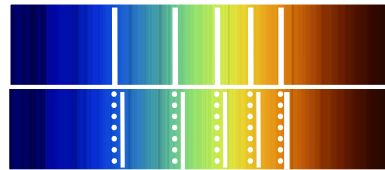
Electronic Factor
(\rightarrow Wavefunction)

Nuclear Size:
Change of Nuclear
Charge Radius

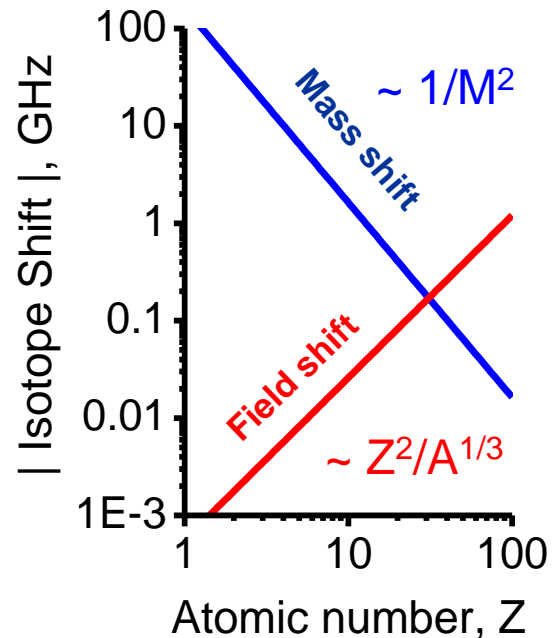
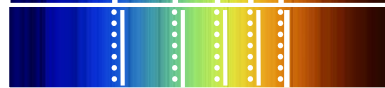
Summary Isotope Shift



Isotope 1



Isotope 2

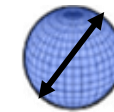


$$\begin{aligned} \delta\nu_{\text{MS}}^{AA'} &= \delta\nu_{\text{NMS}}^{AA'} + \delta\nu_{\text{SMS}}^{AA'} \\ &= (K_{\text{NMS}} + K_{\text{SMS}}) \cdot \frac{M_{A'} - M_A}{M_A M_{A'}} \end{aligned}$$

$$\begin{aligned} \delta\nu_{\text{FS}}^{AA'} &= \frac{2\pi Ze}{3} \Delta |\Psi(0)|^2 \left(\langle r^2 \rangle^{A'} - \langle r^2 \rangle^A \right) \\ &= \underbrace{\frac{2\pi Ze}{3} \Delta |\Psi(0)|^2}_{\text{Electronic Factor}} \underbrace{\delta \langle r^2 \rangle^{AA'}}_{\text{Nuclear Size}} \end{aligned}$$

Electronic Factor
(→ Wavefunction)

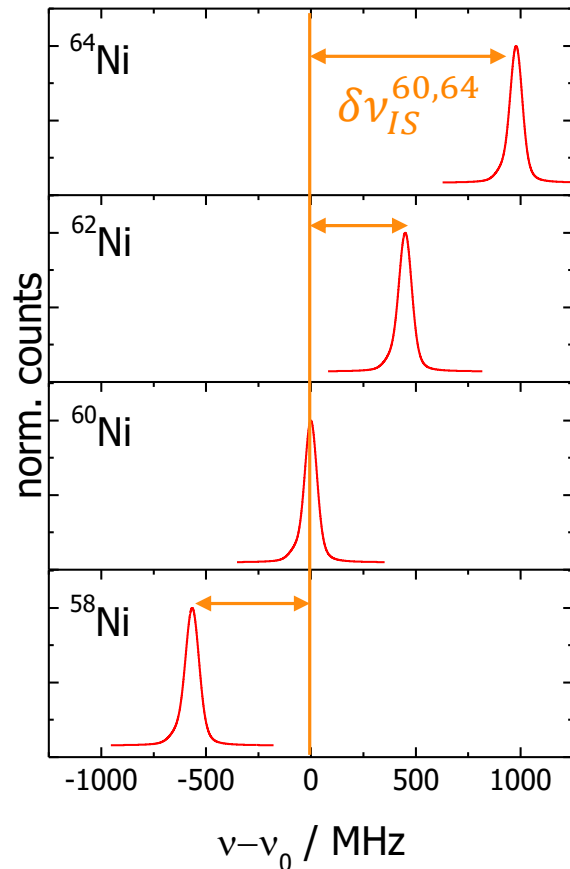
Nuclear Size:
Change of Nuclear
Charge Radius



Summary Isotope Shift (Example: Nickel)



<http://www.alexpetty.com/2014/09/21/the-periodic-table-of-light/>




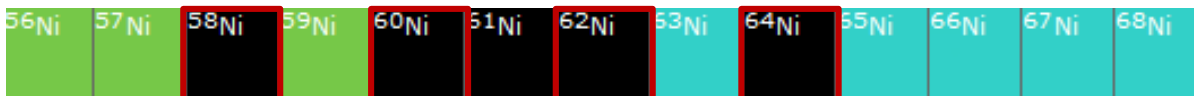
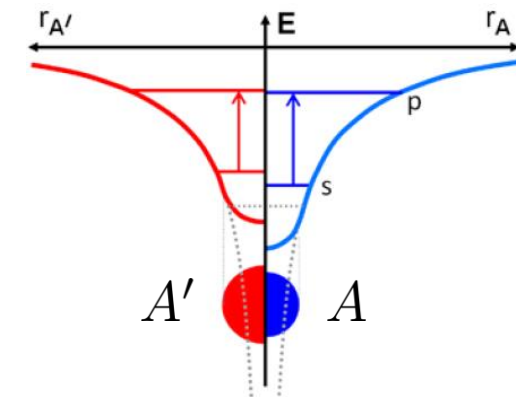
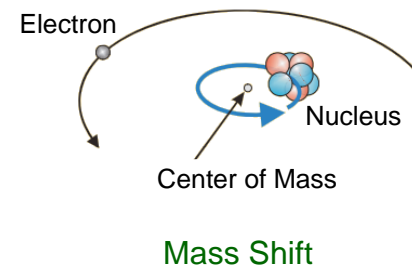
isotope shift

charge radius

$$\delta \nu_{IS}^{AA'} = K \cdot M + F \delta \langle r_c^2 \rangle^{AA'}$$

→ mass shift (cm motion)
→ field shift





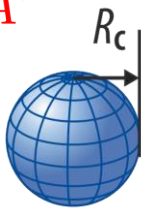
Isotope Shift

isotope shift

charge radius

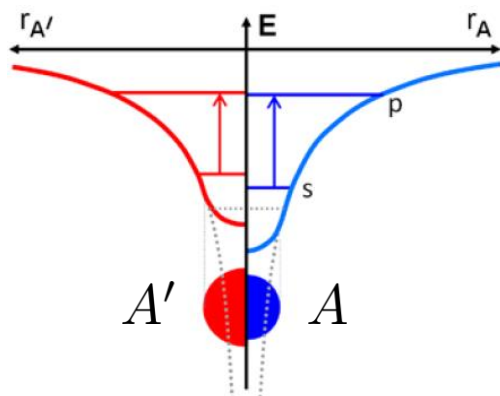
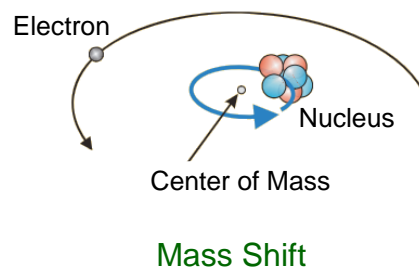
$$\delta\nu_{IS}^{AA'} = K \cdot M + F \delta \langle r_c^2 \rangle^{AA'}$$

← mass shift (cm motion) ← field shift

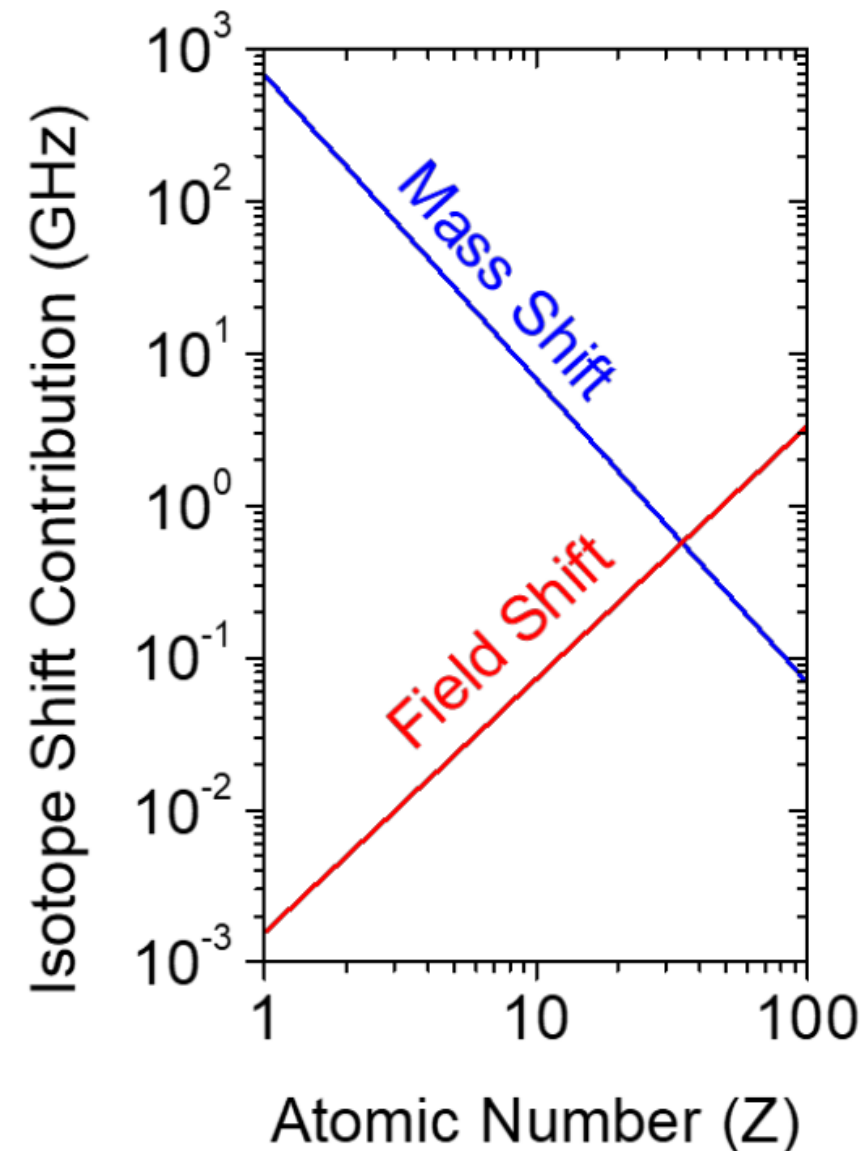


mass shift
(cm motion)

field shift



$$\delta \langle r_c^2 \rangle^{AA'} = \frac{\delta\nu_{IS}^{AA'} - K \cdot M}{F}$$



Absolute Radii of light Isotopes

$$\delta \langle r_c^2 \rangle^{AA'} = \frac{\delta \nu_{IS}^{AA'} - K \cdot M}{F}$$

Theory

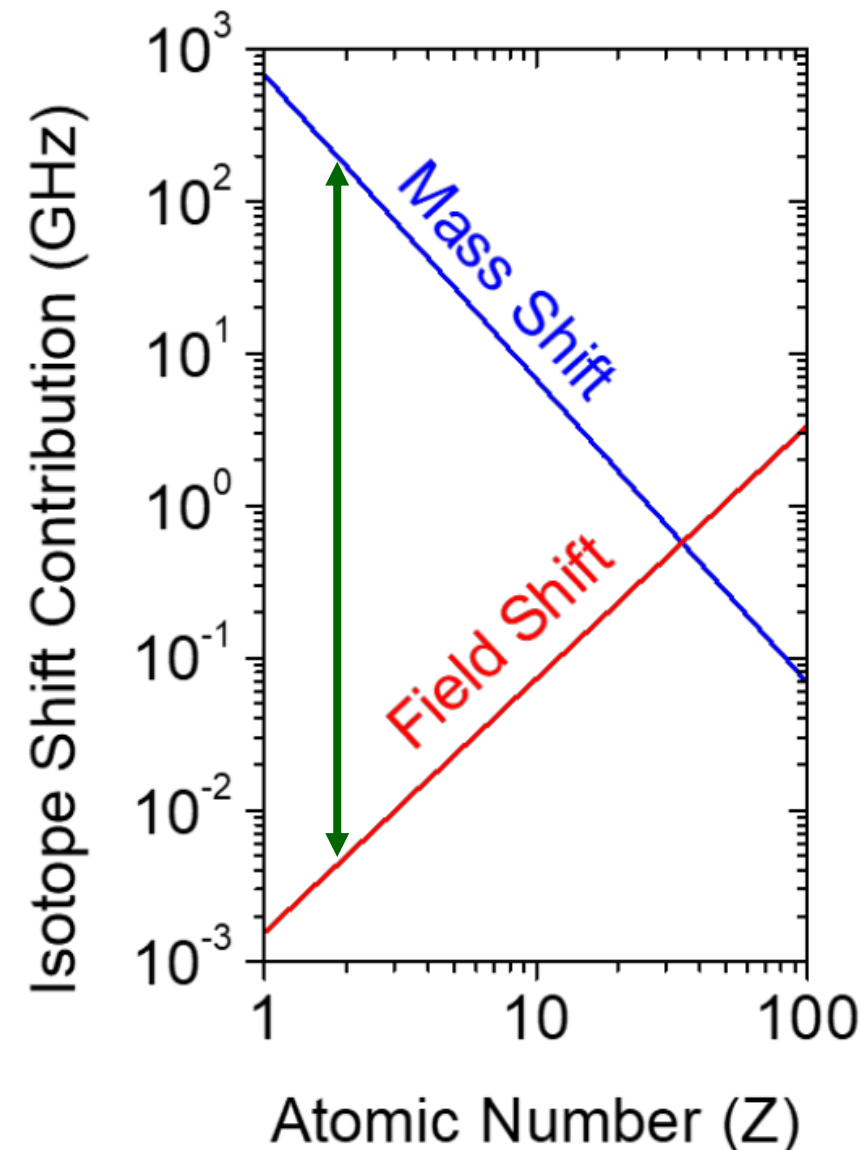
Required Accuracy $< 10^{-5}$

Mass Shift Calculations:

- He (2 e⁻) : F. Marin et al., Z. Phys. D **32**, 285 (1995).
- Li (3 e⁻) : Z.C. Yan and G.W.F. Drake, PRA **61**, 022504 (2000).
- Be (4 e⁻) : M. Puchalski et al., PRA **89**, 012506 (2014).
- B (5 e⁻) : B. Maaß et al, PRL **122**, 182501 (2019)

$$R_c(A) = \underbrace{R_c(A_{\text{ref}})} + \delta \langle r_c^2 \rangle^{A_{\text{ref}}, A}$$

Reference radius required from a different technique !



From where do we get a Reference Radius ?

(1) Elastic Electron Scattering

Form Factor: $|F(q^2)|^2 = \frac{\left(\frac{d\sigma}{d\Omega}\right)_{\text{exp}}}{\left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}}}$ $q = \frac{2E}{\hbar c} \cdot \sin^2(\theta/2)$

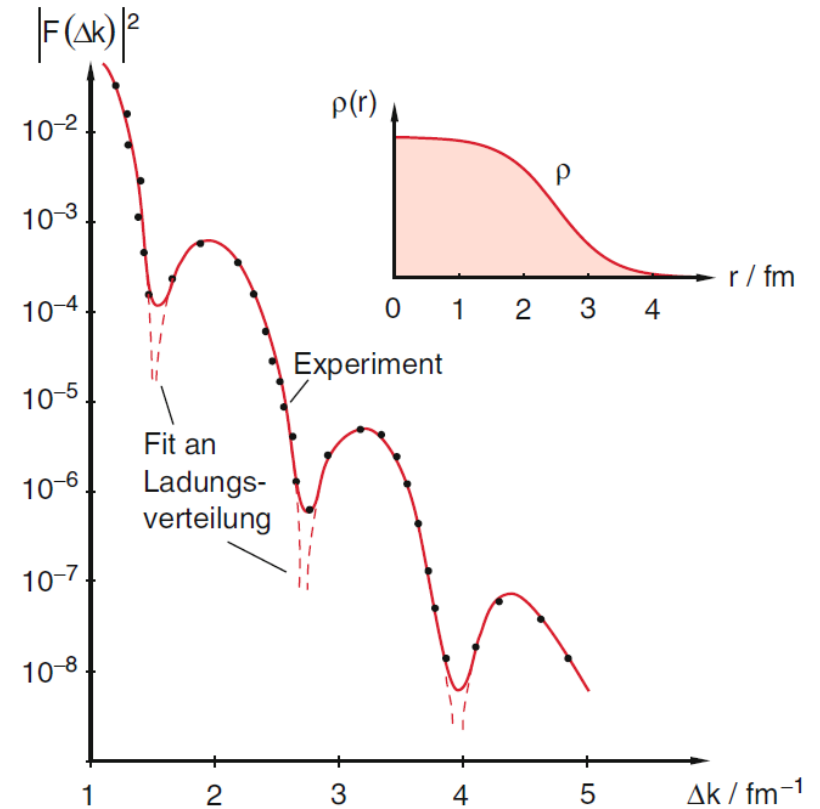
$$\rho_N(r) = \frac{1}{2\pi} \int_0^\infty dq q^2 F(q) \frac{\sin(qr)}{qr}$$

or in more general terms:

$$\rho_N(r) = \frac{1}{(2\pi)^3} \int d^3q e^{-i\vec{q}\cdot\vec{r}} F(q).$$

The form factor is the Fourier transform of the charge distribution.

- To some extent model-dependent
- Contributions of inelastic scattering and higher order contributions must be excluded

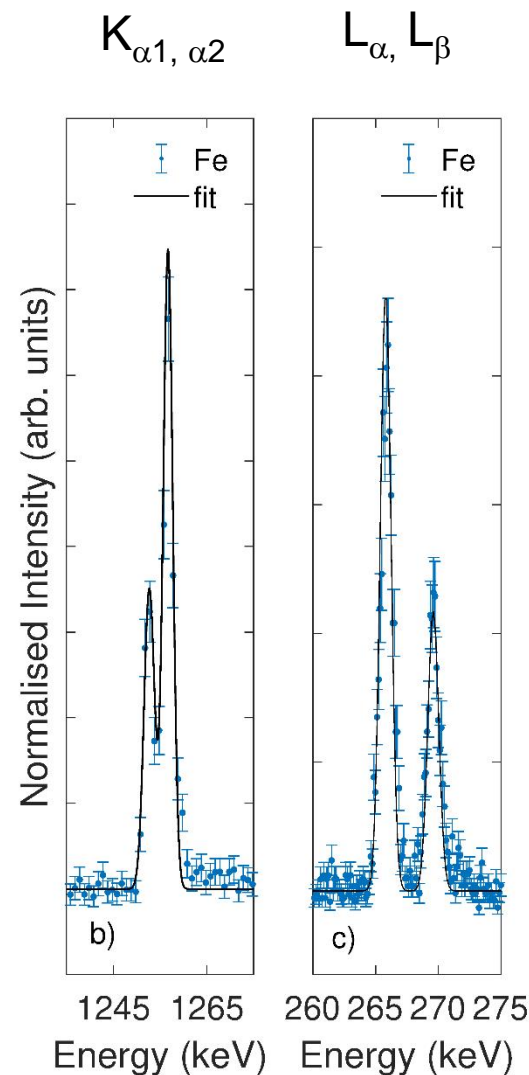
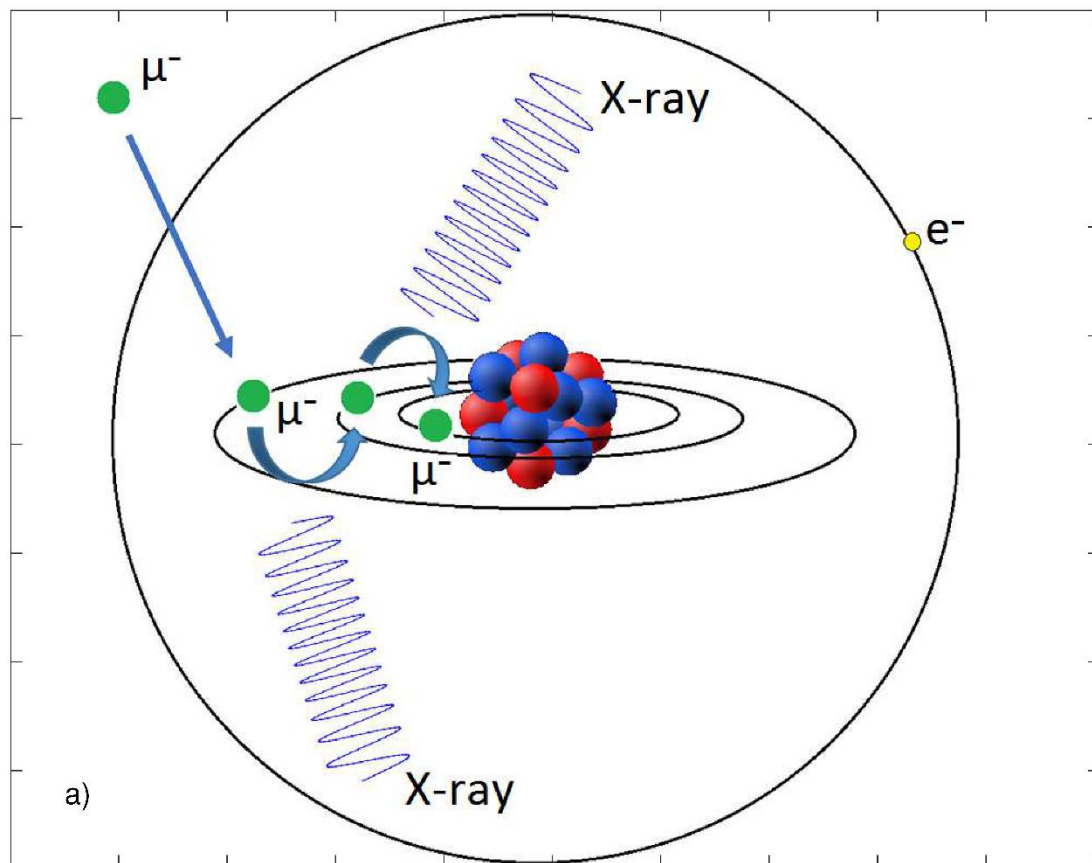


Scattering of 750 MeV electrons on O atoms. The dashed curve is calculated based on the density distribution shown in the inlet.

[Demtröder, Experimentalphysik IV]

From where do we get a Reference Radius ?

(2) Muonic Atoms X-Ray Spectroscopy



Barret Radii:

$$\langle r^k e^{-\alpha r} \rangle = \frac{4\pi}{Ze} \int \rho_c(r) r^k e^{-\alpha r} r^2 dr$$

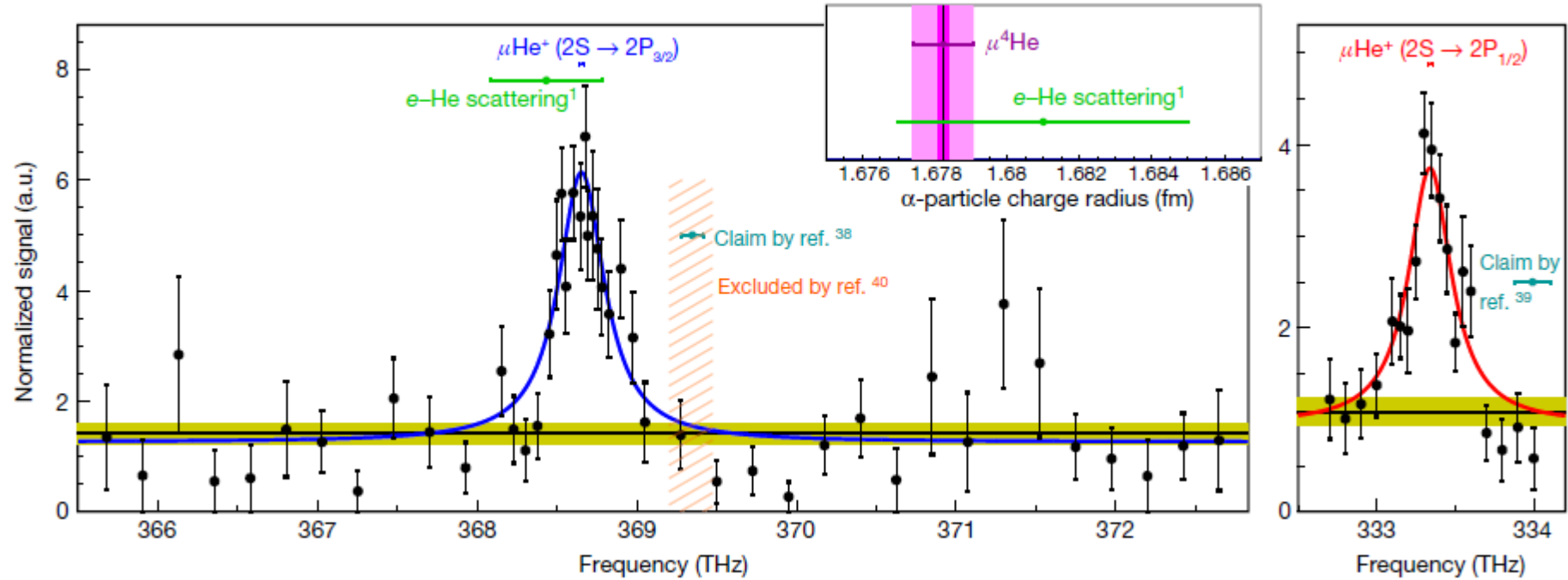
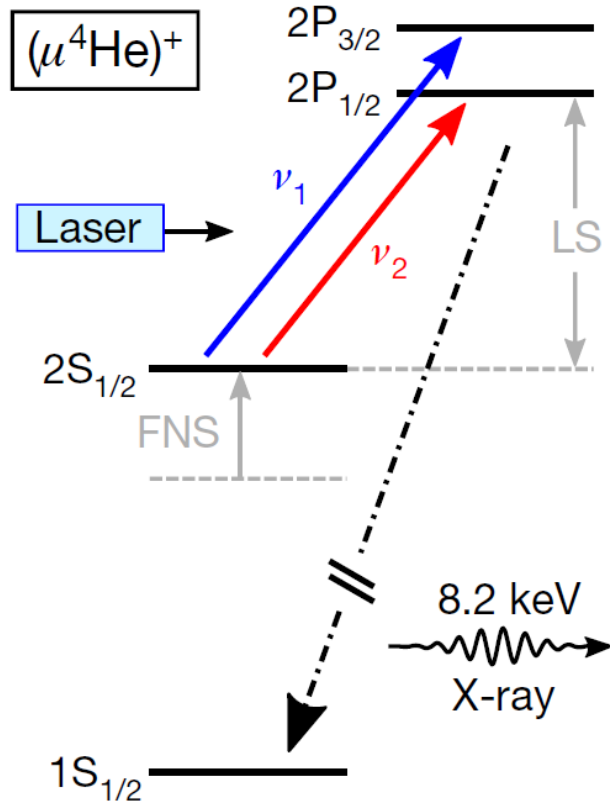
Largest uncertainty contribution:
Nuclear polarization effects

Combination with electron scattering data (ratio of moments) provides reliable reference radii.

From where do we get a Reference Radius ?

(3) Muonic Atoms Laser Spectroscopy

$$\begin{aligned} \Delta E_{2P_{1/2}-2S}^{\text{theo}} &= 1,668.489(14) \text{ meV} \\ &\quad - 106.220(8) \text{ meV fm}^{-2} \times r_{\alpha}^2 + 0.0112 \text{ meV} \\ &\quad + 9.340(250) \text{ meV} \\ &\quad - 0.150(150) \text{ meV.} \end{aligned}$$



$$\Delta E_{2P_{1/2}-2S}^{\text{exp}} = 1,378.521 \pm 0.048 \text{ meV}, \quad r_{\alpha} = 1.67824(13)_{\text{exp}} (82)_{\text{theo}} \text{ fm.}$$

Atomic Theory

Calculate the energy of an atomic state (Bohr):

$$E_n = -\frac{Z^2 \alpha^2 m_e c^2}{n^2} \frac{1}{2}$$

Employ perturbation theory (in α):

$$E_{tot} = E_{NR} + \alpha^2 E_{rel} + \alpha^3 E_{QED} + \dots + \Delta E_{nuc}$$

And again (in μ/m):

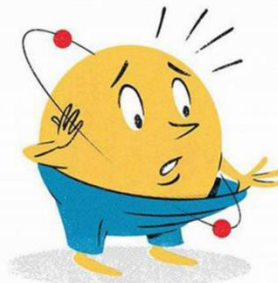
$$E_{NR} = E_{NR}^{(0)} + \left(\frac{\mu}{M}\right) E_{NR}^{(1)} + \left(\frac{\mu}{M}\right)^2 E_{NR}^{(2)} + \dots$$

$$E_{rel} = E_{rel}^{(0)} + \left(\frac{\mu}{M}\right) E_{rel}^{(1)} + \dots$$

$$E_{QED} = E_{QED}^{(0)} + \left(\frac{\mu}{M}\right) E_{QED}^{(1)} + \dots$$

largest
uncertainty

$$\Delta E_{nuc} = \frac{2\pi Z e^2 r_c^2}{3} \sum_i \langle \delta^3(r_i) \rangle = |\Psi(0)|^2$$



Exact energy calculation possible for hydrogen ($1e^-$)

See: proton radius puzzle...

[Udem et al., Phys. Rev. Lett. 79, 2646 (1997)]
[R. Pohl et al., Nature 466, 213 EP (2010)]



Isotope Shift

Subtract Energies for two isotopes, a and b: Only mass-dependent terms stay!

$$\begin{aligned}\Delta E(a - b) = & \left[\left(\frac{\mu}{M} \right)_a - \left(\frac{\mu}{M} \right)_b \right] \left(E_{NR}^{(1)} + \alpha^2 E_{rel}^{(1)} + \alpha^3 E_{QED}^{(1)} \right) \\ & + \left[\left(\frac{\mu}{M} \right)_a^2 - \left(\frac{\mu}{M} \right)_b^2 \right] E_{NR}^{(2)} + \dots + \frac{2\pi Z e^2}{3} |\Psi(0)|^2 [r_{c,a}^2 - r_{c,b}^2]\end{aligned}$$

Isotope Shift

Subtract Energies for two isotopes, a and b: Only mass-dependent terms stay!

$$\Delta E(a - b) = \left[\left(\frac{\mu}{M} \right)_a - \left(\frac{\mu}{M} \right)_b \right] \left(E_{NR}^{(1)} + \alpha^2 E_{rel}^{(1)} + \alpha^3 E_{QED}^{(1)} \right) + \left[\left(\frac{\mu}{M} \right)_a^2 - \left(\frac{\mu}{M} \right)_b^2 \right] E_{NR}^{(2)} + \dots + \frac{2\pi Z e^2}{3} |\Psi(0)|^2 [r_{c,a}^2 - r_{c,b}^2]$$

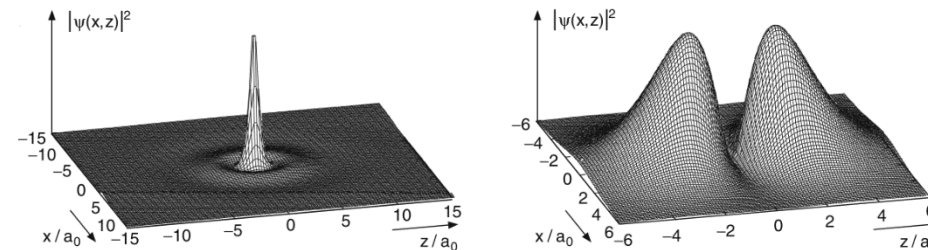
Test a *transition*, not an *energy*:

$$\Delta \nu_{IS} = \delta \nu_{MS} - F_{i \rightarrow f} [r_{c,a}^2 - r_{c,b}^2]$$

$$F_{i \rightarrow f} \propto (|\Psi_f(0)|^2 - |\Psi_i(0)|^2)$$

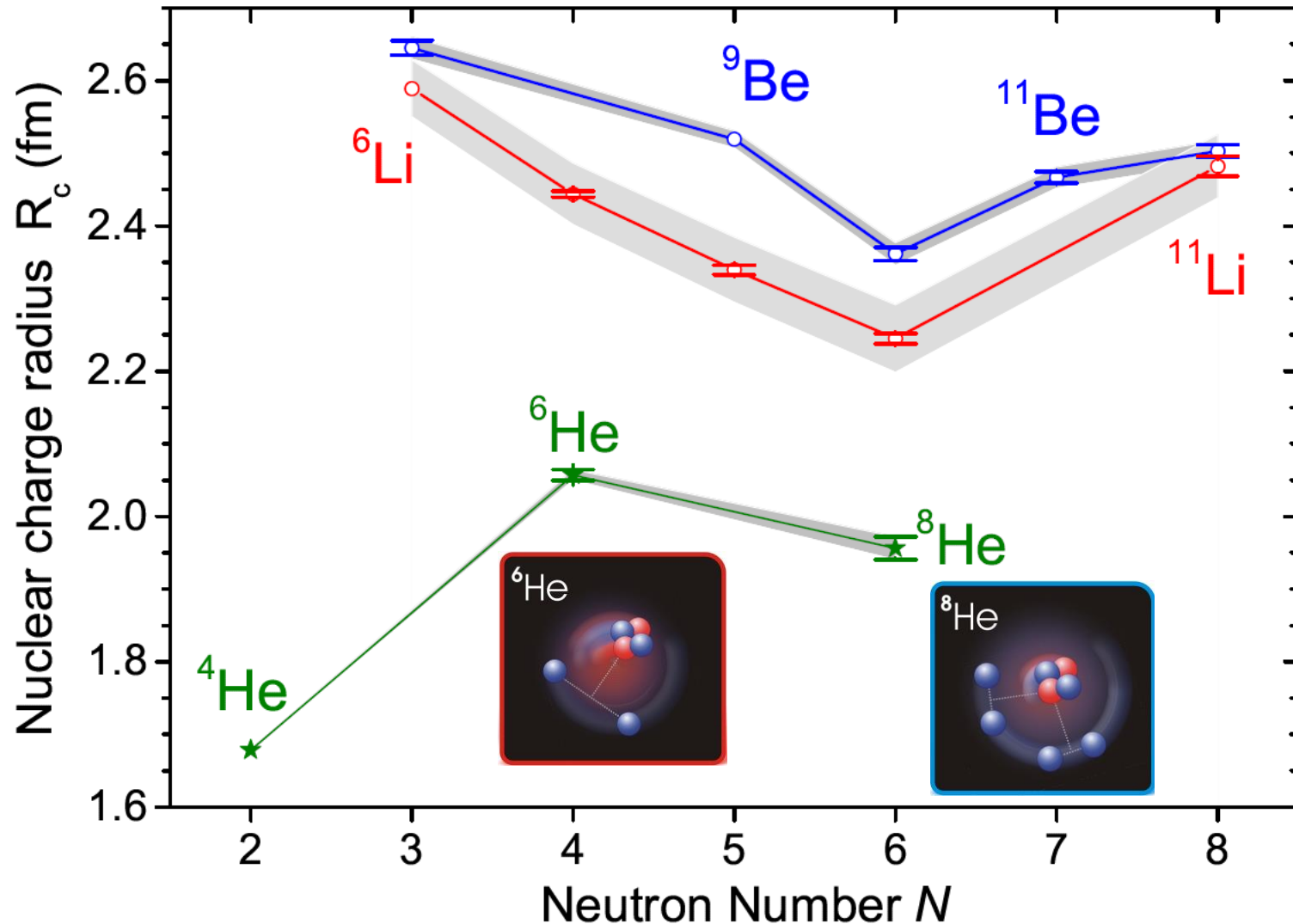
[Drake et al., Can. J. Phys. 83 311 (2005)]
[Lu et al., Rev. Mod. Phys. 85, 1383 (2013)]

Better probe the IS in a S-P transition!



[Demtröder Experimentalphysik 3, Springer, 2005]

Nuclear Radii of the Lightest Isotopes



Uncertainties in charge radii:

$$R_c(A) = \underbrace{R_c(A_{\text{ref}})}_{\text{Grey Regions } \sigma(R_c)} + \underbrace{\delta \langle r_c^2 \rangle^{A_{\text{ref}}, A}}_{\text{Error Bars } \sigma(\delta \nu_{\text{IS}})}$$

$$R_c(^9\text{Be}) = 2.519(12) \text{ fm}$$

Jansen *et al.*, Nucl. Phys. A 188, 337 (1972)

$$R_c(^6\text{Li}) = 2.589(39) \text{ fm}$$

Nörtershäuser *et al.*, Phys. Rev. C 84, 024307 (2011)

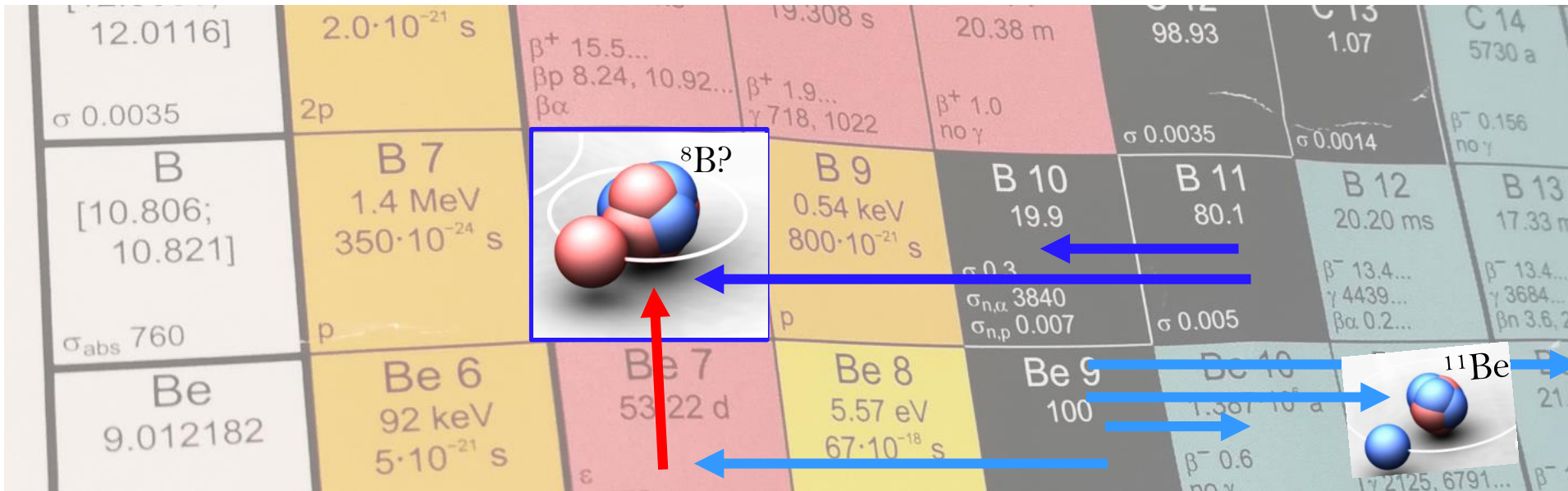
$$r_\alpha = 1.67824(83) \text{ fm}$$

Krauth *et al.*, Nature 589, 527(2021)

from μHe^+ (one-electron system)

“all-optical charge radius”

The Proton-Halo Nucleus ${}^8\text{B}$



$$\delta\nu_{\text{IS}} - \delta\nu_{\text{MS}}^{\text{Theory}} \propto \delta\langle r_c^2 \rangle$$

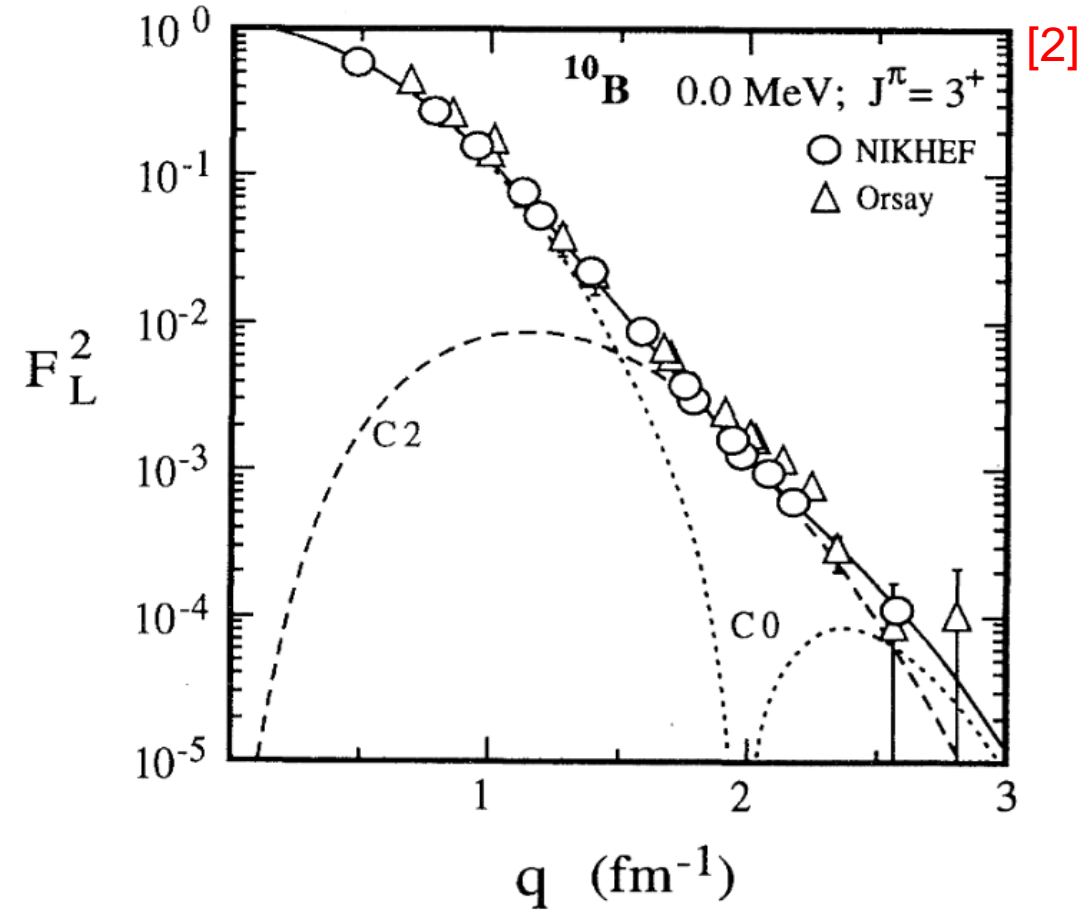
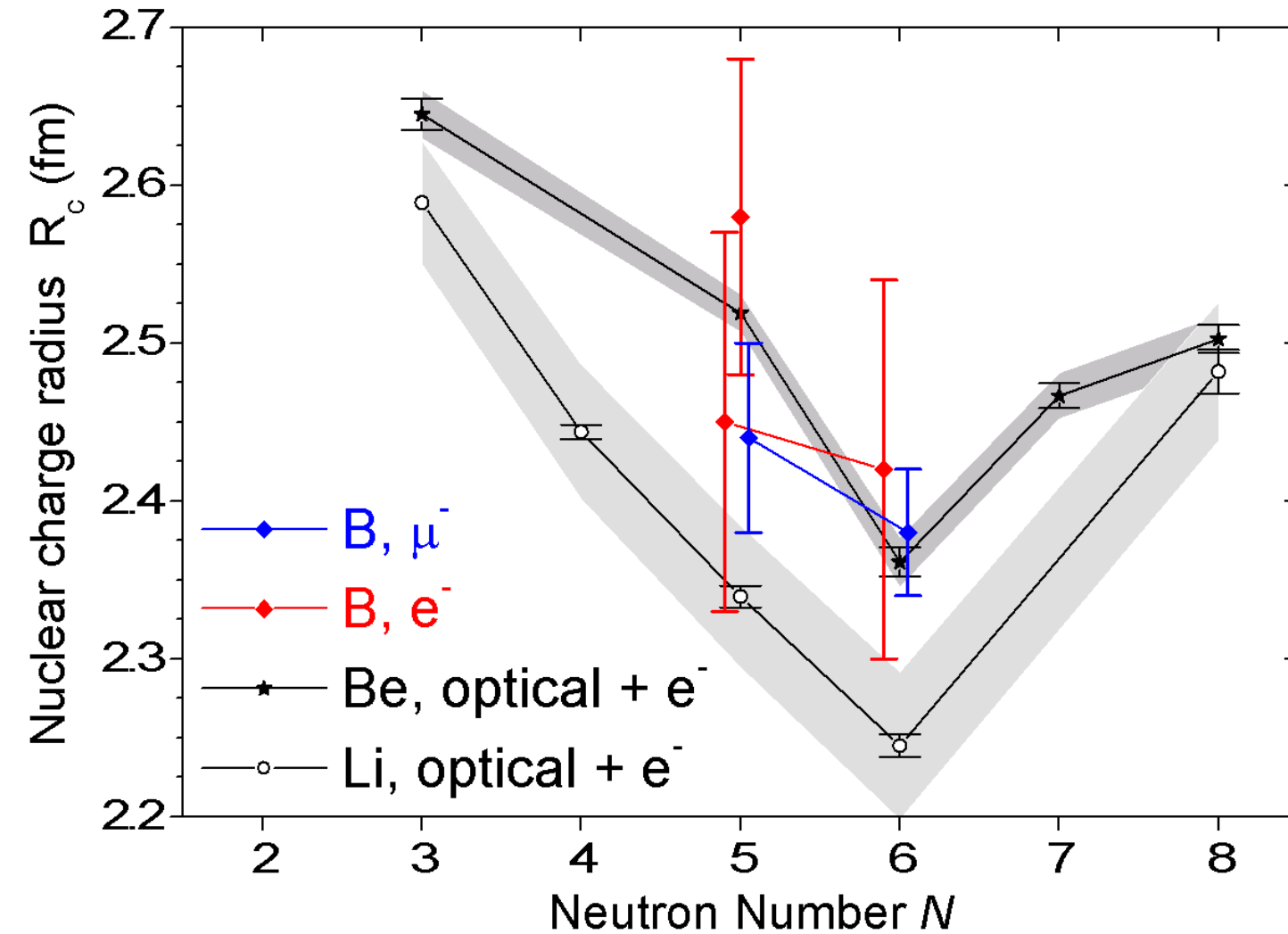
$$R_c(A) = R_c(A_{\text{ref}}) + \delta\langle r_c^2 \rangle^{A_{\text{ref}}, A}$$

„Proton-halo size“: $R_c(\text{p}_{\text{halo}}) = R_c({}^8\text{B}) - R_c({}^7\text{Be})$

Reference Radius required

Conclusion: To gain information about the proton halo of ${}^8\text{B}$, we need reliable reference radii for Be and B on equal footing !

The „Tragedy“ of Boron Reference Radii



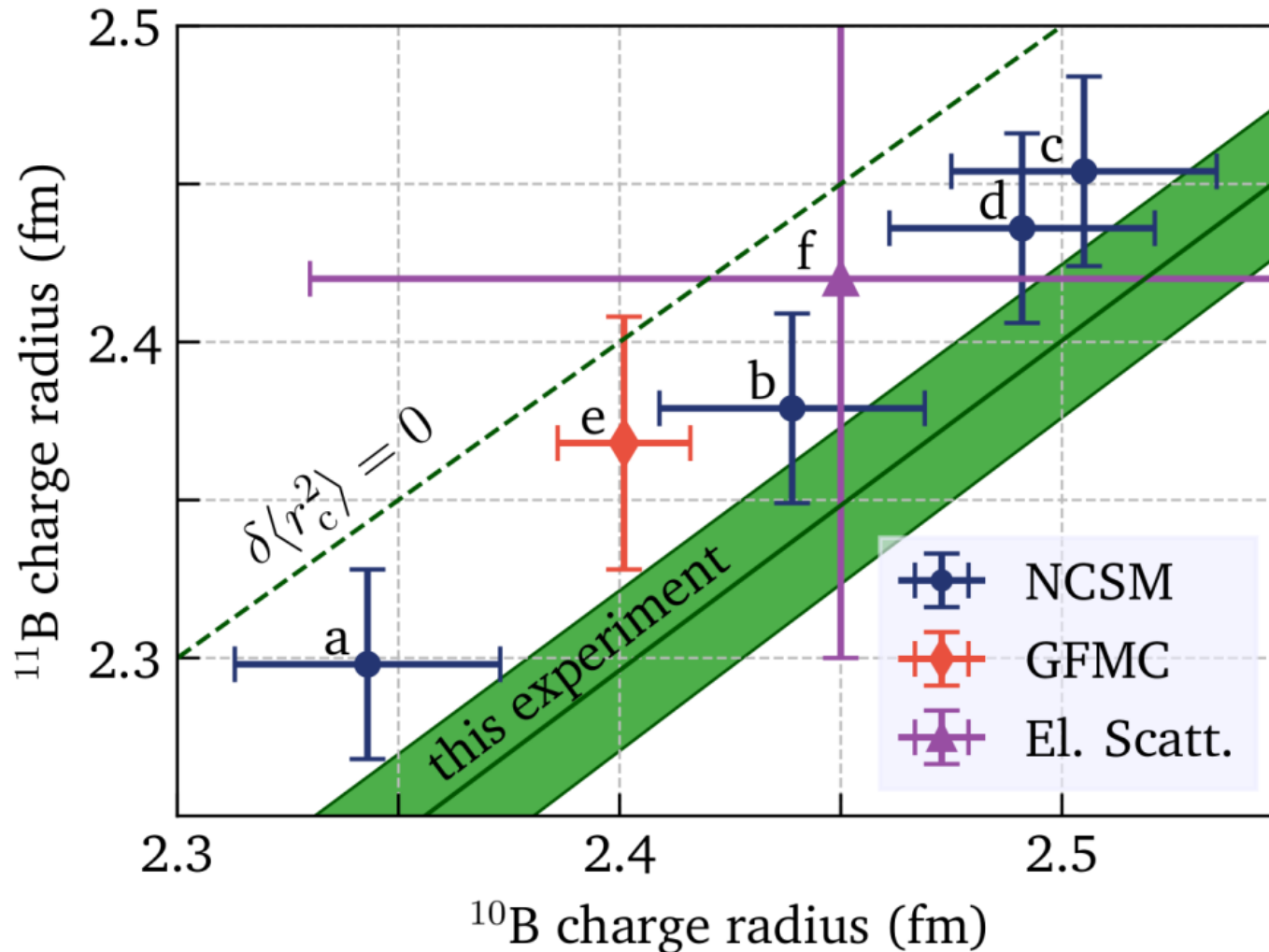
[1] Stovall, Nucl. Phys. 86, 225 (1966)

[2] Cichocki et al., PRC 51, 2406 (1995)

[3] Schaller et al., Nucl. Phys. A 343, 333 (1980)

[4] Olin et al., Nucl. Phys. A 360, 426 (1981)

Isotope Shift Measurement of $^{10,11}\text{B}$



NCSM: No-Core Shell Model

(a) $\text{N}2\text{LO}_{\text{Sat}}$, (b,c) EM, $\text{N}3\text{LO}$ (d) EMN, $\text{N}4\text{LO}$

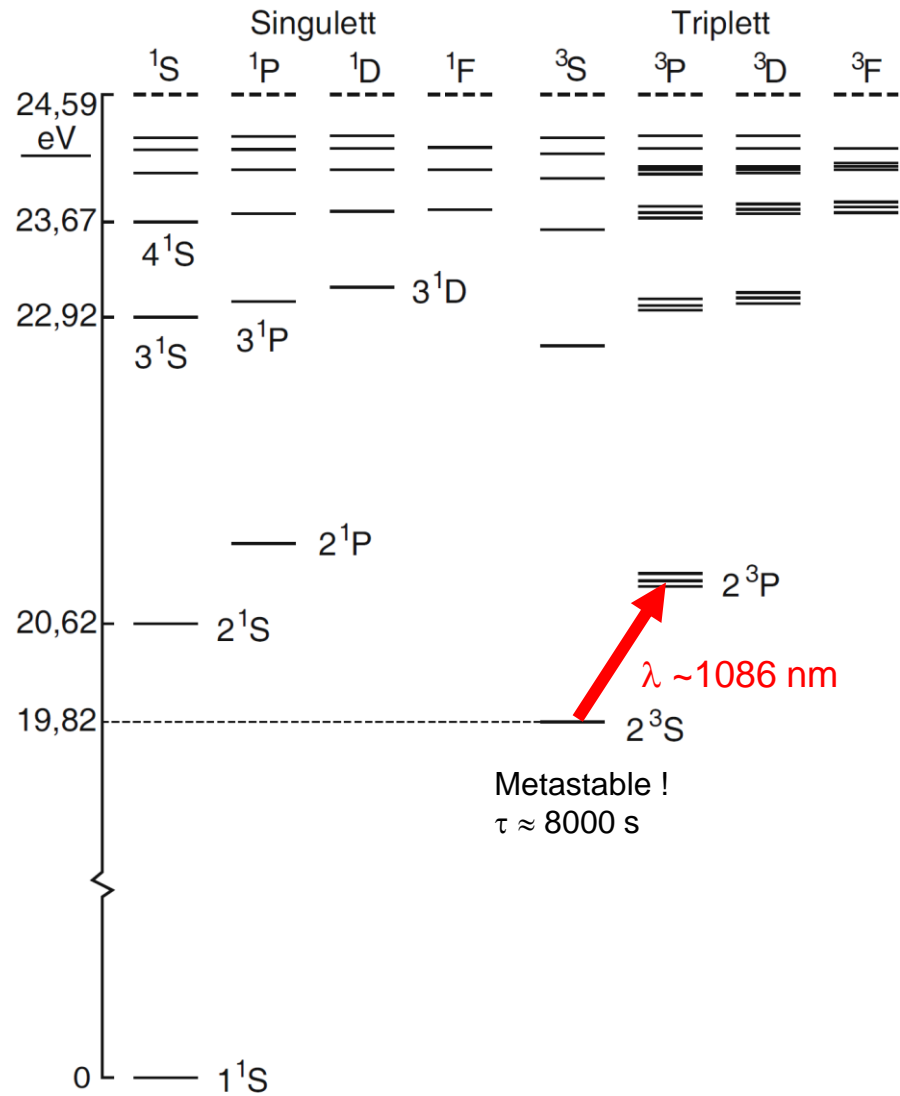
Robert Roth, Thomas Hüther, TU Darmstadt

GFMC: Greens function Monte-Carlo (AV18 + IL7)

R.B. Wiringa, A. Lovato, ANL

Reference radii from eleastic electron scattering have very large uncertainties for both stable isotopes.

Helium



Status of Atomic Theory



Helium

PHYSICAL REVIEW A **103**, 042809 (2021)

Complete $\alpha^7 m$ Lamb shift of helium triplet states

Vojtěch Patkóš¹, Vladimir A. Yerokhin², and Krzysztof Pachucki³

By comparing the theoretical predictions with high-precision experimental results (particularly, the $2^3S - 2^3P$ transition energy), one can determine R . The present theoretical accuracy is, in principle, sufficient for a determination of the nuclear radius with an accuracy of about 1%.

Transition	Theory (MHz)	Difference
$2^3S - 3^3D_1$	786 823 849.540 (52) ^a	-0.462 (76)
$2^3P_0 - 3^3D_1$	510 059 754.863 (16) ^{a,b}	-0.489 (32)
$2^3P - 2^3S$	276 736 495.620 (54)	0.020 (54)

Helium-Like Systems

QED calculations of energy levels of helium-like ions with $5 \leq Z \leq 30$

Vladimir A. Yerokhin,¹ Vojtěch Patkóš,² and Krzysztof Pachucki³

¹Peter the Great St. Petersburg Polytechnic University, Polytekhnicheskaya 29, 195251 St. Petersburg, Russia

²Faculty of Mathematics and Physics, Charles University, Ke Karlovu 3, 121 16 Prague 2, Czech Republic

³Faculty of Physics, University of Warsaw, Pasteura 5, 02-093 Warsaw, Poland

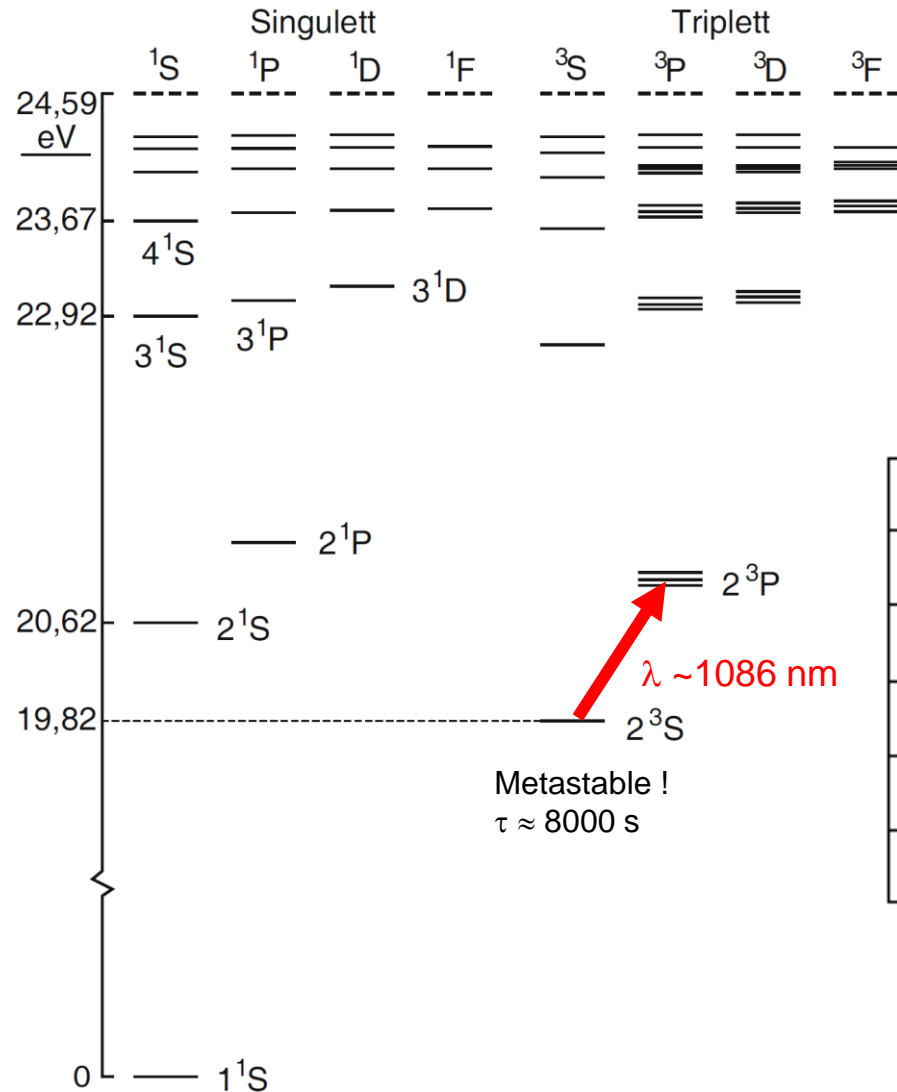
(Dated: June 29, 2022)

arXiv:2206.14161

TABLE X. Comparison of theoretical and experimental $n = 2$ intrashell transition energies, in cm^{-1} .

Z	Theory	Experiment	Difference	Ref.
$2^3S_1 - 2^3P_0$				
5	35 393.6211 (49)	35 393.627 (13)	-0.006 (13)	[47]
	35 393.628 (14) ^a			
$2^3S_1 - 2^3P_2$				
5	35 430.0876 (22)	35 430.084 (9)	0.004 (9)	[47]
	35 430.088 (14) ^a			
$2^3P_0 - 2^3P_1$				
7	8.706 (54)	8.6707 (7)	0.035 (54)	[44]
	8.675 (21) ^a			
	8.6731 (67) ^d			

Helium-Like Systems



Tab. 1: Transition wavelength of different fine structure components of the $1s2s\ ^3S_1 \rightarrow 1s2s\ ^3P_J$ transitions in He-like ions.

Property	Be ²⁺	B ³⁺	C ⁴⁺
Lifetime $\tau\ 1s2s\ ^3S_1$ (ms)	1780	149	21
$\lambda\ ^3S_1 \rightarrow ^3P_1$ (nm)	372.398	282.666	227.863
$\lambda\ ^3S_1 \rightarrow ^3P_0$ (nm)	372.245	282.537	227.797
$\lambda\ ^3S_1 \rightarrow ^3P_2$ (nm)	372.190	282.246	227.159
Laser System	TiSa × 2	Dye × 2	TiSa × 4

Status of Atomic Theory



Helium

PHYSICAL REVIEW A **103**, 042809 (2021)

Complete $\alpha^7 m$ Lamb shift of helium triplet states

Vojtěch Patkóš¹, Vladimir A. Yerokhin², and Krzysztof Pachucki³

By comparing the theoretical predictions with high-precision experimental results (particularly, the $2^3S - 2^3P$ transition energy), one can determine R . The present theoretical accuracy is, in principle, sufficient for a determination of the nuclear radius with an accuracy of about 1%.

Transition	Theory (MHz)	Difference
$2^3S - 3^3D_1$	786 823 849.540 (52) ^a	-0.462 (76)
$2^3P_0 - 3^3D_1$	510 059 754.863 (16) ^{a,b}	-0.489 (32)
$2^3P - 2^3S$	276 736 495.620 (54)	0.020 (54)

Helium-Like Systems

QED calculations of energy levels of helium-like ions with $5 \leq Z \leq 30$

Vladimir A. Yerokhin,¹ Vojtěch Patkóš,² and Krzysztof Pachucki³

¹Peter the Great St. Petersburg Polytechnic University, Polytekhnicheskaya 29, 195251 St. Petersburg, Russia

²Faculty of Mathematics and Physics, Charles University, Ke Karlovu 3, 121 16 Prague 2, Czech Republic

³Faculty of Physics, University of Warsaw, Pasteura 5, 02-093 Warsaw, Poland

(Dated: June 29, 2022)

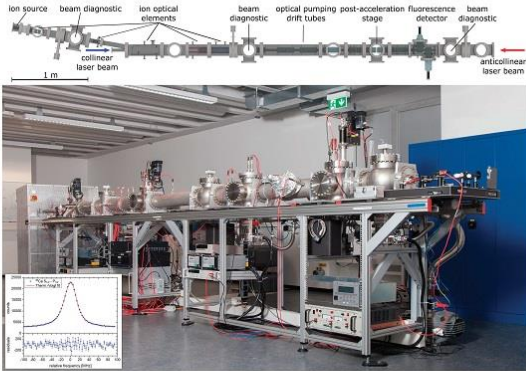
arXiv:2206.14161

TABLE X. Comparison of theoretical and experimental $n = 2$ intrashell transition energies, in cm^{-1} .

Z	Theory	Experiment	Difference	Ref.
5	$2^3S_1 - 2^3P_0$	35 393.627 (13)	-0.006 (13)	[47]
	35 393.6211 (49)			
5	$2^3S_1 - 2^3P_2$	35 430.084 (9)	0.004 (9)	[47]
	35 430.0876 (22)			
7	$2^3P_0 - 2^3P_1$	8.6707 (7)	0.035 (54)	[44]
	8.706 (54)			
	8.675 (21) ^a			
	8.6731 (67) ^d			

Measuring principle @ COALA, TU Darmstadt

scitation.org/journal/rsi

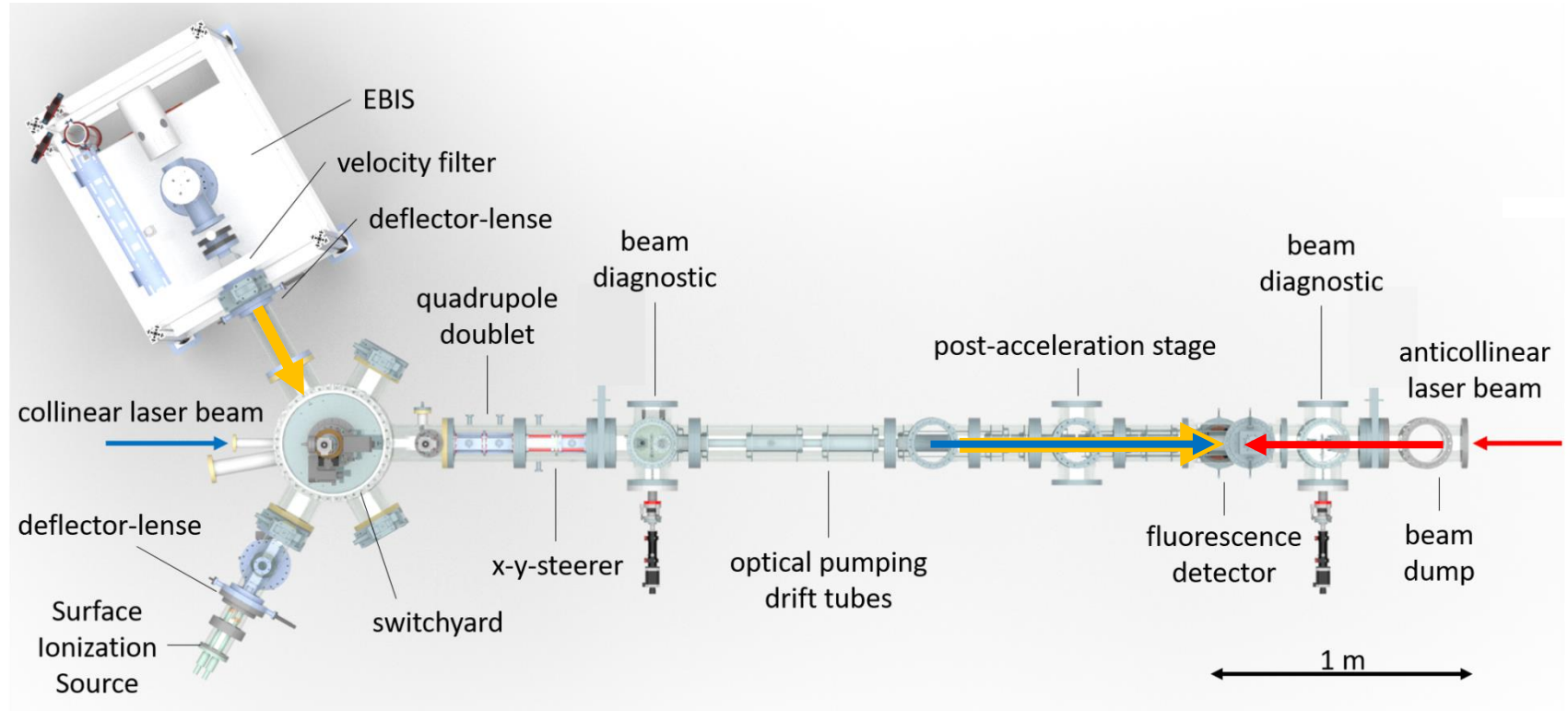


Volume 91, Issue 8, Aug. 2020

A new Collinear Apparatus for Laser Spectroscopy and Applied Science (COALA)

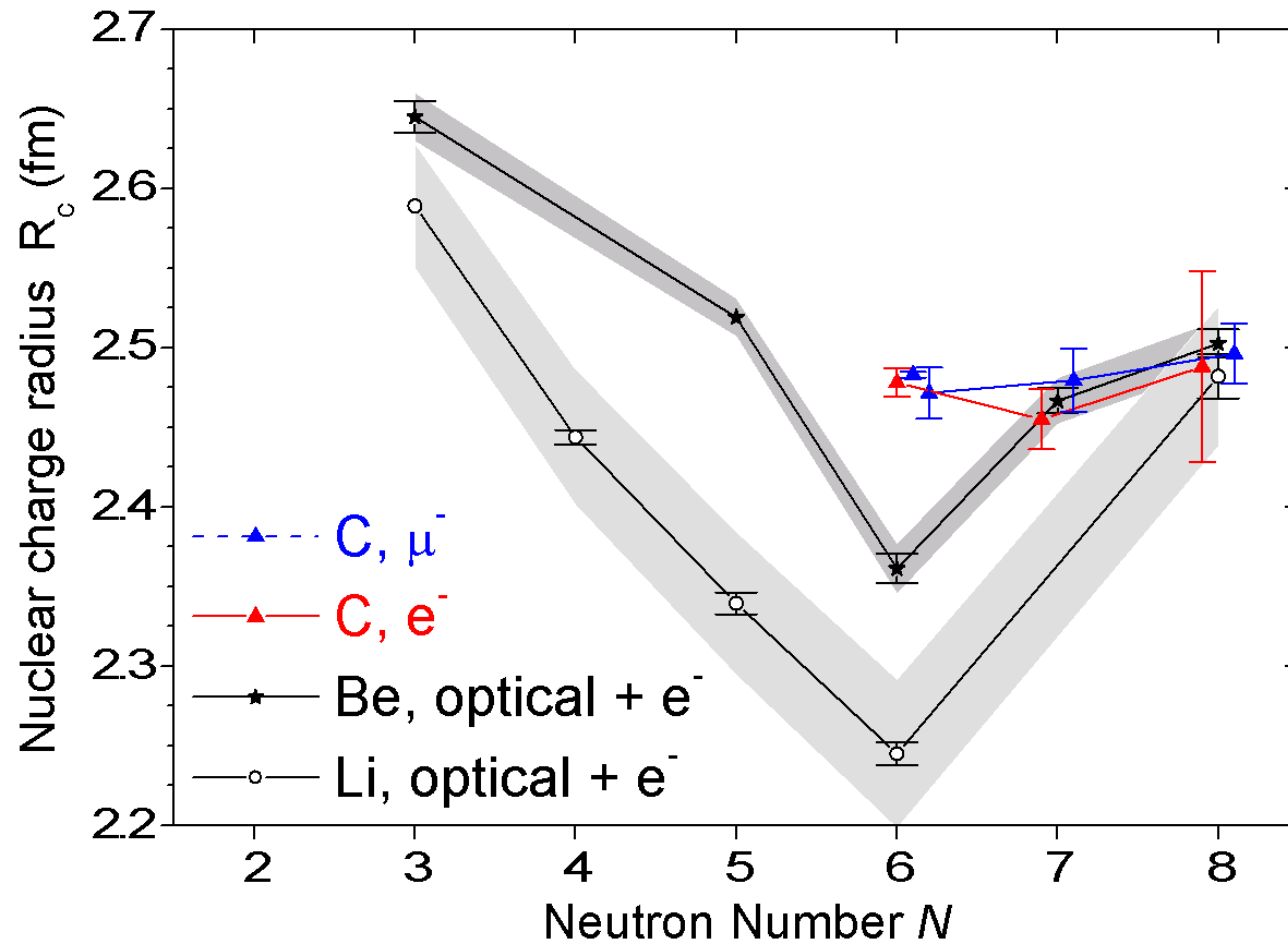
Rev. Sci. Instrum. 91, 081301 (2020); doi.org/10.1063/5.0010903

K. König, J. Krämer, C. Geppert, P. Imgram, B. Maaß, T. Ratajczyk, and W. Nörtershäuser



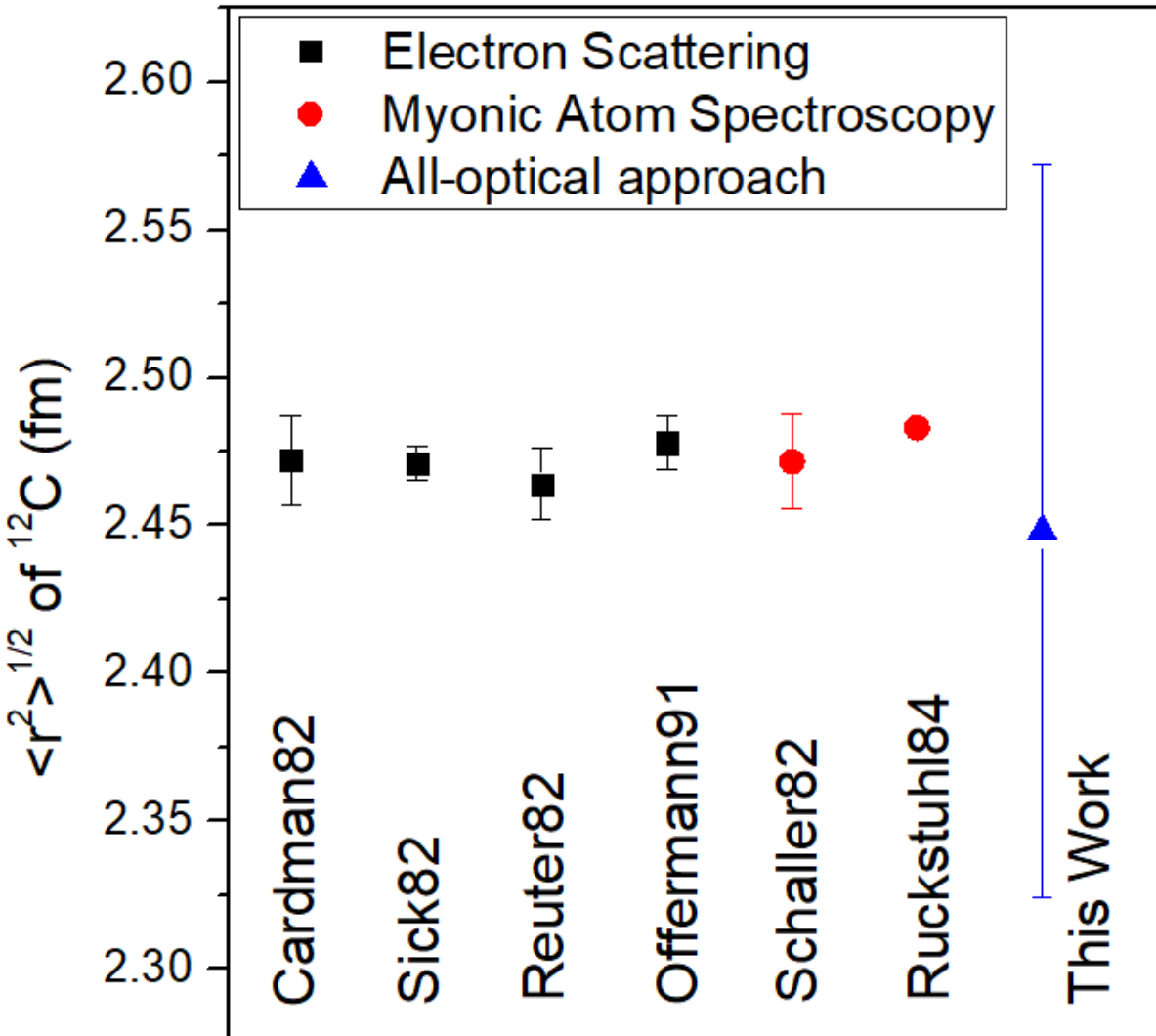
$$\left. \begin{aligned} v_c &= v_0 \gamma (1 + \beta) \\ v_a &= v_0 \gamma (1 - \beta) \end{aligned} \right\} v_c \cdot v_a = v_0^2 \gamma^2 \cdot (1 + \beta)(1 - \beta) = v_0^2$$

^{12}C : Proof of Principle & Test of Theory



- Charge Radius of nucleus very well known
→ Test for theory
- Easy to produce in an EBIS
- Transition wavelength of 227 nm
→ Ti:Sa \times 4 stabilized to frequency comb
- No hyperfine structure → no hyperfine-induced level mixing
- ^{13}C hyperfine structure requires a more elaborated experiment and theory
- Charge radius of ^{13}C can be improved based on the isotope shift method and compared to direct extraction from transition frequency

Extraction of the Nuclear Charge Radius



Current status:

based on Yerokhin, Patkos & Pachucki, arXiv:2206.14161
[Phys. Rev. A **106**, 022815 (2022)]

- uncertainty dominated by theory
 - experimental accuracy sufficient to compete with muonic atom result
- theory needs to be further improved

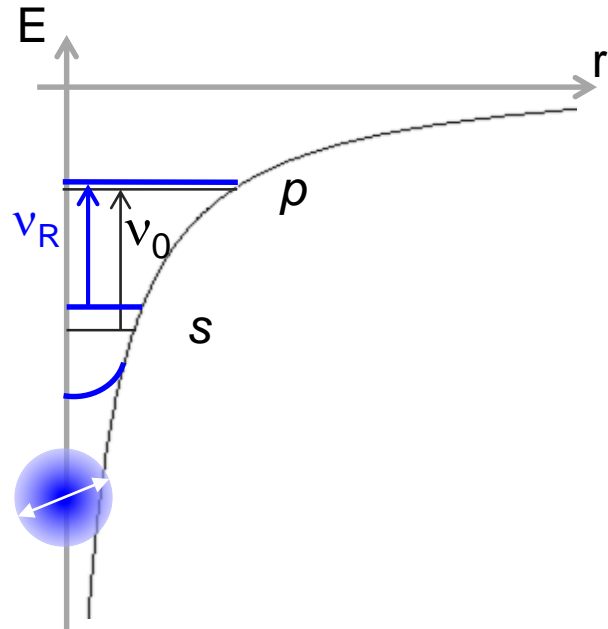
Fine structure splitting:

- measurement will provide $m\alpha^8$ -estimation for helium
- contributes to α -determination

Outlook:

- $^{10,11}\text{B}^{3+}$ and $^{10,11}\text{B}^{2+}$ measurements at KOALA
- ^8B measurements at ANL
- Be^{2+} with external ion injection at KOALA
- Online application for carbon isotopes ??

The Idea of All-Optical Absolute Charge Radii



$$\delta\nu_{FS} = \underbrace{-\frac{Ze^2}{6\epsilon_0} \Delta |\Psi_e(0)|_{i \rightarrow f}^2}_{\text{Electronic Factor} \rightarrow \text{Wavefunction}} \times \langle r_c^2 \rangle$$

$$= F_{i \rightarrow f} \langle r_c^2 \rangle$$

- Measure **transition frequency** ν_R
- Compare with high precision atomic calculation for a point-like nucleus ν_0
- Difference $\nu_R - \nu_0$ is finite-size effect and **proportional to the ms charge radius**
- So far applied **only for H-like systems**, i.e., H, μH and μHe
- Two-electron system requires elaborate QED calculations, which are now in reach
Yerokhin, Patkóš & Pachucki, PRA 98, 032503 (2018)
Patkóš, Yerokhin & Pachucki, PRA 103, 042809 (2021)
- Laser spectroscopy on $^{12,13}\text{C}^{4+}$ performed.

Address He-like systems:
Li⁺, Be²⁺, **B³⁺**, **C⁴⁺**, N⁵⁺

